

# MATH 436 FALL 2012 HOMEWORK 4 SOLUTIONS

DUE NOV. 8 IN CLASS

**Note.** All problem numbers refer to “Updated” version of lecture note.

- **Exercise 3.1.** Consider the Telegrapher’s equation

$$u_{xx} = u_{tt} + \lambda u_t \quad (1)$$

(recall that  $\lambda > 0$ ) over the interval  $x \in [0, L]$  subject to conditions

$$u(0, t) = u(L, t) = 0; \quad u(x, 0) = f(x), \quad u_t(x, 0) = h(x). \quad (2)$$

Use the method of separation of variables to study the limiting behavior of  $u$  as  $t \rightarrow \infty$ .

**Solution.** We separate the variables: Substitute  $X(x)T(t)$  into the equation:

$$T(t)X''(x) = T''(t)X(x) + \lambda T'(t)X(x). \quad (3)$$

Divide both sides by  $X(x)T(t)$  we get

$$\frac{X''}{X} = \frac{T''}{T} + \lambda \frac{T'}{T} = c^2 K. \quad (4)$$

Thus we get the eigenvalue problem:

$$X'' - KX = 0, \quad X(0) = X(L) = 0. \quad (5)$$

and the equation for  $T$ :

$$T'' + \lambda T' - KT = 0. \quad (6)$$

We solve the eigenvalue problem to obtain

$$K_n = -\left(\frac{n\pi}{L}\right)^2, \quad X_n = \sin \frac{n\pi x}{L}, \quad n = 1, 2, 3, \dots \quad (7)$$

Now we expand

$$f(x) = \sum_{n=1}^{\infty} f_n \sin \frac{n\pi x}{L}, \quad h(x) = \sum_{n=1}^{\infty} h_n \sin \frac{n\pi x}{L} \quad (8)$$

where

$$f_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi x}{L} dx, \quad h_n = \frac{2}{L} \int_0^L h(x) \sin \frac{n\pi x}{L} dx. \quad (9)$$

Next we solve the  $T_n$  equation:

$$T_n'' + \lambda T_n' + \left(\frac{n\pi}{L}\right)^2 T_n = 0. \quad (10)$$

Its auxiliary equation is

$$r^2 + \lambda r + \left(\frac{n\pi}{L}\right)^2 = 0 \implies r_{1,2} = \frac{-\lambda \pm \sqrt{\lambda^2 - 4\left(\frac{n\pi}{L}\right)^2}}{2}. \quad (11)$$

Now it is clear that  $\Re r_{1,2} < 0$  for all possible values of  $\lambda$ . As a consequence the solution would read

$$u(x, t) = \sum e^{-\alpha_n t} \left[ a_n \cos \frac{n \pi x}{L} + b_n \sin \frac{n \pi x}{L} \right] \quad (12)$$

with all  $\alpha_n > 0$ . Therefore  $u(x, t) \rightarrow 0$  as  $t \rightarrow \infty$ .

**(To Cody: The following is not required)**

More specifically, we have

- If  $\lambda < \frac{2\pi}{L}$ , then  $\lambda^2 - 4 \left( \frac{n\pi}{L} \right)^2 < 0$  for all  $n \in \mathbb{N}$ . Consequently the real part of all  $\alpha_n$  is  $-\frac{\lambda}{2}$ . Thus  $u(x, t) \rightarrow 0$  like  $e^{-\lambda t/2}$ .
- If  $\lambda = \frac{2\pi}{L}$ ,  $u(x, t) \rightarrow 0$  like  $e^{-\lambda t/2} t$  (since now we have a repeated root).
- If  $\lambda > \frac{2\pi}{L}$ , then

$$\Re \alpha_n \geq \frac{-\lambda + \sqrt{\lambda^2 - 4 \left( \frac{\pi}{L} \right)^2}}{2} \quad (13)$$

which means  $u(x, t) \rightarrow 0$  like  $\exp \left[ \frac{-\lambda + \sqrt{\lambda^2 - 4 \left( \frac{\pi}{L} \right)^2}}{2} t \right]$ .

- **Exercise 3.5.** Consider the boundary value problem for  $u(x, y)$  in the annular region:

$$u_{xx} + u_{yy} = 0 \quad \rho^2 < x^2 + y^2 < 1; \quad u(x, y) = \begin{cases} f & x^2 + y^2 = \rho^2 \\ g & x^2 + y^2 = 1 \end{cases}. \quad (14)$$

Obtain the formula for the solution using separation of variables.

**Solution.** The solution process is identical to that of Laplace equation in the disc until we need to solve the  $R$  equation:

$$r^2 R'' + r R' - n^2 R = 0, \quad R(\rho_1) = f_{ni}, \quad R(\rho_2) = g_{ni}. \quad (15)$$

Here  $f_{ni}, g_{ni}$  comes from expansions

$$f(\theta) = \sum_{n=0}^{\infty} [f_{n1} \cos n\theta + f_{n2} \sin n\theta], \quad g(\theta) = \sum_{n=0}^{\infty} [g_{n1} \cos n\theta + g_{n2} \sin n\theta]. \quad (16)$$

When  $n=0$ , the general solution is

$$R(r) = C_1 + C_2 \log r \quad (17)$$

which leads to

$$C_1 + C_2 \log \rho = f_{0i}, \quad C_1 = g_{0i}. \quad (18)$$

Consequently

$$C_2 = (f_{0i} - g_{0i}) / (\log \rho), \quad C_1 = g_{0i}. \quad (19)$$

When  $n > 0$  the general solution is

$$R(r) = C_1 r^n + C_2 r^{-n}. \quad (20)$$

Then the boundary conditions lead to

$$C_1 \rho^n + C_2 \rho^{-n} = f_{ni}, \quad C_1 + C_2 = g_{ni}. \quad (21)$$

The solution is

$$C_1 = \frac{f_{ni} \rho^n - g_{ni}}{\rho^{2n} - 1}, \quad C_2 = \rho^n \frac{g_{ni} \rho^n - f_{ni}}{\rho^{2n} - 1}. \quad (22)$$

Therefore the solution to the problem is

$$u(r, \theta) = \left[ g_{01} + \frac{f_{01} - g_{01}}{\log \rho} \right] + \sum_{n=1}^{\infty} \left[ \left( \frac{f_{n1} \rho^n - g_{n1}}{\rho^{2n} - 1} r^n + \left( \frac{\rho}{r} \right)^n \frac{g_{n1} \rho^n - f_{n1}}{\rho^{2n} - 1} \right) \cos(n\theta) + \left( \frac{f_{n2} \rho^n - g_{n2}}{\rho^{2n} - 1} r^n + \left( \frac{\rho}{r} \right)^n \frac{g_{n2} \rho^n - f_{n2}}{\rho^{2n} - 1} \right) \sin(n\theta) \right]. \quad (23)$$

- **Exercise 3.8.** Consider the equation

$$r^2 R'' + r R' + (\lambda r^2 - n^2) R = 0, \quad R(0), R'(0) \text{ bounded}, R(1) = 0. \quad (24)$$

- Prove that there are no negative eigenvalues (along the way you will see why the boundedness of  $R'(0)$  is needed.)
- Prove that  $\lambda = 0$  is not an eigenvalue.
- Let  $\lambda_k \neq \lambda_l$  be eigenvalues, prove that the eigenfunctions  $R_k(r), R_l(r)$  satisfy

$$\int_0^1 R_k(r) R_l(r) r \, dr = 0. \quad (25)$$

(Hint: For a) and c), write the equation as

$$\frac{1}{r} \left[ (r R')' - \frac{n^2}{r} R \right] + \lambda R = 0 \quad (26)$$

first.)

**Proof.** We follow the hint and write the equation as

$$\frac{1}{r} \left[ (r R')' - \frac{n^2}{r} R \right] + \lambda R = 0 \quad (27)$$

- $\lambda < 0$  cannot be eigenvalue.

Let  $\lambda$  be an eigenvalue with corresponding eigenfunction  $R$ . Multiply both sides of the equation by  $r R$  and integrate from 0 to 1 we have

$$\begin{aligned} 0 &= \int_0^1 \left\{ \frac{1}{r} \left[ (r R')' - \frac{n^2}{r} R \right] + \lambda R \right\} R r \, dr \\ &= \int_0^1 \left[ (r R')' - \frac{n^2}{r} R \right] R \, dr + \int_0^1 \lambda R^2 \, dr \\ &= \int_0^1 (r R')' R \, dr - \int_0^1 \left( \frac{n^2}{r} - \lambda \right) R^2 \, dr \\ &= r R' R|_0^1 - \int_0^1 r (R')^2 - \int_0^1 \left( \frac{n^2}{r} - \lambda \right) R^2 \, dr. \end{aligned} \quad (28)$$

As  $R(1) = 0$ ,  $R'(0)$   $R(0)$  bounded, the boundary terms vanish. Consequently

$$\int_0^1 \left( \frac{n^2}{r} - \lambda \right) R^2 dr \leq 0 \quad (29)$$

which is not possible for nonzero  $R$  if  $\lambda < 0$ .

**Remark 1.** Note that in fact the above also shows  $\lambda = 0$  is not an eigenvalue.

- $\lambda = 0$  is not an eigenvalue.

When  $\lambda = 0$  the equation becomes

$$r^2 R'' + r R' - n^2 R = 0 \quad (30)$$

which general solution is

$$R(r) = \begin{cases} C_1 + C_2 \log r & n = 0 \\ C_1 r^n + C_2 r^{-n} & n \neq 0 \end{cases}. \quad (31)$$

The boundary condition at 0 leads to  $C_2 = 0$ . Then the other boundary condition  $R(1) = 0$  requires  $C_1 = 0$ .

- If

$$\frac{1}{r} \left[ (r R_k')' - \frac{n^2}{r} R_k \right] + \lambda_k R_k = 0, \quad \frac{1}{r} \left[ (r R_l')' - \frac{n^2}{r} R_l \right] + \lambda_l R_l = 0 \quad (32)$$

then we can multiply the first equation by  $r R_l$  and integrate from 0 to 1:

$$\begin{aligned} 0 &= \int_0^1 \left[ (r R_k')' - \frac{n^2}{r} R_k \right] R_l + \lambda_k \int_0^1 R_k R_l r dr \\ &= \int_0^1 (r R_k')' R_l - \int_0^1 \frac{n^2}{r} R_k R_l + \lambda_k \int_0^1 R_k R_l r dr \\ &= [r R_k' R_l] \Big|_0^1 - \int_0^1 r R_k' R_l' - \int_0^1 \frac{n^2}{r} R_k R_l + \lambda_k \int_0^1 R_k R_l r dr \\ &= 0 - [r R_k R_l'] \Big|_0^1 + \int_0^1 R_k (r R_l')' - \int_0^1 \frac{n^2}{r} R_k R_l + \lambda_k \int_0^1 R_k R_l r dr \\ &= \int_0^1 \left[ (r R_l')' - \frac{n^2}{r} R_l \right] R_k dr + \lambda_k \int_0^1 R_k R_l r dr \\ &= (\lambda_k - \lambda_l) \int_0^1 R_k R_l r dr. \end{aligned} \quad (33)$$

Since  $\lambda_k - \lambda_l \neq 0$  we conclude

$$\int_0^1 R_k R_l r dr = 0 \quad (34)$$

which is the desired orthogonality relation.  $\square$