# Asymptotics of One-Dimensional Lévy Approximations 

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#### Abstract

For arbitrary Borel probability measures on the real line, necessary and sufficient conditions are presented that characterize best purely atomic approximations relative to the classical Lévy probability metric, given any number of atoms, and allowing for additional constraints regarding locations or weights of atoms. The precise asymptotics (as the number of atoms goes to infinity) of the approximation error is identified for the important special cases of best uniform (i.e. all atoms having equal weight) and best (i.e. unconstrained) approximations, respectively. When compared to similar results known for other probability metrics, the results for Lévy approximations are more complete and require fewer assumptions.


Keywords Best (uniform) approximation • Lévy probability metric • Inverse function • Inverse measure • Approximation error • Asymptotic point distribution

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## 1 Introduction

Let $\mathcal{P}$ be the set of all Borel probability measures on the real line, and denote the support of $\mu \in \mathcal{P}$ by supp $\mu$. For each positive integer $n$, let $\mathcal{P}_{n}^{*}=\{\mu \in \mathcal{P}$ : \# supp $\mu \leq n\}$. Recall that $\mathcal{P}$ endowed with the topology of weak convergence is a Polish space that contains $\mathcal{P}_{\infty}^{*}:=\bigcup_{n} \mathcal{P}_{n}^{*}=\{\mu \in \mathcal{P}: \# \operatorname{supp} \mu<\infty\}$ as a dense subspace [11, Ch. 11]. Many different metrics (and metric-like quantities [15]) on $\mathcal{P}$ or parts thereof have been studied extensively, as they play important roles in probability theory and statistics [30,31]. Given a specific probability metric $d$ and $\mu \in \mathcal{P} \backslash \mathcal{P}_{\infty}^{*}$, it is natural to ask whether there exists, for every $n$, a best $d$-approximation $\delta_{\bullet}^{\bullet}, n$ of $\mu$ in $\mathcal{P}_{n}^{*}$, i.e. $d\left(\mu, \delta_{\bullet}^{\bullet}, n\right)=\inf \left\{d(\mu, v): v \in \mathcal{P}_{n}^{*}\right\}$, perhaps with additional desirable properties

[^0]such as, e.g. all atoms having equal weight; see Sect. 2 for precise terminology and notation. Provided they exist, how can such best $d$-approximations be characterized and found systematically? How fast do they converge to $\mu$, i.e. at what rate does the approximation error $d\left(\mu, \delta_{\bullet}^{\bullet, n}\right)$ tend to 0 as $n \rightarrow \infty$ ? Questions like these, regarding the approximation in $\mathcal{P}$ by elements of $\mathcal{P}_{\infty}^{*}$, continue to attract interest in a wide variety of contexts; see, e.g. $[5,6,9,10,23,24,28]$ and the many references therein.

Denoting the distribution function of $\mu \in \mathcal{P}$ by $F_{\mu}$, that is, $\left.\left.F_{\mu}(x)=\mu(]-\infty, x\right]\right)$ for all $x \in \mathbb{R}$, recall the Kantorovich (or Wasserstein; cf. [6, p. 4] and [15]) metric, given by

$$
\begin{equation*}
d_{\mathrm{W}}(\mu, \nu)=\int_{\mathbb{R}}\left|F_{\mu}(x)-F_{v}(x)\right| \mathrm{d} x=\int_{[0,1]}\left|F_{\mu}^{-1}(y)-F_{v}^{-1}(y)\right| \mathrm{d} y \tag{1.1}
\end{equation*}
$$

where $F_{\mu}^{-1}$ is an inverse of $F_{\mu}$; see Sect. 2 for details. Note that strictly speaking $d_{\mathrm{W}}$ is not defined on all of $\mathcal{P} \times \mathcal{P}$, but only on $\mathcal{P}_{1} \times \mathcal{P}_{1}$, with $\mathcal{P}_{1}=\{\mu \in \mathcal{P}$ : $\left.\int_{\mathbb{R}}|x| \mathrm{d} \mu(x)<+\infty\right\} \supset \mathcal{P}_{\infty}^{*}$. The metric space $\left(\mathcal{P}_{1}, d_{\mathrm{W}}\right)$ is complete and separable, though its metric topology is finer than the subspace topology inherited from $\mathcal{P}$. Due to its simplicity and functional-analytic flavour, the metric $d_{\mathrm{W}}$ figures prominently in many applied areas, e.g. image compression, signal processing, mathematical finance and optimal transport [23,27,32,35,38]. A vast literature exists addressing the basic questions mentioned earlier relative to $d_{\mathrm{W}}$, as well as many generalizations thereof, notably to multi-dimensional settings [6,13,21,22,28].

Another important notion of distance, the Prokhorov metric is given by

$$
\begin{equation*}
d_{\mathrm{P}}(\mu, \nu)=\inf \left\{y \in \mathbb{R}^{+}: \mu(B) \leq \nu\left(B^{y}\right)+y \forall \text { Borel sets } B \subset \mathbb{R}\right\} \quad \forall \mu, v \in \mathcal{P}, \tag{1.2}
\end{equation*}
$$

where $B^{y}=\{x \in \mathbb{R}: \operatorname{dist}(x, B)<y\}$. Note that $d_{\mathrm{P}}$ is defined on all of $\mathcal{P} \times \mathcal{P}$, unlike $d_{\mathrm{W}}$, and metrizes precisely the topology of weak convergence [11,15]. Also, $d_{\mathrm{P}}(\mu, \nu) \leq 1$ for all $\mu, \nu \in \mathcal{P}$. A general theory of best $d_{\mathrm{P}}$-approximation in $\mathcal{P}$ by elements of $\mathcal{P}_{\infty}^{*}$ has been initiated in [20], where the authors observe that some aspects of the theory are "more difficult [than the corresponding theory for $d_{\mathrm{w}}$ ] ...mainly due to the lack of suitable scaling properties [of $\left.d_{\mathrm{P}}\right]$ ".

In a spirit similar to $[18,37]$, this article addresses the approximation problem relative to the classical Lévy metric,

$$
\begin{equation*}
d_{1}(\mu, \nu)=\inf \left\{y \in \mathbb{R}^{+}: F_{\mu-}(\cdot-y)-y \leq F_{\nu} \leq F_{\mu}(\cdot+y)+y\right\} \quad \forall \mu, \nu \in \mathcal{P} \tag{1.3}
\end{equation*}
$$

where $F_{\mu-}(x)=\lim _{\varepsilon \downarrow 0} F_{\mu}(x-\varepsilon)=\mu(]-\infty, x[)$. Note that $d_{1} \leq 1$, similarly to $d_{\mathrm{p}}$. The values of $d_{\mathrm{W}}, d_{\mathrm{P}}$ and $d_{1}$ are not completely unrelated, since $d_{1}(\mu, \nu) \leq d_{\mathrm{P}}(\mu, \nu) \leq$ $\sqrt{d_{\mathrm{W}}(\mu, v)}$ for all $\mu, v \in \mathcal{P}_{1}$; see, e.g. [4,11,15]. When compared to $d_{\mathrm{W}}$ and $d_{\mathrm{P}}$, the metric $d_{1}$ is particularly attractive: On the one hand, it is a bona fide metric [3, p. 100] metrizing the topology of weak convergence on all of $\mathcal{P}$ (similar to $d_{\mathrm{P}}$, but unlike $d_{\mathrm{W}}$ ). On the other hand, its definition (1.3) is considerably easier to work with than (1.2). Although computing $d_{1}$ for concrete problems may still be "not easy" [15, p. 423] (cf. also [34]), especially when compared to (1.1), the main (asymptotic) results of this article suggest that nevertheless $d_{1}$ often is more benign than both $d_{\mathrm{W}}$ and $d_{\mathrm{P}}$, in that
fewer assumptions (or no assumptions at all, as in Theorem 4.1 and Proposition 4.3) are needed to draw analogous or perhaps even stronger conclusions. With all technical details deferred to later sections, this is illustrated here for two familiar (absolutely continuous) distributions-standard normal and 1-Pareto.

Let $\mu$ be the standard normal distribution. By a celebrated asymptotic result for best $d_{\mathrm{W}}$-approximations [18, Thm. 6.2],

$$
\begin{equation*}
\lim _{n \rightarrow \infty} n d_{\mathrm{W}}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=\sqrt{\frac{\pi}{2}}=1.253 \tag{1.4}
\end{equation*}
$$

whereas by [20, Ex. 5.1],

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \frac{n}{\sqrt{\log n}} d_{\mathrm{p}}\left(\mu, \delta_{\bullet}^{\bullet}, n\right)=\sqrt{2} \tag{1.5}
\end{equation*}
$$

Note that (1.4) yields the faster decay of the approximation (or quantization) error $d\left(\mu, \delta_{\bullet}^{\bullet, n}\right)$, whereas only (1.5) involves a probability metric that actually metrizes the topology of weak convergence. As it turns out, for the Lévy metric these two desirable properties can be achieved simultaneously: Theorem 4.1, one of the main results of this article, yields

$$
\begin{equation*}
\lim _{n \rightarrow \infty} n d_{1}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=-\sqrt{\frac{\pi}{2}} \mathrm{Li}_{1 / 2}\left(-\frac{1}{\sqrt{2 \pi}}\right)=0.3931 \tag{1.6}
\end{equation*}
$$

where $\mathrm{Li}_{1 / 2}$ denotes the polylogarithm of order $\frac{1}{2}$.
An interesting variant of (1.4)-(1.6) considers best uniform approximations of $\mu \in \mathcal{P}$, that is, best approximations of $\mu$ by $v \in \mathcal{P}_{n}^{*}$, subject to the additional requirement that $n v(\{x\})$ is a (positive) integer for every $x \in \operatorname{supp} \nu$. Best uniform (or, more generally, best constrained) approximations have recently attracted considerable interest, not least in view of potential applications in stochastic processes and differential equations [ $7,8,16,17,36,37$ ]; they may also be viewed as deterministic analogues of (random) empirical measures [6,9,14]. With $\delta_{\bullet}^{u_{n}}$ denoting a best uniform $d$-approximation of $\mu$, trivially $d\left(\mu, \delta_{\bullet}^{\bullet, n}\right) \leq d\left(\mu, \delta_{\bullet}^{u_{n}}\right)$. For $\mu$ being the standard normal distribution, [37, Ex. 5.18] reports that

$$
d_{\mathrm{W}}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\mathcal{O}\left(\frac{\sqrt{\log n}}{n}\right) \quad \text { as } n \rightarrow \infty
$$

and this bound is sharp; cf. also [7,16]. Although the authors do not know of any analogous result regarding best uniform $d_{\mathrm{p}}$-approximations, (1.5) makes it clear that $d_{\mathrm{P}}\left(\mu, \delta_{\bullet}^{u_{n}}\right)$ is at least $\mathcal{O}\left(n^{-1} \sqrt{\log n}\right)$ as $n \rightarrow \infty$, if not larger. By contrast, Theorem 3.3, another main result of this article, simply yields

$$
\lim _{n \rightarrow \infty} n d_{1}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\frac{1}{2}
$$

which represents a faster and more precise rate than its $d_{\mathrm{W}}$ - and $d_{\mathrm{p}}$-counterparts.

For a second illustrative example, let $\mu$ be the 1-Pareto distribution, i.e. $F_{\mu}(x)=$ $1-x^{-1}$ for all $x \geq 1$. Since $\mu \notin \mathcal{P}_{1}$, clearly $\mu$ is not amenable to $d_{\mathrm{W}}$-approximation, whereas [20, Thm. 5.2] yields

$$
\lim _{n \rightarrow \infty} \sqrt{n} d_{\mathrm{p}}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=\frac{1}{\sqrt{2}}
$$

For the Lévy metric, this article again provides faster, more precise rates, namely

$$
\begin{equation*}
n d_{1}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\frac{1}{2}-\frac{1}{8} n^{-2}+\mathcal{O}\left(n^{-3}\right) \quad \text { as } n \rightarrow \infty \tag{1.7}
\end{equation*}
$$

as well as

$$
\begin{equation*}
n d_{1}\left(\mu, \delta_{\bullet}^{\bullet}, n\right)=\frac{\pi}{8}+\frac{\pi^{2}(6-\pi)}{3 \times 2^{10}} n^{-2}+\mathcal{O}\left(n^{-3}\right) \quad \text { as } n \rightarrow \infty . \tag{1.8}
\end{equation*}
$$

Thus, the results of this article make the case that although the Lévy metric $d_{1}$, unlike $d_{\mathrm{W}}$ and $d_{\mathrm{P}}$, does not extend to higher dimensions in a straightforward way, its usage for one-dimensional probabilities often leads to simpler and stronger results.

This article is organized as follows. Section 2 first introduces all required terminology and notation and then reviews basic facts pertaining to approximations in $\mathcal{P}$ relative to the Lévy metric. Utilizing the latter, Sects. 3 and 4 specifically study best uniform and best (unconstrained) approximations, respectively, and in particular the asymptotics of the approximation error as $n \rightarrow \infty$. Also, under a mild assumption the atoms of (asymptotically) best approximations conform to an asymptotic point distribution, as shown by Theorem 4.5.

## 2 Lévy Probability Metrics

This section reviews basic facts regarding the approximation in $\mathcal{P}$ by measures with finite support, relative to the Lévy probability metric(s). The stated results are straightforward extensions of [4,36], and the reader is referred to these references for further details and elementary proofs. The following, mostly standard notations are used throughout. The sets of all positive integers, nonnegative integers, integers, positive real numbers and real numbers are denoted $\mathbb{N}, \mathbb{N}_{0}, \mathbb{Z}, \mathbb{R}^{+}$and $\mathbb{R}$, respectively. Numerical values of real numbers are displayed to four correct significant decimal digits. For every $x \in \mathbb{R}$ and non-empty $A \subset \mathbb{R}, \operatorname{dist}(x, A)=\inf _{a \in A}|x-a|, \operatorname{diam} A=\sup _{a, b \in A}|a-b|$, and $\mathbf{1}_{A}$ is the indicator function of $A$; also, $\lfloor x\rfloor$ denotes the largest integer not larger than $x$, and the cardinality of $A$ is \#A. If the domain of a function $f$ contains $A$, then $f(A)=\{f(a): a \in A\}$. Lebesgue measure on the real line is denoted $\lambda$.

Since non-decreasing functions play a crucial role in what follows, first a few basic properties of such functions are recorded. Throughout, denote by $\overline{\mathbb{R}}=\mathbb{R} \cup\{-\infty,+\infty\}$ the extended real line with its usual order and topology, and by $\mathcal{F}$ the family of all functions $f: \mathbb{R} \rightarrow \overline{\mathbb{R}}$ that are non-decreasing and right-continuous. Given $f \in \mathcal{F}$, let $f( \pm \infty)=\lim _{x \rightarrow \pm \infty} f(x) \in \overline{\mathbb{R}}$, and for every $x \in \overline{\mathbb{R}}$ let $f_{-}(x)=\lim _{\varepsilon \downarrow 0} f(x-\varepsilon)$.

Note that $f_{-}(x) \leq f(x) \leq f_{-}(y)$ whenever $x<y$; in particular, $f_{-}(x)=f(x)$ if and only if $f$ is continuous at $x$. With every $f \in \mathcal{F}$ associate its (upper) inverse function $f^{-1}: \mathbb{R} \rightarrow \overline{\mathbb{R}}$ given by

$$
f^{-1}(x)=\sup \{y \in \mathbb{R}: f(y) \leq x\} \quad \forall x \in \mathbb{R} ;
$$

here and throughout the convention $\sup \varnothing=-\infty($ and $\inf \varnothing=+\infty)$ is adhered to. Importantly, $\mathcal{F}$ is closed under inversion and composition.

Proposition 2.1 Let $f, g \in \mathcal{F}$. Then, $f^{-1} \circ g \in \mathcal{F}$, and $\left(f^{-1}\right)^{-1}=f$.
Given $f, g \in \mathcal{F}$ and $\epsilon>0$, let

$$
d_{\epsilon}(f, g)=\inf \left\{y \in \mathbb{R}^{+}: f_{-}(\cdot-y / \epsilon)-y \leq g \leq f(\cdot+y / \epsilon)+y\right\} \in[0,+\infty] .
$$

Motivated for $\epsilon=1$ by (1.3), this definition enables a unified treatment of all $\epsilon$-Lévy probability metrics later in this section. It is readily checked that $d_{\epsilon}$ indeed satisfies the axioms of a metric on $\mathcal{F}$, except that $d_{\epsilon}(f, g)$ may equal $+\infty$. Also, $d_{\epsilon}$ is compatible with inversion.

Proposition 2.2 Let $f, g \in \mathcal{F}$ and $\epsilon>0$. Then, $d_{\epsilon}\left(f^{-1}, g^{-1}\right)=\epsilon d_{1 / \epsilon}(f, g)$.
Given $f, g \in \mathcal{F}$, note that $\epsilon \mapsto d_{\epsilon}(f, g)$ is non-decreasing and continuous on $\mathbb{R}^{+}$. Consequently, the limits of $d_{\epsilon}(f, g)$ exist as $\epsilon \rightarrow 0$ or $\epsilon \rightarrow+\infty$. For instance, if $f, g \in \mathcal{F}$ are bounded, then simply

$$
\begin{aligned}
\lim _{\epsilon \rightarrow 0} d_{\epsilon}(f, g) & =\lim \sup _{|x| \rightarrow+\infty}|f(x)-g(x)| \\
& =\max \{|f(-\infty)-g(-\infty)|,|f(+\infty)-g(+\infty)|\}
\end{aligned}
$$

but also

$$
\lim _{\epsilon \rightarrow+\infty} d_{\epsilon}(f, g)=\sup _{x \in \mathbb{R}}|f(x)-g(x)|=\|f-g\|_{\infty}
$$

here, as usual, $\|h\|_{\infty}=\operatorname{ess} \sup |h|=\inf \left\{y \in \mathbb{R}^{+}: \lambda(\{|h| \geq y\})=0\right\}$ for every measurable function $h: \mathbb{R} \rightarrow \overline{\mathbb{R}}$.

Given $f \in \mathcal{F}$, let $I \subset \overline{\mathbb{R}}$ be any interval with the property that

$$
\begin{equation*}
f_{-}(\sup I-x),-f(\inf I+x)<+\infty \quad \text { for some } x \in \mathbb{R} \tag{2.1}
\end{equation*}
$$

and consider the auxiliary function $\ell_{f, I}: \mathbb{R} \rightarrow \mathbb{R}$, introduced in [4], with

$$
\ell_{f, I}(x)=\inf \left\{y \in \mathbb{R}^{+}: f_{-}(\sup I-y)-y \leq x \leq f(\inf I+y)+y\right\} \quad \forall x \in \mathbb{R}
$$

also, let $\ell_{f, I}^{*}=\inf \left\{y \in \mathbb{R}^{+}: f_{-}(\sup I-y)-y \leq f(\inf I+y)+y\right\}$. For any sequence $\left(I_{k}\right)_{k \in \mathbb{N}}$ of intervals in $\overline{\mathbb{R}}$, write $\lim _{k \rightarrow \infty} I_{k}=I$ if $\lim _{k \rightarrow \infty} \inf I_{k}=\inf I$ and $\lim _{k \rightarrow \infty} \sup I_{k}=\sup I$.

Proposition 2.3 Let $f \in \mathcal{F}$, and let $I \subset \overline{\mathbb{R}}$ be an interval satisfying (2.1).
(i) The function $\ell_{f, I}$ is Lipschitz continuous and nonnegative;
(ii) $\ell_{f, I}(x) \geq \ell_{f, I}^{*} \geq 0$ for all $x \in \mathbb{R}$;
(iii) $\ell_{f, I}^{*} \leq \frac{1}{2} \lambda(I)$, and $\ell_{f, I}^{*}=0$ if and only if $f_{-}(\sup I) \leq f(\inf I)$;
(iv) If $\left(I_{k}\right)_{k \in \mathbb{N}}$ is a sequence of intervals in $\overline{\mathbb{R}}$ with $\lim _{k \rightarrow \infty} I_{k}=I$, then $I_{k}$ satisfies (2.1) for all sufficiently large $k$, and

$$
\lim _{k \rightarrow \infty} \ell_{f, I_{k}}^{*}=\ell_{f, I}^{*} \text { as well as } \lim _{k \rightarrow \infty} \ell_{f, I_{k}}(x)=\ell_{f, I}(x) \quad \forall x \in \mathbb{R}
$$

Remark 1 (i) If $f=F_{\mu}$ (respectively, $f=F_{\mu}^{-1}$ ) for some $\mu \in \mathcal{P}$, then every (respectively, every bounded) interval $I \subset \overline{\mathbb{R}}$ satisfies (2.1). Given $f \in \mathcal{F}$, note that $f=F_{\mu}$ for some (necessarily unique) $\mu \in \mathcal{P}$ if and only if $f(-\infty)=0$ and $f(+\infty)=1$; similarly, $f=F_{\mu}^{-1}$ for some $\mu \in \mathcal{P}$ if and only if $f_{-}(0)=-\infty$, $f(1)=+\infty$, and $f(] 0,1[) \subset \mathbb{R}$.
(ii) The function $\ell_{f, I}$ may not attain a minimum value, or when it does, that minimum value may be larger than $\ell_{f, I}^{*}$. However, mild additional assumptions guarantee that $\ell_{f, I}(x)=\ell_{f, I}^{*}$ for some $x \in \mathbb{R}$; see [4, Prop. 3.3].

For every $\epsilon>0$, consider the $\epsilon$-Lévy metric on $\mathcal{P}$ given by

$$
d_{\epsilon}(\mu, \nu)=d_{\epsilon}\left(F_{\mu}, F_{\nu}\right) \quad \forall \mu, v \in \mathcal{P}
$$

The metric $d_{\epsilon}$ is complete, separable and induces the topology of weak convergence (for an authoritative account on the family $\left(d_{\epsilon}\right)_{\epsilon>0}$, the reader may want to consult [31, Sec. 4.2]; see also [34]). Note that $\epsilon \mapsto d_{\epsilon}(\mu, v)$ is non-decreasing with $\lim _{\epsilon \rightarrow 0} d_{\epsilon}(\mu, \nu)=0$, whereas

$$
\lim _{\epsilon \rightarrow+\infty} d_{\epsilon}(\mu, \nu)=\left\|F_{\mu}-F_{\nu}\right\|_{\infty} \quad \forall \mu, \nu \in \mathcal{P}
$$

often referred to as the uniform or Kolmogorov metric, yields a complete yet nonseparable metric on $\mathcal{P}$ and induces a finer topology [4, Sec. 5]. For any $\mu \in \mathcal{P}$ and $\epsilon>0$, and with the dilation $T_{\epsilon}: x \mapsto \epsilon x$, notice the simple but useful identity

$$
\begin{equation*}
d_{\epsilon}(\mu, v)=d_{1}\left(\mu \circ T_{\epsilon}^{-1}, v \circ T_{\epsilon}^{-1}\right) \quad \forall \mu, v \in \mathcal{P} . \tag{2.2}
\end{equation*}
$$

To study finitely supported (and hence purely atomic) $d_{\epsilon}$-approximations of any $\mu \in$ $\mathcal{P}$, this article employs the following notations: for every $n \in \mathbb{N}$, let $\Xi_{n}=\left\{x \in \mathbb{R}^{n}\right.$ : $\left.x_{, 1} \leq \ldots \leq x_{, n}\right\}, \Pi_{n}=\left\{p \in \mathbb{R}^{n}: p_{, j} \geq 0, \sum_{j=1}^{n} p_{, j}=1\right\}$, and for each $x \in \Xi_{n}$ and $p \in \Pi_{n}$ let $\delta_{x}^{p}=\sum_{j=1}^{n} p_{, j} \delta_{x_{, j}}$. For convenience, $x_{, 0}:=-\infty$ and $x_{, n+1}:=+\infty$ for every $x \in \Xi_{n}$, as well as $P_{, i}:=\sum_{j=1}^{i} p_{, j}$ for $i=0, \ldots, n$ and every $p \in \Pi_{n}$; note that $P_{, 0}=0$ and $P_{, n}=1$. Henceforth, usage of the symbol $\delta_{x}^{p}$ tacitly assumes that $x \in \Xi_{n}$ and $p \in \Pi_{n}$, for some $n \in \mathbb{N}$ either specified explicitly or else clear from the context. Utilizing (2.2) and [4, Lem. 3.4], the value of $d_{\epsilon}\left(\mu, \delta_{x}^{p}\right)$ allows for simple explicit expressions.

Proposition 2.4 Let $\mu \in \mathcal{P}, \epsilon>0$, and $n \in \mathbb{N}$. For every $x \in \Xi_{n}$ and $p \in \Pi_{n}$,

$$
\begin{equation*}
d_{\epsilon}\left(\mu, \delta_{x}^{p}\right)=\epsilon \max _{j=0}^{n} \ell_{F_{\mu} / \epsilon,\left[x_{, j}, x_{, j+1}\right]}\left(P_{, j} / \epsilon\right)=\max _{j=1}^{n} \ell_{\epsilon F_{\mu}^{-1},\left[P_{, j-1}, P_{, j}\right]}\left(\epsilon x_{, j}\right) \tag{2.3}
\end{equation*}
$$

For every $\mu \in \mathcal{P}, \epsilon>0$ and $n \in \mathbb{N}$, (2.3) suggests considering the following quantities: given $x \in \Xi_{n}$, let

$$
\begin{gathered}
\ell_{x}^{\bullet}=\epsilon \max \left\{\ell_{F_{\mu} / \epsilon,[-\infty, x, 1]}(0), \ell_{F_{\mu} / \epsilon,\left[x_{1}, x_{, 2}\right]}^{*}, \ldots, \ell_{F_{\mu} / \epsilon,\left[x_{, j-1}, x_{, j}\right]}^{*}, \ldots\right. \\
\left.\ldots, \ell_{F_{\mu} / \epsilon,\left[x, n-1, x_{n}\right]}^{*}, \ell_{F_{\mu} / \epsilon,[x, n,+\infty]}(1 / \epsilon)\right\}
\end{gathered}
$$

and given $p \in \Pi_{n}$, let

$$
\ell_{\bullet}^{p}=\max _{j=1}^{n} \ell_{\epsilon F_{\mu}^{-1},\left[P_{, j-1}, P_{, j}\right]}^{*}
$$

Notice that while $\ell_{x}^{\bullet}$ and $\ell_{\bullet}^{p}$ do depend on $\mu, \epsilon$, and, implicitly, also $n$, in order to keep notations simple, this dependence is not displayed explicitly. By Proposition 2.3(iv), $p \mapsto \ell_{\bullet}^{p}$ is continuous on $\Pi_{n}$, and hence

$$
\ell_{\bullet}^{\bullet, n}=\min _{p \in \Pi_{n}} \ell_{\bullet}^{p}
$$

is well defined (for a constructive alternative definition of $\ell_{\bullet}^{\bullet, n}$, see [4, Sec. 3]). The quantities $\ell_{x}^{\bullet}, \ell_{\bullet}^{p}$ and $\ell_{\bullet}^{\bullet, n}$ control the minimization of $(x, p) \mapsto d_{\epsilon}\left(\mu, \delta_{x}^{p}\right)$, with or without constraints, in a sense made precise by Proposition 2.5. To formulate the result, call $\delta_{x}^{p}$ a best $d_{\epsilon}$-approximation of $\mu \in \mathcal{P}$, given $x \in \Xi_{n}$ if

$$
d_{\epsilon}\left(\mu, \delta_{x}^{p}\right) \leq d_{\epsilon}\left(\mu, \delta_{x}^{q}\right) \quad \forall q \in \Pi_{n} .
$$

Similarly, call $\delta_{x}^{p}$ a best $d_{\epsilon}$-approximation of $\mu$, given $p \in \Pi_{n}$ if

$$
d_{\epsilon}\left(\mu, \delta_{x}^{p}\right) \leq d_{\epsilon}\left(\mu, \delta_{y}^{p}\right) \quad \forall y \in \Xi_{n} .
$$

Denote by $\delta_{x}^{\bullet}$ and $\delta_{\bullet}^{p}$ any best $d_{\epsilon}$-approximation of $\mu$, given $x$ and $p$, respectively. Best $d_{\epsilon}$-approximations, given $p=u_{n}:=\left(n^{-1}, \ldots, n^{-1}\right)$, are referred to as best uniform $d_{\epsilon}$-approximations, and denoted $\delta_{\bullet}^{u_{n}}$. Finally, $\delta_{x}^{p}$ is a best $d_{\epsilon}$-approximation of $\mu \in \mathcal{P}$, denoted $\delta_{\bullet}^{\bullet, n}$, if

$$
d_{\epsilon}\left(\mu, \delta_{x}^{p}\right) \leq d_{\epsilon}\left(\mu, \delta_{y}^{q}\right) \quad \forall y \in \Xi_{n}, q \in \Pi_{n} .
$$

Notice that usage of the symbols $\delta_{x}^{\bullet}, \delta_{\bullet}^{p}$ and $\delta_{\bullet}^{\bullet}, n$ always refers to specific $\mu \in \mathcal{P}$, $\epsilon>0$ and $n \in \mathbb{N}$, all of which are usually clear from the context.

Proposition 2.5 Let $\mu \in \mathcal{P}, \epsilon>0$ and $n \in \mathbb{N}$.
(i) For every $x \in \Xi_{n}$, there exists a best $d_{\epsilon}$-approximation of $\mu$, given $x$, and $d_{\epsilon}\left(\mu, \delta_{x}^{\bullet}\right)=\ell_{x}^{\bullet}$. Moreover, $d_{\epsilon}\left(\mu, \delta_{x}^{p}\right)=\ell_{x}^{\bullet}$ with $p \in \Pi_{n}$ if and only if

$$
\begin{equation*}
\epsilon \ell_{F_{\mu} / \epsilon,\left[x_{, j}, x_{j+1}\right]}\left(P_{, j} / \epsilon\right) \leq \ell_{x}^{\bullet} \quad \forall j=0, \ldots, n \tag{2.4}
\end{equation*}
$$

(ii) For every $p \in \Pi_{n}$, there exists a best $d_{\epsilon}$-approximation of $\mu$, given $p$ and $d_{\epsilon}\left(\mu, \delta_{\bullet}^{p}\right)=\ell_{\bullet}^{p}$. Moreover, $d_{\epsilon}\left(\mu, \delta_{x}^{p}\right)=\ell_{\bullet}^{p}$ with $x \in \Xi_{n}$ if and only if

$$
\begin{equation*}
\left.\ell_{\epsilon F_{\mu}^{-1},[P, j-1}, P, j\right], \tag{2.5}
\end{equation*}
$$

(iii) There exists a best $d_{\epsilon}$-approximation of $\mu$, and $d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=\ell_{\bullet}^{\bullet, n}$. Moreover, $d_{\epsilon}\left(\mu, \delta_{x}^{p}\right)=\ell_{\bullet}^{\bullet, n}$ with $x \in \Xi_{n}, p \in \Pi_{n}$ if and only if (2.4) and (2.5) hold with $\ell_{\bullet}^{\bullet, n}$ instead of $\ell_{x}^{\bullet}$ and $\ell_{\bullet}^{p}$, respectively.

The following two examples illustrate Proposition 2.5. Notice that in either example the sequences $\left(d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)\right)$ and $\left(d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)\right)$ both converge to 0 at the same rate $\left(n^{-1}\right)$. As demonstrated in Sects. 3 and 4 for best uniform and best $d_{\epsilon}$-approximations, respectively, this rate is not specific to these examples, but rather indicative of much more general mechanisms.
Example 1 Consider the exponential distribution $\exp (a)$ with parameter $a>0$, i.e. let $F_{\mu}(x)=1-e^{-a x}$ for all $x \geq 0$. From Proposition 2.5, it is easily deduced that $\delta_{x}^{u_{n}}$ with $x \in \Xi_{n}$ is a best uniform $d_{\epsilon}$-approximation of $\mu$ if and only if, for every $j=1, \ldots, n$,

$$
x_{, j} \in\left[-\frac{1}{a} \log \left(1-\frac{j}{n}+\ell_{\bullet}^{u_{n}}\right)-\frac{\ell_{\bullet}^{u_{n}}}{\epsilon},-\frac{1}{a} \log \left(1-\frac{j-1}{n}-\ell_{\bullet}^{u_{n}}\right)+\frac{\ell_{\bullet}^{u_{n}}}{\epsilon}\right],
$$

with $\ell_{\bullet}^{u_{n}}=d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)$ being the unique solution of $n \ell\left(e^{2 a \ell / \epsilon}+1\right)=1$. A straightforward analysis of the latter equation yields the asymptotic equality

$$
\begin{equation*}
n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\frac{1}{2}-\frac{a}{4 \epsilon} n^{-1}+\mathcal{O}\left(n^{-2}\right) \quad \text { as } n \rightarrow \infty . \tag{2.6}
\end{equation*}
$$

A best $d_{\epsilon}$-approximation of $\mu$ also exists and in fact is unique, with

$$
x_{, j}=-\frac{1}{a} \log \left(\frac{e^{2 a \ell_{\bullet}^{\bullet n}(n-j) / \epsilon}-1}{e^{2 a \ell_{:}^{\bullet, n} n / \epsilon}-1}+\ell_{\bullet}^{\bullet, n}\right)-\frac{\ell_{\bullet}^{\bullet, n}}{\epsilon}, \quad P_{, j}=\frac{1-e^{-2 a \ell_{\bullet}^{, n} j / \epsilon}}{1-e^{-2 a \ell \ell_{\bullet}^{n} n / \epsilon}},
$$

for all $j=1, \ldots, n$, where $\ell_{\bullet}^{\bullet, n}=d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)$ solves $\ell e^{2 n a \ell / \epsilon}=\ell+\tanh (a \ell / \epsilon)$. Similarly to before, an analysis of this equation yields

$$
\begin{equation*}
n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=c_{1}-\frac{a^{2} c_{1}^{2}}{6 \epsilon(a+\epsilon)} n^{-2}+\mathcal{O}\left(n^{-4}\right) \quad \text { as } n \rightarrow \infty \tag{2.7}
\end{equation*}
$$

with $c_{1}=\frac{1}{2} \epsilon \log (1+a / \epsilon) / a<\frac{1}{2}$. Notice that $d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet}\right)<d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)$ for every $n \geq 2$; see also Fig. 1 .


Fig. 1 Illustrating $d_{1}$-approximations of the standard exponential distribution (dotted curve) with $n=4$ atoms: while the best approximation $\delta_{\bullet}^{\bullet}, 4$ (solid line) is unique, best uniform approximations $\delta_{\bullet}^{u_{4}}$ (broken lines) are not; see Example 1

Example 2 Fix $b>1$, and let $F_{\mu}(x)=\frac{\log x}{\log b}$ for all $1 \leq x \leq b$. Usually referred to as Benford's law base $b$, this distribution has many interesting properties; see, e.g. [2,4] and the references therein. As in the previous example, best uniform $d_{\epsilon}$-approximations of $\mu$ are non-unique, yet $\ell_{\bullet}^{u_{n}}$ is the unique solution of $b^{1-\ell}-b^{1+\ell-1 / n}=2 \ell / \epsilon$, which in turn yields

$$
\begin{equation*}
n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=c_{2}-\frac{c_{2}^{2}}{b \epsilon} n^{-1}+\mathcal{O}\left(n^{-2}\right) \quad \text { as } n \rightarrow \infty \tag{2.8}
\end{equation*}
$$

with $c_{2}=\frac{1}{2} \epsilon b \log b /(1+\epsilon b \log b)<\frac{1}{2}$. Also similarly to Example 1, best $d_{\epsilon}$-approximations of $\mu$ are unique, $\ell_{\bullet}^{\bullet}, n$ solves $b^{2 n \ell}(\ell+\epsilon \sinh (\ell \log b))=\ell+$ $\epsilon b \sinh (\ell \log b)$, and a straightforward analysis yields

$$
\begin{equation*}
n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=c_{3}+\frac{(b-1) c_{2}^{2} c_{3}^{2} b^{2 c_{3}-2}}{3 \epsilon} n^{-2}+\mathcal{O}\left(n^{-4}\right) \quad \text { as } n \rightarrow \infty \tag{2.9}
\end{equation*}
$$

with $c_{3}=\frac{\log (1+\epsilon b \log b)-\log (1+\epsilon \log b)}{2 \log b}<c_{2}$; notice again $d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)<$ $d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)$ for every $n \geq 2$.

Remark 2 The reader may find it instructive to compare the quantitative results in Examples 1 and 2 to their $d_{\mathrm{W}}$ - and $d_{\mathrm{P}}$-counterparts. On the one hand, for $\mu=\exp (a)$
[37, Ex. 5.8] shows that

$$
n d_{\mathrm{W}}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\frac{1}{4 a} \log n+\mathcal{O}(1) \quad \text { as } n \rightarrow \infty
$$

whereas by [18, Ex. 5.7],

$$
n d_{\mathrm{W}}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=\frac{n}{a} \log \frac{n+1}{n}=\frac{1}{a}+\mathcal{O}\left(n^{-1}\right) \quad \text { as } n \rightarrow \infty .
$$

While the authors are not aware of any result regarding best uniform $d_{\mathrm{p}}$-approximations, it follows from [20, Prop. 2.1] that

$$
n d_{\mathrm{P}}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=\frac{1}{2 a} \log n-\frac{1}{2 a} \log \log n+\mathcal{O}(1) \quad \text { as } n \rightarrow \infty
$$

On the other hand, for $\mu$ being Benford's law base $b>1$,

$$
\begin{aligned}
& n d_{\mathrm{W}}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\frac{n}{\log b} \tanh \left(\frac{\log b}{4 n}\right)=\frac{1}{4}+\mathcal{O}\left(n^{-2}\right) \quad \text { as } n \rightarrow \infty, \\
& n d_{\mathrm{W}}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=\frac{1}{\log b} \tanh \left(\frac{\log b}{4}\right) \quad \forall n \in \mathbb{N},
\end{aligned}
$$

as recorded in [4], and

$$
n d_{\mathrm{p}}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=\frac{b-1}{2}+\mathcal{O}\left(n^{-1}\right) \quad \text { as } n \rightarrow \infty
$$

again by [20, Prop. 2.1]. When compared to Lévy approximations, therefore, the asymptotics of best (or best uniform) $d_{\mathrm{W}}$ - and $d_{\mathrm{P}}$-approximations of $\mu$ are more sensitive to the tail behaviour of $\mu$. Though perhaps not obvious from (1.1)-(1.3), this is not too surprising either: informally, finiteness of the integral in (1.1) alone necessitates moderate tails for $\mu, \nu$, and taking the infimum over all Borel sets in (1.2) makes $d_{\mathrm{P}}$ pick up small deviations near infinity. By contrast, the infimum in (1.3) is, in the parlance of (1.2), taken over unbounded intervals only, which in turn causes $d_{\epsilon}$ to be much less sensitive to heavy or irregular tails, an informal observation that is corroborated rigorously by the results in the remaining sections of this article.

## 3 Best Uniform Lévy Approximations

This section provides a detailed asymptotic analysis of $d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)$ for any $\mu \in$ $\mathcal{P}$. Notice the uniform bound $n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right) \leq \frac{1}{2}$, due to Proposition 2.3. Thus, $d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right) \rightarrow 0$ as $n \rightarrow \infty$ at an (upper) rate not slower than $\left(n^{-1}\right)$. Except for trivial cases, this rate is sharp.

Lemma 3.1 Let $\mu \in \mathcal{P}$ and $\epsilon>0$. Then, $\lim \sup _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)>0$ unless $\mu=\delta_{a}$ for some $a \in \mathbb{R}$.

Proof Throughout the proofs of this section, write $g=F_{\mu}^{-1}$ for convenience, and let $\omega_{n}=n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)$ for all $n \in \mathbb{N}$, as well as $\omega^{-}=\liminf _{n \rightarrow \infty} \omega_{n}$ and $\omega^{+}=$ $\lim \sup _{n \rightarrow \infty} \omega_{n}$. Since $\delta_{a} \circ T_{\epsilon}^{-1}=\delta_{\epsilon a}$, by (2.2), it suffices to consider the case of $\epsilon=1$. Fix any $0<x<y<1$. Assume that $\omega^{+}=0$, i.e. $\lim _{n \rightarrow \infty} \omega_{n}=0$. Note that, for every $n \in \mathbb{N}$,

$$
\begin{equation*}
g_{-}\left(\frac{j-2 \omega_{n}}{n}\right)-\frac{2 \omega_{n}}{n} \leq g\left(\frac{j-1+2 \omega_{n}}{n}\right)+\frac{2 \omega_{n}}{n} \quad \forall j=1, \ldots, n, \tag{3.1}
\end{equation*}
$$

by the definition of $\omega_{n}$. Also, observe that for all sufficiently large $n$,

$$
\frac{j+2 \omega_{n+1}}{n+1}<\frac{j-2 \omega_{n}}{n} \text { and } \frac{j+2 \omega_{n}}{n}<\frac{j+1-2 \omega_{n+1}}{n+1}
$$

for every $j=\lfloor n x\rfloor, \ldots,\lfloor n y\rfloor+1$, which, together with (3.1), yields

$$
\begin{aligned}
g(y)-g(x) \leq & g\left(\frac{\lfloor n y\rfloor+1+2 \omega_{n}}{n}\right)-g_{-}\left(\frac{\lfloor n x\rfloor-2 \omega_{n}}{n}\right) \\
= & \sum_{j=\lfloor n x\rfloor}^{\lfloor n y\rfloor+1}\left(g\left(\frac{j+2 \omega_{n}}{n}\right)-g_{-}\left(\frac{j-2 \omega_{n}}{n}\right)\right) \\
& +\sum_{j=\lfloor n x\rfloor+1}^{\lfloor n y\rfloor+1}\left(g_{-}\left(\frac{j-2 \omega_{n}}{n}\right)-g\left(\frac{j-1+2 \omega_{n}}{n}\right)\right) \\
\leq & \sum_{j=\lfloor n x\rfloor}^{\lfloor n y\rfloor+1}\left(g_{-}\left(\frac{j+1-2 \omega_{n+1}}{n+1}\right)-g\left(\frac{j+2 \omega_{n+1}}{n+1}\right)\right) \\
& +\frac{4 \omega_{n}}{n}(\lfloor n y\rfloor-\lfloor n x\rfloor+1) \\
\leq & \frac{4 \omega_{n+1}}{n+1}(\lfloor n y\rfloor-\lfloor n x\rfloor+2)+\frac{4 \omega_{n}}{n}(\lfloor n y\rfloor-\lfloor n x\rfloor+1) \\
\leq & 12 \omega_{n+1}+8 \omega_{n} .
\end{aligned}
$$

Since $\lim _{n \rightarrow \infty} \omega_{n}=0$ by assumption, and $0<x<y<1$ have been arbitrary, $g(0)=g_{-}(1)=a$ for some $a \in \mathbb{R}$, that is, $\mu=\delta_{a}$.

For the subsequent finer analysis, the following terminology is useful: for every $f \in \mathcal{F}$, let $G_{f}$ be the growth set of $f$, i.e. let

$$
G_{f}=\{x \in \mathbb{R}: f(x-\varepsilon)<f(x+\varepsilon) \forall \varepsilon>0\} .
$$

Note that $G_{f}$ is closed in $\mathbb{R}$, and $G_{f} \neq \varnothing$, unless $f$ is constant. For example, $G_{F_{\mu}}=\operatorname{supp} \mu$ and $\{0,1\} \subset G_{F_{\mu}^{-1}} \subset[0,1]$ for every $\mu \in \mathcal{P}$. Also, $f(x) \in \mathbb{R}$ whenever $f^{-1}(-\infty)<x<f^{-1}(+\infty)$. With $\left.\left.\lambda_{f}(]-\infty, f^{-1}(-\infty)\right]\right):=$ $\lambda_{f}\left(\left[f^{-1}(+\infty),+\infty[):=0\right.\right.$ and

$$
\left.\left.\lambda_{f}(] x, y\right]\right):=f(y)-f(x) \quad \forall f^{-1}(-\infty)<x \leq y<f^{-1}(+\infty)
$$

therefore, $\lambda_{f}$ is a $\sigma$-finite positive Borel measure concentrated on $G_{f}$. For example, $\lambda_{\mathrm{id}_{\mathbb{R}}}=\lambda$, and $\lambda_{F_{\mu}}=\mu$ for every $\mu \in \mathcal{P}$. Also, $\mu^{-1}:=\lambda_{F_{\mu}^{-1}}$ is a positive Borel measure supported on $G_{F_{\mu}^{-1}} \subset[0,1]$, referred to as the inverse measure of $\mu$; see, e.g. [6,37]. For convenience, write $G_{F_{\mu}}$ and $G_{F_{\mu}^{-1}}$ simply as $G_{\mu}$ and $G_{\mu^{-1}}$, respectively. Note that $\mu^{-1}(\mathbb{R})=\mu^{-1}(] 0,1[)=\operatorname{diam} G_{\mu}$, and hence $\mu^{-1}=0$ precisely if $\mu=\delta_{a}$ for some $a \in \mathbb{R}$. When rephrased utilizing this terminology, Lemma 3.1 has the following corollary.

Proposition 3.2 For every $\mu \in \mathcal{P}$ and $\epsilon>0$, the following are equivalent:
(i) $\lim _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=0$;
(ii) $d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=0$ for every $n \in \mathbb{N}$;
(iii) $\mu=\delta_{a}$ for some $a \in \mathbb{R}$;
(iv) $\mu^{-1}=0$.

The first main result in this section asserts that $\left(n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)\right)$ does converge, to an easily determined limit, if $\mu^{-1}$ is absolutely continuous. The result is reminiscent of a theorem regarding best uniform $d_{\mathrm{W}}$-approximations [37, Thm. 5.15] (see also [7,16]), but unlike in that theorem, no integrability assumption on $\mathrm{d} \mu^{-1} / \mathrm{d} \lambda$ is needed, and the limit in question always is finite. When formulating the result, it is helpful to use the function $\Omega: \mathbb{R} \rightarrow \mathbb{R}$ with

$$
\Omega(x)=\frac{x}{2+2|x|} \quad \forall x \in \mathbb{R}
$$

Plainly, $\Omega$ is an increasing $C^{1}$-function, with $|\Omega(x)| \leq \frac{1}{2}|x|$ for all $x \in \mathbb{R}$, and $\Omega( \pm \infty)= \pm \frac{1}{2}$. While the appearance of $\Omega$ in the following theorem is a simple consequence of the bound (3.3), the reader may find it curious to note that $2 \Omega$ plays a prominent role in the theory of random walks [12].

Theorem 3.3 Let $\mu \in \mathcal{P}$ and $\epsilon>0$. If $\mu^{-1}$ is absolutely continuous (w.r.t. $\lambda$ ), then

$$
\begin{equation*}
\lim _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\left\|\Omega\left(\epsilon \frac{\mathrm{d} \mu^{-1}}{\mathrm{~d} \lambda}\right)\right\|_{\infty} \tag{3.2}
\end{equation*}
$$

Proof Since $\frac{\mathrm{d}\left(\mu \circ T_{\epsilon}^{-1}\right)^{-1}}{\mathrm{~d} \lambda}=\epsilon \frac{\mathrm{d} \mu^{-1}}{\mathrm{~d} \lambda}$, it is enough to prove (3.2) for $\epsilon=1$. Using the same symbols as in the proof of Lemma 3.1, for every $n \in \mathbb{N}$ let

$$
J_{n, j}=\left[\frac{j-1+\omega_{n}}{n}, \frac{j-\omega_{n}}{n}\right] \quad \forall j=1, \ldots, n .
$$

Note that $\omega_{n}<\frac{1}{2}$ for every $n$ since $g$ is continuous. Moreover,

$$
g\left(\frac{j-\omega_{n}}{n}\right)-g\left(\frac{j-1+\omega_{n}}{n}\right) \leq \frac{2 \omega_{n}}{n} \quad \forall j=1, \ldots, n,
$$

and consequently, by the absolute continuity of $g$,

$$
\begin{align*}
\frac{2 \omega_{n}}{1-2 \omega_{n}} & \geq \frac{g\left(\left(j-\omega_{n}\right) / n\right)-g\left(\left(j-1+\omega_{n}\right) / n\right)}{\left(1-2 \omega_{n}\right) / n} \\
& =\frac{1}{\lambda\left(J_{n, j}\right)} \int_{J_{n, j}} g^{\prime} \mathrm{d} \lambda \quad \forall j=1, \ldots, n \tag{3.3}
\end{align*}
$$

Equality holds on the left in (3.3) for at least one $j$, and for that $j$,

$$
\frac{2 \omega_{n}}{1-2 \omega_{n}}=\frac{1}{\lambda\left(J_{n, j}\right)} \int_{J_{n, j}} g^{\prime} \mathrm{d} \lambda \leq\left\|g^{\prime}\right\|_{\infty}
$$

from which it is clear that

$$
\begin{equation*}
\omega^{+} \leq \frac{\left\|g^{\prime}\right\|_{\infty}}{2+2\left\|g^{\prime}\right\|_{\infty}}=\left\|\Omega\left(g^{\prime}\right)\right\|_{\infty} \leq \frac{1}{2} \tag{3.4}
\end{equation*}
$$

Since (3.2) trivially holds when $\omega^{-}=\frac{1}{2}$, henceforth assume $\omega^{-}<\frac{1}{2}$, and pick $n_{1}<n_{2}<\ldots$ so that $\lim _{k \rightarrow \infty} \omega_{n_{k}}=\omega^{-}$. Given any $0<x<1$, let $j_{k}(x)=$ $\left\lfloor n_{k} x\right\rfloor+1 \in\left\{1, \ldots, n_{k}\right\}$, and note that $J_{n_{k}, j_{k}(x)} \subset\left[x-1 / n_{k}, x+1 / n_{k}\right]$, but also

$$
\frac{\lambda\left(J_{n_{k}, j_{k}(x)}\right)}{\lambda\left(\left[x-1 / n_{k}, x+1 / n_{k}\right]\right)}=\frac{1}{2}-\omega_{n_{k}} \xrightarrow{k \rightarrow \infty} \frac{1}{2}-\omega^{-}>0 .
$$

Thus, the sequence $\left(J_{n_{k}}, j_{k}(x)\right)_{k \in \mathbb{N}}$ shrinks to $x$ nicely, and by [33, Thm. 7.10],

$$
g^{\prime}(x)=\lim _{k \rightarrow \infty} \frac{1}{\lambda\left(J_{n_{k}, j_{k}(x)}\right)} \int_{J_{n_{k}}, j_{k}(x)} g^{\prime} \mathrm{d} \lambda \leq \lim _{k \rightarrow \infty} \frac{2 \omega_{n_{k}}}{1-2 \omega_{n_{k}}}=\frac{2 \omega^{-}}{1-2 \omega^{-}}
$$

for $\lambda$-almost every $x \in[0,1]$. It follows that $\omega^{-} \geq \Omega\left(g^{\prime}\right)$ holds $\lambda$-almost everywhere, and hence $\omega^{-} \geq\left\|\Omega\left(g^{\prime}\right)\right\|_{\infty}$. Together with (3.4), this completes the proof.

If $\mu^{-1}$ is singular, then the asymptotic behaviour of $\left(d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)\right)$ is controlled by a different mechanism. To prepare for the general result, observe that $\lim \inf _{n \rightarrow \infty} \operatorname{dist}(n x, \mathbb{Z})=0$ for every $x \in \mathbb{R}$, whereas

$$
\lim \sup _{n \rightarrow \infty} \operatorname{dist}(n x, \mathbb{Z})= \begin{cases}\frac{1}{2}(q-1) / q & \text { if } x=p / q \text { with coprime }  \tag{3.5}\\ \frac{1}{2} & p \in \mathbb{Z}, q \in \mathbb{N}, q \text { odd } \\ \frac{\text { otherwise }}{}\end{cases}
$$

in particular, $\lim \sup _{n \rightarrow \infty} \operatorname{dist}(n x, \mathbb{Z}) \geq \frac{1}{3}$ unless $x \in \mathbb{Z}$. Defining $\iota: \mathbb{R} \rightarrow \mathbb{N}_{0} \cup\{+\infty\}$ as

$$
\iota(x)=2 \inf \left\{n \in \mathbb{N}_{0}:(2 n+1) x \in \mathbb{Z}\right\}
$$

notice that the right-hand side in (3.5) is nothing other than $\Omega \circ \iota(x)$. With this, consider the very simple example of $\mu_{a}=a \delta_{-1}+(1-a) \delta_{1}$ for some $0<a<1$, for which $G_{\mu_{a}^{-1}}=\{0, a, 1\}$. It is readily confirmed that $n d_{\epsilon}\left(\mu_{a}, \delta_{\bullet}^{u_{n}}\right)=\operatorname{dist}(n a, \mathbb{Z})$ for all sufficiently large $n$, and hence ${\lim \inf _{n \rightarrow \infty}} n d_{\epsilon}\left(\mu_{a}, \delta_{\bullet}^{u_{n}}\right)=0$, as well as

$$
\lim \sup _{n \rightarrow \infty} n d_{\epsilon}\left(\mu_{a}, \delta_{\bullet}^{u_{n}}\right)=\Omega \circ \iota(a)=\max \Omega \circ \iota\left(G_{\mu_{a}^{-1}}\right)
$$

This equality is but one manifestation of a general principle.
Lemma 3.4 Let $\mu \in \mathcal{P}$ and $\epsilon>0$. If $\mu^{-1}$ is singular (w.r.t. $\lambda$ ), then

$$
\begin{equation*}
\limsup _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\sup \Omega \circ \iota\left(G_{\mu^{-1}}\right) . \tag{3.6}
\end{equation*}
$$

Proof Using the same symbols as in previous proofs, write $G_{\mu^{-1}}$ simply as $G$, and let $2 m=\sup \iota(G)$ with the appropriate $m \in \mathbb{N}_{0} \cup\{+\infty\}$; also, let $G^{*}$ be the set of atoms of $\mu^{-1}$, i.e. $G^{*}=\left\{0<x<1: g_{-}(x)<g(x)\right\}$. Assume first that $m \in \mathbb{N}_{0}$. Since $m=0$ implies $G=\{0,1\}$, or equivalently $\mu^{-1}=0$, and (3.6) is correct in this case by Proposition 3.2, henceforth assume $m \geq 1$. Then, $\mu^{-1}$ is concentrated on finitely many atoms, thus

$$
G^{*}=G=\left\{0, \frac{k_{1}}{2 m_{1}+1}, \ldots, \frac{k_{l}}{2 m_{l}+1}, 1\right\},
$$

with the appropriate positive integers $l, k_{1}, \ldots, k_{l}, m_{1}, \ldots, m_{l}$, where the numbers $k_{i}, 2 m_{i}+1$ are coprime for all $i$, and $\max _{i=1}^{l} m_{i}=m$. As seen in the example above, for all sufficiently large $n$,

$$
\omega_{n}=\max _{i=1}^{l} \operatorname{dist}\left(\frac{n k_{i}}{2 m_{i}+1}, \mathbb{Z}\right),
$$

and hence

$$
\omega^{+}=\max _{i=1}^{l} \Omega \circ \iota\left(\frac{k_{i}}{2 m_{i}+1}\right)=\max _{i=1}^{l} \Omega\left(2 m_{i}\right)=\Omega(2 m),
$$

so again (3.6) is correct. It remains to consider the case of $m=+\infty$. Here, it is convenient to consider two subcases, depending on whether $\iota\left(G^{*}\right)$ is unbounded or not. In the former case, fix $a \in \mathbb{R}^{+}$, and pick $x \in G^{*}$ with $\iota(x) \geq a$. Moreover, pick $b>3$, and recall that $y_{n}:=(\operatorname{dist}(n x, \mathbb{Z})-1 / b) / n>0$ for infinitely many $n \in \mathbb{N}$. Since $x$ is an atom of $\mu^{-1}$, for every $c \in \mathbb{R}^{+}$clearly

$$
\begin{equation*}
g_{-}\left(x+\frac{1}{b n}\right)-g\left(x-\frac{1}{b n}\right) \geq \frac{c}{b n} \text { for all sufficiently large } n . \tag{3.7}
\end{equation*}
$$

Choosing $c=b / \epsilon$ in (3.7), note that for infinitely many $n$,

$$
\begin{aligned}
g_{-} & \left(\frac{\lfloor n x\rfloor+1-n y_{n}}{n}\right)-g\left(\frac{\lfloor n x\rfloor+n y_{n}}{n}\right) \\
& =g_{-}\left(x+\frac{\max \{1-2 \operatorname{dist}(n x, \mathbb{Z}), 0\}}{n}+\frac{1}{b n}\right) \\
& -g\left(x+\frac{\min \{1-2 \operatorname{dist}(n x, \mathbb{Z}), 0\}}{n}-\frac{1}{b n}\right) \\
& \geq g_{-}\left(x+\frac{1}{b n}\right)-g\left(x-\frac{1}{b n}\right) \geq \frac{1}{n \epsilon} \geq \frac{2 y_{n}}{\epsilon},
\end{aligned}
$$

and consequently $\omega_{n} \geq n y_{n}$. It follows that

$$
\omega^{+} \geq \lim \sup _{n \rightarrow \infty}\left(\operatorname{dist}(n x, \mathbb{Z})-\frac{1}{b}\right)=\Omega \circ \iota(x)-\frac{1}{b} \geq \Omega(a)-\frac{1}{b}
$$

Since $a, b>3$ have been arbitrary, $\omega^{+}=\frac{1}{2}=\Omega(2 m)$. Finally, assume that $\iota\left(G^{*}\right)$ is bounded, and hence $G^{*}$ is finite, possibly empty. Since $m=+\infty$, clearly $G \backslash G^{*} \neq \varnothing$, and every $x \in G \backslash G^{*}$ is a continuity point of $g$, as well as an accumulation point of $G$. By [33, Thm. 7.15],

$$
\lim _{\varepsilon \downarrow 0} \frac{g_{-}(x+\varepsilon)-g(x-\varepsilon)}{2 \varepsilon}=+\infty \text { for } \mu^{-1} \text {-almost every } 0<x<1
$$

From this, it is clear that, given any $b, c>3$, there exists $x \in G \backslash \mathbb{Q}$ for which (3.7) holds. With $\iota(x)=a=+\infty$, the same argument as before shows that $\omega^{+}=\frac{1}{2}$, i.e. (3.6) is correct in this case also.

Combining Theorem 3.3 and Lemma 3.4 yields a sharp (upper) rate for $\left(d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)\right)$, for arbitrary $\mu \in \mathcal{P}$. To formulate the result, recall that every $\sigma$-finite Borel measure $\rho$ on the real line can be written uniquely as $\rho=\rho_{\mathrm{A}}+\rho_{\mathrm{S}}$, where $\rho_{\mathrm{A}}$ and $\rho_{\mathrm{S}}$ are absolutely continuous and singular (w.r.t. $\lambda$ ), respectively.

Theorem 3.5 Let $\mu \in \mathcal{P}$ and $\epsilon>0$. Then,

$$
\begin{equation*}
\lim _{\sup _{n \rightarrow \infty}} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\max \left\{\left\|\Omega\left(\epsilon \frac{\mathrm{d}\left(\mu^{-1}\right)_{\mathrm{A}}}{\mathrm{~d} \lambda}\right)\right\|_{\infty}, \sup \Omega \circ \iota\left(G_{\left.\left(\mu^{-1}\right)_{\mathrm{s}}\right)}\right)\right\} \tag{3.8}
\end{equation*}
$$

Proof Since there is nothing to prove otherwise, assume that $\left(\mu^{-1}\right)_{\mathrm{A}} \neq 0$ and $\left(\mu^{-1}\right)_{\mathrm{S}} \neq 0$. In analogy to the proof of Lemma 3.4, let $g=F_{\mu}^{-1}=g_{\mathrm{A}}+g_{\mathrm{S}}$, with $g_{\mathrm{A}}, g_{\mathrm{S}} \in \mathcal{F}$ such that $\lambda_{g_{\mathrm{A}}}=\left(\mu^{-1}\right)_{\mathrm{A}}$ and $\lambda_{g_{\mathrm{S}}}=\left(\mu^{-1}\right)_{\mathrm{S}}$, as well as $2 m=\sup \iota\left(G_{\mathrm{S}}\right) \in \mathbb{N}_{0} \cup\{+\infty\}$, where $G_{\mathrm{S}}=G_{\left(\mu^{-1}\right)_{\mathrm{S}}}$ for convenience. Since $g_{-}(y)-g(x) \geq g_{S_{-}}(y)-g_{\mathrm{S}}(x)$ for all $0<x<y<1$, Lemma 3.4 yields $\omega^{+} \geq \sup \Omega \circ \iota\left(G_{\mathrm{S}}\right)=\Omega(2 m)$. Thus, (3.8) clearly is correct when $m=+\infty$, and only the case of $m \in \mathbb{N}$ remains to be considered (note that $m=0$ is impossible,
as it would imply $\left.\left(\mu^{-1}\right)_{\mathrm{S}}=0\right)$. In this case, $G_{\mathrm{S}}$ is finite, say, $G_{\mathrm{S}}=\left\{0, x_{1}, \ldots, x_{l}, 1\right\}$ with $l \in \mathbb{N}$ and $0<x_{1}<\ldots<x_{l}<1$. With $J_{n, j}$ as in the proof of Theorem 3.3, and letting $K_{n}=\left\{\left\lfloor n x_{i}\right\rfloor: i=1, \ldots, l\right\} \subset\{1, \ldots, n-1\}$ for $n \geq 1 / x_{1}$, observe that

$$
\begin{aligned}
g_{\mathrm{A}}\left(\frac{j-\omega_{n}}{n}\right)-g_{\mathrm{A}}\left(\frac{j-1+\omega_{n}}{n}\right) & =g_{-}\left(\frac{j-\omega_{n}}{n}\right)-g\left(\frac{j-1+\omega_{n}}{n}\right) \\
& \leq \frac{2 \omega_{n}}{n \epsilon} \quad \forall j \notin K_{n},
\end{aligned}
$$

and consequently

$$
\frac{1}{\lambda\left(J_{n, j}\right)} \int_{J_{n, j}} \epsilon g_{\mathrm{A}}^{\prime} \mathrm{d} \lambda \leq \frac{2 \omega_{n}}{1-2 \omega_{n}} \quad \forall j \notin K_{n} .
$$

If $\omega^{-}<\frac{1}{2}$, then the same argument as in the proof of Theorem 3.3 shows that

$$
\epsilon g_{\mathrm{A}}^{\prime}(x) \leq \frac{2 \omega^{-}}{1-2 \omega^{-}} \text {for } \lambda \text {-almost every } x
$$

since clearly $j_{n}(x) \notin K_{n}$ whenever $x \notin G_{S}$ and $n$ is sufficiently large. Thus, $\omega^{-} \geq\left\|\Omega\left(\epsilon g_{\mathrm{A}}^{\prime}\right)\right\|_{\infty} ;$ trivially, the latter also holds when $\omega^{-}=\frac{1}{2}$. In summary, $\omega^{+} \geq \max \left\{\left\|\Omega\left(\epsilon g_{\mathrm{A}}^{\prime}\right)\right\|_{\infty}, \Omega(2 m)\right\}=: \omega$; note that $\omega$ simply equals the right-hand side in (3.8).

The reverse inequality is non-trivial only when $\omega<\frac{1}{2}$. In this case, assume $m \in \mathbb{N}$ as before, and pick any $z$ with $\omega<z<\frac{1}{2}$. Then, for all sufficiently large $n$,

$$
g_{\mathrm{A}}\left(\frac{j-z}{n}\right)-g_{\mathrm{A}}\left(\frac{j-1+z}{n}\right) \leq \frac{2 z}{\epsilon n} \quad \forall j=1, \ldots, n,
$$

but also, since $G_{\mathrm{S}}$ is finite,

$$
g_{\mathrm{S}-}\left(\frac{j-z}{n}\right)-g_{\mathrm{S}}\left(\frac{j-1+z}{n}\right)=0 \quad \forall j=1, \ldots, n .
$$

Thus, $\omega_{n} \leq z$ for all sufficiently large $n$, and since $z>\omega$ was arbitrary, $\omega^{+} \leq \omega$.
Corollary 3.6 Let $\mu \in \mathcal{P}$ and $\epsilon>0$.
(i) If $\lim \sup _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)<\frac{1}{3}$, then $\left(n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)\right)$ converges, and $\mu^{-1}$ is absolutely continuous (w.r.t. $\lambda$ ).
(ii) If $\lim \sup _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right) \geq \frac{1}{3}$, then either $\left(n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)\right)$ converges, or
$\lim \sup _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right) \in\left\{\frac{1}{3}, \frac{2}{5}, \frac{3}{7}, \ldots, \frac{1}{2}\right\}=\{\Omega(2 m): m \in \mathbb{N} \cup\{+\infty\}\}$.

Remark 3 (i) The proof given above shows that, for every $\mu \in \mathcal{P}$ and $\epsilon>0$,

$$
\liminf _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right) \geq\left\|\Omega\left(\epsilon \frac{\mathrm{d}\left(\mu^{-1}\right)_{\mathrm{A}}}{\mathrm{~d} \lambda}\right)\right\|_{\infty}
$$

(ii) Let $\mu \in \mathcal{P}$ be non-atomic. Then, the right-hand side in (3.8) tends to $\frac{1}{2}$ as $\epsilon \rightarrow+\infty$. This is consistent with the fact that $n \min _{x \in \Xi_{n}}\left\|F_{\mu}-F_{\delta_{x_{n}}}^{u_{n}}\right\|_{\infty}=\frac{1}{2}$ for all $n \in \mathbb{N}$ whenever $\mu$ is non-atomic [4, Cor. 5.5].

The following example illustrates the results of this section. In particular, it demonstrates that all situations allowed by Theorem 3.5 and Corollary 3.6 do occur. It also shows that (3.2) may fail when $\mu^{-1}$ is not absolutely continuous; similarly, (3.6) may fail when $\mu^{-1}$ is not singular.
Example 3 For every $0 \leq a<1<b$, consider $\mu_{a, b}=a \delta_{-1}+(1-a) U_{1, b}$, where $U_{1, b}$ denotes the uniform distribution (normalized Lebesgue measure) on $[1, b]$. Note that $\mu_{a, 1}:=\lim _{b \downarrow 1} \mu_{a, b}=\mu_{a}$, with $\mu_{a}$ considered prior to Lemma 3.4. Since $\left(\mu_{a, b}^{-1}\right)_{\mathrm{A}}=$ $(b-1) U_{a, 1}$ and $\left(\mu_{a, b}^{-1}\right)_{S}=2 \delta_{a}$ provided that $a>0$, Theorem 3.5 yields

$$
\limsup _{n \rightarrow \infty} n d_{\epsilon}\left(\mu_{a, b}, \delta_{\bullet}^{u_{n}}\right)=\max \left\{\Omega\left(\epsilon \frac{b-1}{1-a}\right), \Omega \circ \iota(a)\right\},
$$

whereas by direct inspection,

$$
\liminf _{n \rightarrow \infty} n d_{\epsilon}\left(\mu_{a, b}, \delta_{\bullet}^{u_{n}}\right)=\Omega\left(\epsilon \frac{b-1}{1-a}\right)<\frac{1}{2}
$$

On the one hand, if $a=a_{m}=m /(2 m+1)$ for some $m \in \mathbb{N}_{0}$, then $\Omega \circ \iota\left(a_{m}\right)=a_{m}$, and since $b \mapsto \Omega\left(\epsilon(b-1) /\left(1-a_{m}\right)\right)$ is increasing continuously from 0 to $\frac{1}{2}$, there exists a unique $b_{m}$ with $\Omega\left(\epsilon\left(b_{m}-1\right) /\left(1-a_{m}\right)\right)=a_{m}$. Thus, $\left(n d_{\epsilon}\left(\mu_{a_{m}, b}, \delta_{\bullet}^{u_{n}}\right)\right)$ converges precisely if $b \geq b_{m}$, whereas for $b<b_{m}$ the lim inf can have any value between 0 and $a_{m}$. On the other hand, if $a=\frac{1}{2}$ then $\Omega \circ \iota(a)=\frac{1}{2}$, and again $\liminf _{n \rightarrow \infty} n d_{\epsilon}\left(\mu_{1 / 2, b}, \delta_{\bullet}^{u_{n}}\right)=\Omega(2 \epsilon(b-1))$ can have any value between 0 and $\frac{1}{2}$. Except for the case of $\lim _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\frac{1}{2}$, which occurs, e.g. for the exponential distributions in Example 1, every possible situation allowed by Theorem 3.5 can be observed in this example by choosing $a, b$ appropriately; see also Fig. 2. Notice that $\mu_{a, b}^{-1}$ is absolutely continuous precisely if $a=0$ and is singular only if $b=1$. While Theorem 3.3 and Lemma 3.4 thus cannot in general be reversed, clearly (3.2) and (3.6) may fail if $\mu^{-1}$ is not absolutely continuous and singular, respectively.

Example 4 Let $\mu$ be a normal distribution with variance $\sigma^{2}>0$. With $\phi$ denoting the standard normal distribution function, it is readily deduced from Proposition 2.5 that $\ell_{\bullet}^{u_{n}}=d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)$ is the unique solution of $\phi^{-1}(1-\ell)-\phi^{-1}(1+\ell-1 / n)=2 \ell / \sqrt{\epsilon^{2} \sigma^{2}}$. Utilizing the familiar fact [1, Sec. 26.2]

$$
1-\phi(x)=\frac{e^{-x^{2} / 2}}{\sqrt{2 \pi}}\left(x^{-1}-x^{-3}+\mathcal{O}\left(x^{-5}\right)\right) \quad \text { as } x \rightarrow+\infty
$$

Fig. 2 Solid black lines indicate, for any $\mu \in \mathcal{P}$, the possible values of the limit inferior and the limit superior of $\left(n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)\right)$; see
Theorems 3.3 and 3.5.
Example 3 demonstrates that all possible values may indeed occur

a straightforward analysis of this equation yields

$$
\begin{equation*}
n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\frac{1}{2}-\frac{1}{2 \epsilon \sqrt{2 \sigma^{2}}} \cdot \frac{\sqrt{\log n}}{n}+\mathcal{O}\left(n^{-1}\right) \quad \text { as } n \rightarrow \infty \tag{3.9}
\end{equation*}
$$

which sharpens (3.2). Notice that convergence occurs here at a slightly slower rate than has so far been observed in this article for examples of absolutely continuous $\mu^{-1}$; cf. [16] and [37, Ex. 5.18].

As Examples 1, 2 and 4 suggest, the results of this section, notably Theorem 3.3, can be refined by imposing further assumptions on $\mu$. For instance, assume that $g=F_{\mu}^{-1}$ is $C^{2}$ on $] 0,1\left[\right.$ and that $g, g^{\prime} \neq 0$ both are convex, with $\lim \sup _{x \uparrow 1}(1-$ $x) g^{\prime \prime}(x) / g^{\prime}(x)<+\infty$ (all mentioned examples meet these requirements; for normal distributions, only the right half of $] 0,1[$ has to be considered due to symmetry). Letting $c=\left\|\Omega\left(\epsilon g^{\prime}\right)\right\|_{\infty}>0$ for convenience, it is straightforward to show that, as a refinement of (3.2),

$$
\begin{equation*}
n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=c-\frac{2 c^{2}}{\epsilon} e_{n}+o\left(e_{n}\right) \quad \text { as } n \rightarrow \infty \tag{3.10}
\end{equation*}
$$

where $\lim _{n \rightarrow \infty} e_{n}=0$, and more specifically,

$$
e_{n}= \begin{cases}\frac{1}{g_{-}^{\prime}(1)}+n \frac{g(1-(1-c) / n)-g(1-c / n)}{(1-2 c) g_{-}^{\prime}(1)^{2}} & \text { if } c<\frac{1}{2} \\ \frac{1}{g^{\prime}(1-1 /(2 n))} & \text { if } c=\frac{1}{2}\end{cases}
$$

in particular, if $g_{-}^{\prime \prime}(1)<+\infty$ then simply

$$
e_{n}=\frac{g_{-}^{\prime \prime}(1)}{2 g_{-}^{\prime}(1)^{2}} n^{-1}
$$

As the reader may want to check, for exponential, Benford and normal distributions, the asymptotic equalities (2.6), (2.8) and (3.9), respectively, all are (slightly sharper than, but certainly) consistent with (3.10).

Example 5 Let $\mu$ be the Cantor distribution, i.e. the $\log 2 / \log 3$-dimensional Hausdorff measure on the classical Cantor middle thirds set $C$. Thus, $G_{\mu}=C$, and since $\operatorname{diam} C=1$, the measure $\mu^{-1}$, referred to as the inverse Cantor distribution [6, Ex. A.11], is a probability measure as well. Both $\mu, \mu^{-1}$ are singular: while $\mu$ is nonatomic, $\mu^{-1}$ is purely atomic; in fact, $\mu^{-1}\left(\left\{i 2^{-m}\right\}\right)=3^{-m}$ for every $m \in \mathbb{N}$ and every odd $1 \leq i<2^{m}$. Obviously, Lemma 3.4 applies to both distributions, showing that

$$
\lim \sup _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=\lim \sup _{n \rightarrow \infty} n d_{\epsilon}\left(\mu^{-1}, \delta_{\bullet}^{u_{n}}\right)=\frac{1}{2} .
$$

By contrast, it is straightforward to check that $\liminf _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=0$. For the inverse Cantor distribution, an elementary analysis [4,36] yields $\frac{1}{216} \leq$ $\lim _{\inf _{n \rightarrow \infty}} n d_{1}\left(\mu^{-1}, \delta_{\bullet}^{u_{n}}\right) \leq \frac{1}{3}$, but the authors do not know the precise value of $\lim \inf _{n \rightarrow \infty} n d_{\epsilon}\left(\mu^{-1}, \delta_{\bullet}^{u_{n}}\right)$ for any $\epsilon>0$.

## 4 Best (Unconstrained) Lévy Approximations

This section studies the asymptotics of $d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)$ as $n \rightarrow \infty$. The following theorem is the section's main result and a counterpart to Theorems 3.3 and 3.5. It asserts that the sequence $\left(n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)\right)$ always converges, to a limit smaller than $\frac{1}{2}$ that is easily expressed in terms of $\Omega$ and $\left(\mu^{-1}\right)_{\mathrm{A}}$, the absolutely continuous part of $\mu^{-1}$.

Theorem 4.1 Let $\mu \in \mathcal{P}$ and $\epsilon>0$. Then,

$$
\begin{equation*}
\lim _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=\int \Omega\left(\epsilon \frac{\mathrm{d}\left(\mu^{-1}\right)_{\mathrm{A}}}{\mathrm{~d} \lambda}\right) \mathrm{d} \lambda \tag{4.1}
\end{equation*}
$$

As pointed out already in Introduction, Theorem 4.1 may be regarded as an analogue for $d_{\epsilon}$ of (the one-dimensional version of) a classical $d_{\mathrm{W}}$-quantization theorem [18, Thm. 6.2], but unlike that theorem, it does not impose a moment assumption on $\left(\mu^{-1}\right)_{\mathrm{A}}$. For $d_{\mathrm{W}}$-approximations, such an assumption is known to be essential, as the $d_{\mathrm{W}}$-approximation error for $\mu \in \mathcal{P}_{1}$ can decay arbitrarily slowly without it; see [18, Ex. 6.4] and [37, Thm. 5.33]. In the light of Theorem 4.1, this may be viewed as an artefact specific to $d_{\mathrm{W}}$ that does not exist for $d_{\epsilon}$. Also, recall that unlike $d_{\mathrm{W}}$, the metric $d_{\epsilon}$ metrizes precisely the topology of weak convergence on all of $\mathcal{P}$, and so does $d_{\mathrm{p}}$. As far as the authors have been able to ascertain, however, all known results pertaining
to the asymptotics of the $d_{\mathrm{p}}$-approximation error also impose additional assumptions [20, Sec. 4], and despite the similarities between $d_{\mathrm{P}}$ and $d_{\epsilon},[20$, Sec. 5] suggests that the $d_{\mathrm{p}}$-approximation error can decay arbitrarily slowly as well.

When proving Theorem 4.1, the following observation, a direct consequence of Proposition 2.5 together with the argument establishing [4, Thm. 3.9], is helpful; its routine verification is left to the interested reader.

Proposition 4.2 Let $\mu \in \mathcal{P}, \epsilon>0$, and $n \in \mathbb{N}$. With $\ell=d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)$, there exists $p \in \Pi_{n}$ such that for every $j=1, \ldots, n$,

$$
\begin{aligned}
& \operatorname{dist}\left(P_{, j}, G_{\mu^{-1}}\right) \leq \ell \quad \text { and } \\
& \qquad \mu^{-1}(] P_{, j-1}+\ell, P_{, j}-\ell[) \leq \frac{2 \ell}{\epsilon} \leq \mu^{-1}\left(\left[P_{, j-1}+\ell, P_{, j}-\ell\right]\right)
\end{aligned}
$$

Proof of Theorem 4.1 Throughout this proof, for convenience let $g=F_{\mu}^{-1}$ and $G=$ $G_{\mu^{-1}}$ as before, but also $\ell_{n}=d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)$ and $\omega_{n}=n \ell_{n}$ for all $n \in \mathbb{N}$, as well as $\omega^{-}=\lim _{\inf }^{n \rightarrow \infty} \omega_{n}$ and $\omega^{+}=\lim \sup _{n \rightarrow \infty} \omega_{n}$. Again it suffices to consider the case of $\epsilon=1$. Note that $\ell_{n}=0$ for some (and hence all sufficiently large) $n \in \mathbb{N}$ if and only if $G$ is finite, in which case (4.1) clearly is correct. Thus, assume $G$ to be infinite from now on, and consequently $\ell_{n}>0$ for all $n \in \mathbb{N}$.

Given $n \in \mathbb{N}$, choose $p_{n} \in \Pi_{n}$ as in Proposition 4.2 and notice that $\ell_{n}>0$ implies $\min _{j=1}^{n}\left(P_{n, j}-P_{n, j-1}\right) \geq 2 \ell_{n}>0$; in particular, $P_{n, j-1}<P_{n, j}$ for all $j=1, \ldots, n$. Consequently, for every $x \in\left[0,1\left[\right.\right.$ there exists a unique $j_{n}(x) \in\{1, \ldots, n\}$ with $P_{n, j_{n}(x)-1} \leq x<P_{n, j_{n}(x)}$. For convenience, let $J_{n, j}=\left[P_{n, j-1}+\ell_{n}, P_{n, j}-\ell_{n}\right]$ for all $j=1, \ldots, n$, and hence $\lambda_{n, j}:=\lambda\left(J_{n, j}\right)=P_{n, j}-P_{n, j-1}-2 \ell_{n}$. Next, recall that the set $U:=[0,1] \backslash G$ is open, possibly empty. If $U \neq \varnothing$, let $I_{1}, I_{2}, \ldots$ be its (at most countably many) connected components, that is, the disjoint open intervals with endpoints in $G$ and $U=\bigcup_{k} I_{k}$. Thus, for every $x \in U$ there exists a unique $k(x) \in\{1,2, \ldots\}$ with $x \in I_{k(x)}$. Finally, consider the subset $G^{\dagger}$ of $G$ defined as

$$
\begin{equation*}
G^{\dagger}=\{x \in G: I \cap G=\{x\} \text { for some interval } I \text { with } \lambda(I)>0\} . \tag{4.2}
\end{equation*}
$$

Notice that $\{0,1\} \subset G^{\dagger}$, and $G^{\dagger}$ is (at most) countable. Utilizing Proposition 4.2, it is readily checked that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left[P_{n, j_{n}(x)-1}, P_{n, j_{n}(x)}\right]=\{x\} \quad \forall x \in G \backslash G^{\dagger}, \tag{4.3}
\end{equation*}
$$

but also

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left[P_{n, j_{n}(x)-1}, P_{n, j_{n}(x)}\right]=I_{k(x)} \quad \forall x \in U \tag{4.4}
\end{equation*}
$$

With these preparations, the proof is now carried out in three separate steps for the reader's convenience.
Step I: Assume $\mu^{-1}$ is absolutely continuous.
Proposition 4.2 with $\epsilon=1$ yields $\mu^{-1}\left(J_{n, j}\right)=2 \ell_{n}$ or, equivalently,

$$
\begin{equation*}
g\left(P_{n, j}-\ell_{n}\right)-\ell_{n}=g\left(P_{n, j-1}+\ell_{n}\right)+\ell_{n} \quad \forall j=1, \ldots, n, \tag{4.5}
\end{equation*}
$$

and since $\mu^{-1}$ is absolutely continuous also $\lambda_{n, j}>0$. Fix any $0<a<1$, and recalling that $g$ is differentiable $\lambda$-almost everywhere, with $g^{\prime} \geq 0$ integrable over every compact subinterval of $] 0,1\left[\right.$, pick a nonnegative continuous function $f_{a}$ : $] 0,1\left[\rightarrow \mathbb{R}\right.$ with $\int_{[0,1]}\left|g^{\prime}-f_{a}\right| \mathrm{d} \lambda<a$ (notice that $g^{\prime}, f_{a}$ may not be integrable over $[0,1])$. If $U \neq \varnothing$, then also pick $k_{a} \in \mathbb{N}$ large enough to ensure $\lambda\left(\bigcup_{k>k_{a}} I_{k}\right)<a$, and for every $k=1, \ldots, k_{a}$ pick a continuous function $\left.e_{k}:\right] 0,1\left[\rightarrow[0,1]\right.$ with $e_{k}(x)=1$ for all $x \in] 0,1\left[\backslash I_{k}\right.$ such that $\int_{I_{k}} f_{a} e_{k}<a \lambda\left(I_{k}\right)$. Let $f=f_{a} \prod_{k=1}^{k_{a}} e_{k}$; in case $U=\varnothing$, simply let $f=f_{a}$. Clearly, $f$ is nonnegative and continuous on $] 0,1[$, with

$$
\begin{equation*}
\int_{[0,1]}\left|g^{\prime}-f\right| \mathrm{d} \lambda=\int_{G}\left|g^{\prime}-f_{a}\right| \mathrm{d} \lambda+\int_{U}|f| \mathrm{d} \lambda \leq \int_{[0,1]}\left|g^{\prime}-f_{a}\right| \mathrm{d} \lambda<a \tag{4.6}
\end{equation*}
$$

since $g^{\prime}$ vanishes on $U$. Next, deduce from (4.5) that

$$
\begin{aligned}
\frac{1}{\lambda_{n, j}} \int_{J_{n, j}} f \mathrm{~d} \lambda & =\frac{2 \ell_{n}}{\lambda_{n, j}}+\frac{1}{\lambda_{n, j}} \int_{J_{n, j}}\left(f-g^{\prime}\right) \mathrm{d} \lambda \\
& =\frac{2 \ell_{n}-\int_{J_{n, j}}\left(g^{\prime}-f\right) \mathrm{d} \lambda}{P_{n, j}-P_{n, j-1}-2 \ell_{n}} \quad \forall j=1, \ldots, n,
\end{aligned}
$$

and consequently, for every $j=1, \ldots, n$,

$$
\begin{equation*}
2 \Omega\left(\frac{1}{\lambda_{n, j}} \int_{J_{n, j}} f \mathrm{~d} \lambda\right)\left(P_{n, j}-P_{n, j-1}\right)=2 \ell_{n}-\frac{\int_{J_{n, j}}\left(g^{\prime}-f\right) \mathrm{d} \lambda}{1+\frac{1}{\lambda_{n, j}} \int_{J_{n, j}} f \mathrm{~d} \lambda} . \tag{4.7}
\end{equation*}
$$

Summing (4.7) over $j=1, \ldots, n$ yields

$$
\begin{equation*}
\omega_{n}-\int_{[0,1]} h_{n} \mathrm{~d} \lambda=\frac{1}{2} \sum_{j=1}^{n} \frac{\int_{J_{n, j}}\left(g^{\prime}-f\right) \mathrm{d} \lambda}{1+\frac{1}{\lambda_{n, j}} \int_{J_{n, j}} f \mathrm{~d} \lambda} \tag{4.8}
\end{equation*}
$$

with the piecewise constant nonnegative function $h_{n}:[0,1[\rightarrow \mathbb{R}$ given by

$$
h_{n}(x)=\Omega\left(\frac{1}{\lambda_{n, j_{n}(x)}} \int_{J_{n, j_{n}(x)}} f \mathrm{~d} \lambda\right) \quad \forall x \in[0,1[.
$$

Recall that $f \geq 0$, and so the right-hand side in (4.8) is bounded, for every $n \in \mathbb{N}$, by $\frac{1}{2} \sum_{j=1}^{n} \int_{J_{n, j}}\left|g^{\prime}-f\right| \mathrm{d} \lambda \leq \frac{1}{2} \int_{[0,1]}\left|g^{\prime}-f\right| \mathrm{d} \lambda<\frac{1}{2} a$. Deduce from (4.3) and the continuity of $f$ that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} h_{n}(x)=\Omega(f(x)) \quad \forall x \in G \backslash G^{\dagger} \tag{4.9}
\end{equation*}
$$

Similarly, (4.4) and the choice of the functions $e_{k}$ for $k=1, \ldots, k_{a}$ imply that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} h_{n}(x)=\Omega\left(\frac{1}{\lambda\left(I_{k(x)}\right)} \int_{I_{k(x)}} f \mathrm{~d} \lambda\right) \leq \Omega(a)<\frac{a}{2} \quad \forall x \in \bigcup_{k=1}^{k_{a}} I_{k} \tag{4.10}
\end{equation*}
$$

The elementary estimate, valid for all $n \in \mathbb{N}$,

$$
\begin{aligned}
\left|\omega_{n}-\int_{[0,1]} \Omega\left(g^{\prime}\right) \mathrm{d} \lambda\right| \leq & \left|\omega_{n}-\int_{[0,1]} h_{n} \mathrm{~d} \lambda\right|+\int_{\bigcup_{k=1}^{k_{a}} I_{k}} h_{n} \mathrm{~d} \lambda+\int_{\bigcup_{k>k_{a}} I_{k}} h_{n} \mathrm{~d} \lambda \\
& +\int_{[0,1]}\left|\Omega(f)-\Omega\left(g^{\prime}\right)\right| \mathrm{d} \lambda+\int_{G}\left|h_{n}-\Omega(f)\right| \mathrm{d} \lambda
\end{aligned}
$$

together with (4.8), Fatou's lemma and (4.10), the choice of $k_{a}$ and the fact that $0 \leq h_{n} \leq \frac{1}{2}$, estimate (4.6) as well as the dominated convergence theorem yield

$$
\limsup _{n \rightarrow \infty}\left|\omega_{n}-\int_{[0,1]} \Omega\left(g^{\prime}\right) \mathrm{d} \lambda\right| \leq \frac{a}{2}+\frac{a}{2}+\frac{a}{2}+\frac{a}{2}+0=2 a .
$$

Since $0<a<1$ has been arbitrary, $\lim _{n \rightarrow \infty} \omega_{n}=\int_{[0,1]} \Omega\left(g^{\prime}\right) \mathrm{d} \lambda$, i.e. (4.1) holds.
Step II: Assume $\mu^{-1}$ is singular.
Given any $0<a<1$, let $U_{a}=\{x \in[0,1]: \operatorname{dist}(x, G) \geq a\}$. Note that $U_{a} \subset$ $[a, 1-a]$ is a compact, possibly empty subset of $U$, so $U_{a} \cap G=\varnothing$. Assume for the time being that all atoms of $\mu^{-1}$ in $[a, 1-a]$ are small in that

$$
\begin{equation*}
\mu^{-1}(\{x\}) \leq a^{2} \quad \forall x \in[a, 1-a] . \tag{4.11}
\end{equation*}
$$

Recall that $2 \ell_{n} \leq \mu^{-1}\left(J_{n, j}\right)$ for all $j=1, \ldots, n$, by Proposition 4.2, and correspondingly

$$
\omega_{n} \leq \sum_{j=1}^{n} \Omega\left(\frac{\mu^{-1}\left(J_{n, j}\right)}{\lambda_{n, j}}\right)\left(P_{n, j}-P_{n, j-1}\right)=\int_{[0,1]} \widetilde{h}_{n} \mathrm{~d} \lambda,
$$

with the piecewise constant function $\widetilde{h}_{n}:[0,1[\rightarrow \mathbb{R}$ given by

$$
\widetilde{h}_{n}(x)=\Omega\left(\frac{\mu^{-1}\left(J_{n, j_{n}(x)}\right)}{\lambda_{n, j_{n}(x)}}\right) \quad \forall x \in[0,1[.
$$

First, observe that if $x \in U_{a}$, then (4.4) and (4.11) imply that

$$
\limsup _{n \rightarrow \infty} \mu^{-1}\left(J_{n, j_{n}(x)}\right) \leq \mu^{-1}\left(I_{k(x)}\right) \leq 2 a^{2}
$$

whereas clearly $\lim _{n \rightarrow \infty} \lambda_{n, j_{n}(x)}=\lambda\left(I_{k(x)}\right) \geq 2 a$. Thus,

$$
\begin{equation*}
\lim \sup _{n \rightarrow \infty} \widetilde{h}_{n}(x) \leq \Omega(a)<\frac{a}{2} \quad \forall x \in U_{a} \tag{4.12}
\end{equation*}
$$

Next, notice that if $x \in G \backslash G^{\dagger}$, then $\left(\left[P_{n, j_{n}(x)-1}, P_{n, j_{n}(x)}\right]\right)$ shrinks to $x$ nicely, and hence

$$
\lim _{n \rightarrow \infty} \frac{\mu^{-1}\left(\left[P_{n, j_{n}(x)-1}, P_{n, j_{n}(x)}\right]\right)}{P_{n, j_{n}(x)}-P_{n, j_{n}(x)-1}}=0 \quad \text { for } \lambda \text {-almost every } x \in G
$$

by [33, Thm. 7.13]. Thus, $\lim _{n \rightarrow \infty} \ell_{n} /\left(P_{n, j_{n}(x)}-P_{n, j_{n}(x)-1}\right)=0$ for $\lambda$-almost every $x \in G$, which in turn shows that $\left(J_{n, j_{n}(x)}\right)$ shrinks to $x$ nicely as well. Applying [33, Thm. 7.13] once more yields $\lim _{n \rightarrow \infty} \mu^{-1}\left(J_{n, j_{n}(x)}\right) / \lambda_{n, j_{n}(x)}=0$ for $\lambda$-almost every $x \in G$, and thus

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \widetilde{h}_{n}(x)=0 \text { for } \lambda \text {-almost every } x \in G \tag{4.13}
\end{equation*}
$$

Recalling that $G^{\dagger}$ is countable, deduce from (4.12) and (4.13) that

$$
\begin{align*}
\omega^{+} & \leq \lim \sup _{n \rightarrow \infty} \int_{[0,1]} \widetilde{h}_{n} \mathrm{~d} \lambda \\
& =\lim \sup _{n \rightarrow \infty}\left(\int_{U_{a}} \widetilde{h}_{n} \mathrm{~d} \lambda+\int_{G} \widetilde{h}_{n} \mathrm{~d} \lambda+\int_{[0,1] \backslash\left(U_{a} \cup G\right)} \widetilde{h}_{n} \mathrm{~d} \lambda\right) \\
& \leq \frac{a}{2}+\frac{1}{2}\left(1-\lambda\left(U_{a} \cup G\right)\right) . \tag{4.14}
\end{align*}
$$

In summary, (4.14) holds provided that $\mu$ satisfies (4.11).
To conclude the argument in the case of $\mu^{-1}$ being singular, given $0<b<1$, pick $0<a<b$ so small that $\lambda\left(U_{a} \cup G\right)>1-b$. Noting that the set $G_{a}:=\{x \in$ $\left.[a, 1-a]: \mu^{-1}(\{x\})>a^{2}\right\}$ is finite, consider $\widetilde{g} \in \mathcal{F}$ given by

$$
\tilde{g}=g-\sum_{x \in G_{a}} \mu^{-1}(\{x\}) \mathbf{1}_{[x,+\infty[ },
$$

as well as the unique $\tilde{\mu} \in \mathcal{P}$ with $F_{\tilde{\mu}}^{-1}=\widetilde{g}$. Crucially, (4.11) holds with $\tilde{\mu}$ instead of $\mu$. Moreover, notice that $\widetilde{G}:=G_{\widetilde{\mu}^{-1}} \supset G \backslash G_{a}$, and clearly $\widetilde{U}_{a} \supset U_{a}$, where $\widetilde{U}_{a}=\{x \in[0,1] \dot{\sim} \operatorname{dist}(x, \widetilde{G}) \geq a\}$. Thus, $\widetilde{U}_{a} \cup \widetilde{G} \supset\left(U_{a} \cup G\right) \backslash G_{a}$, and (4.14) applied to $\widetilde{\mu}$, with $\widetilde{\ell}_{n}:=d_{1}\left(\tilde{\mu}, \delta_{\bullet}^{\bullet, n}\right)$, yields

$$
\lim \sup _{n \rightarrow \infty} n \tilde{\ell}_{n} \leq \frac{a}{2}+\frac{1}{2}\left(1-\lambda\left(\widetilde{U}_{a} \cup \widetilde{G}\right)\right)<b
$$

Finally, let $m_{a}=\# G_{a}$ and observe that $\ell_{n+m_{a}} \leq \widetilde{\ell}_{n}$ for all $n \in \mathbb{N}$, so

$$
\omega_{n} \leq n \tilde{\ell}_{n-m_{a}}=\left(n-m_{a}\right) \tilde{\ell}_{n-m_{a}}+m_{a} \widetilde{\ell}_{n-m_{a}}
$$

for all $n>m_{a}$. Since $\lim _{n \rightarrow \infty} \tilde{\ell_{n}}=0$, clearly $\omega^{+} \leq \lim \sup _{n \rightarrow \infty} n \tilde{\ell}_{n}<b$, and since $0<b<1$ has been arbitrary, $\omega^{+}=0$. Thus, (4.1) holds, with vanishing right-hand side, whenever $\mu^{-1}$ is singular.

Step III: Let $\mu \in \mathcal{P}$ be arbitrary.
As in the proof of Theorem 3.5, write $g=g_{\mathrm{A}}+g_{\mathrm{S}}$ with $g_{\mathrm{A}}, g_{\mathrm{S}} \in \mathcal{F}$ such that $\lambda_{g_{\mathrm{A}}}=\left(\mu^{-1}\right)_{\mathrm{A}}$ and $\lambda_{g \mathrm{~S}}=\left(\mu^{-1}\right)_{\mathrm{S}}$. Let $\mu^{\langle\mathrm{A}\rangle}$ and $\mu^{\langle\mathrm{S}\rangle}$ be the (uniquely determined) probability measures with $\left(\mu^{\langle\mathrm{A}\rangle}\right)^{-1}=\left(\mu^{-1}\right)_{\mathrm{A}}$ and $\left(\mu^{\langle\mathrm{S}\rangle}\right)^{-1}=\left(\mu^{-1}\right)_{\mathrm{S}}$, respectively (notice that in general $\mu^{\langle\mathrm{A}\rangle} \neq \mu_{\mathrm{A}}$ and $\mu^{\langle\mathrm{S}\rangle} \neq \mu_{\mathrm{S}}$ ). Also, for every $n \in \mathbb{N}$ let $\ell_{n}^{\langle\mathrm{A}\rangle}=$ $d_{1}\left(\mu^{\langle\mathrm{A}\rangle}, \delta_{\bullet}^{\bullet}, n\right)$ and $\ell_{n}^{\langle\mathrm{S}\rangle}=d_{1}\left(\mu^{\langle\mathrm{S}\rangle}, \delta_{\bullet}^{\bullet, n}\right)$. Given any $m, n \in \mathbb{N}$, pick $p_{m}^{\langle\mathrm{A}\rangle} \in \Pi_{m}$ and $p_{n}^{\langle\mathrm{S}\rangle} \in \Pi_{n}$ as in Proposition 4.2. By considering the joint partition of [0, 1] generated by $\left\{P_{m, i}^{\langle\mathrm{A}\rangle}: i=0, \ldots, m\right\}$ and $\left\{P_{n, j}^{\langle\mathrm{S}\rangle}: j=0, \ldots, n\right\}$, it is readily seen that $\ell_{m+n} \leq$ $\ell_{m}^{\langle\mathrm{A}\rangle}+\ell_{n}^{(\mathrm{S}\rangle}$. For every $0<a<1$ and $n \in \mathbb{N}$, therefore

$$
\omega_{n} \leq n \ell_{\lfloor(1-a) n\rfloor+\lfloor a n\rfloor} \leq \frac{1+\lfloor(1-a) n\rfloor}{1-a} \ell_{\lfloor(1-a) n\rfloor}^{\langle\mathrm{A}\rangle}+\frac{1+\lfloor a n\rfloor}{a} \ell_{\lfloor a n\rfloor}^{\langle\mathrm{S}\rangle},
$$

and applying Steps I and II to $\mu^{\langle\mathrm{A}\rangle}$ and $\mu^{\langle\mathrm{S}\rangle}$, respectively, yields

$$
\omega^{+} \leq \frac{1}{1-a} \int \Omega\left(\frac{\mathrm{~d}\left(\mu^{\langle\mathrm{A}\rangle}\right)_{\mathrm{A}}^{-1}}{\mathrm{~d} \lambda}\right) \mathrm{d} \lambda=\frac{1}{1-a} \int \Omega\left(g_{\mathrm{A}}^{\prime}\right) \mathrm{d} \lambda .
$$

(recall that $\lim _{n \rightarrow \infty} \ell_{n}^{\langle\mathrm{A}\rangle}=\lim _{n \rightarrow \infty} n \ell_{n}^{\langle\mathrm{S}\rangle}=0$ ). Since $0<a<1$ has been arbitrary, $\omega^{+} \leq \int \Omega\left(g_{\mathrm{A}}^{\prime}\right) \mathrm{d} \lambda$. To obtain a lower bound for $\omega^{-}$, recall from Proposition 4.2 that

$$
g_{-}\left(P_{n, j}-\ell_{n}\right)-\ell_{n} \leq g\left(P_{n, j-1}+\ell_{n}\right)+\ell_{n} \quad \forall j=1, \ldots, n,
$$

and since $g_{-}=g_{\mathrm{A}}+g_{\mathrm{S}_{-}}$,

$$
\begin{aligned}
g_{\mathrm{A}}\left(P_{n, j}-\ell_{n}\right)-\ell_{n} \leq & g_{\mathrm{A}}\left(P_{n, j-1}+\ell_{n}\right)+\ell_{n} \\
& -\max \left\{0,\left(g_{\mathrm{S}-}\left(P_{n, j}-\ell_{n}\right)-g_{\mathrm{S}}\left(P_{n, j-1}+\ell_{n}\right)\right)\right\} \\
\leq & g_{\mathrm{A}}\left(P_{n, j-1}+\ell_{n}\right)+\ell_{n} \quad \forall j=1, \ldots, n,
\end{aligned}
$$

from which it is clear that $\ell_{n}^{\langle\mathrm{A}\rangle} \leq \ell_{n}$ for every $n \in \mathbb{N}$. Applying Step I to $\mu^{\langle\mathrm{A}\rangle}$ yields $\omega_{-} \geq \liminf _{n \rightarrow \infty} n \ell_{n}^{\langle\mathrm{A}\rangle}=\int \Omega\left(g_{\mathrm{A}}^{\prime}\right) \mathrm{d} \lambda$. Hence, $\lim _{n \rightarrow \infty} \omega_{n}=\int \Omega\left(g_{\mathrm{A}}^{\prime}\right) \mathrm{d} \lambda$, and the proof is complete.

Along the lines of the above proof, and by considering the absolutely continuous part of $\mu$ rather than of $\mu^{-1}$, the following dual version of Theorem 4.1 can be established; the routine details are left to the interested reader.

Proposition 4.3 Let $\mu \in \mathcal{P}$ and $\epsilon>0$. Then,

$$
\lim _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=\epsilon \int \Omega\left(\frac{1}{\epsilon} \cdot \frac{\mathrm{~d} \mu_{\mathrm{A}}}{\mathrm{~d} \lambda}\right) \mathrm{d} \lambda
$$

Notice that Theorem 4.1 and Proposition 4.3 together imply the familiar fact [6, Ch. A] that $\mu^{-1}$ is singular if and only if $\mu$ is singular, and hence yield a direct analogue of Proposition 3.2 in the context of best approximations.

Proposition 4.4 For every $\mu \in \mathcal{P}$ and $\epsilon>0$, the following are equivalent:
(i) $\lim _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=0$;
(ii) $\mu_{\mathrm{A}}=0$;
(iii) $\left(\mu^{-1}\right)_{\mathrm{A}}=0$.

Just as in the case of best uniform approximations, Theorem 4.1 can be refined through further assumptions on $\mu$. For instance, if $g=F_{\mu}^{-1}$ is $C^{4}$ on $] 0,1\left[\right.$, if both $g, g^{\prime} \neq 0$ are convex and if $\Omega\left(\epsilon g^{\prime}\right)$ has a $C^{3}$-extension to $\mathbb{R}$, then mild boundedness assumptions on $g$ and its derivatives (ensuring all relevant integrals are finite) guarantee that, as a refinement of (4.1),

$$
\begin{equation*}
n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=c_{1}+\frac{c_{1}^{2} c_{2}}{12} n^{-2}+o\left(n^{-2}\right) \quad \text { as } n \rightarrow \infty \tag{4.15}
\end{equation*}
$$

where $c_{1}=\int \Omega\left(\epsilon g^{\prime}\right) \mathrm{d} \lambda$ and

$$
c_{2}=\int \frac{2\left(1+\epsilon g^{\prime}\right)\left(g^{\prime \prime}\right)^{2}-\left(2+\epsilon g^{\prime}\right) g^{\prime} g^{\prime \prime \prime}}{\left(1+\epsilon g^{\prime}\right)^{2}\left(g^{\prime}\right)^{2}} \mathrm{~d} \lambda
$$

When compared to $\left(n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)\right)$, therefore, not only does $\left(n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet}, n\right)\right)$ converge to a smaller value (unless $g^{\prime}$ is constant), but also it converges at the rate $\left(n^{-2}\right)$ which often is faster than the rate in (3.10). For example, (4.15) applies to exponential as well as Benford distributions, and the reader may want to check that (2.7) and (2.9) both are consistent with it. If $\mu$ is a normal distribution with variance $\sigma^{2}>0$, then $\Omega\left(\epsilon g^{\prime}\right)$ does not have even a $C^{1}$-extension to $\mathbb{R}$, and correspondingly $c_{2}=-\infty$, which suggests that $\left(n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)\right)$ converges at a slower rate. This indeed is the case: an elementary albeit lengthy analysis yields

$$
n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=-\epsilon \sqrt{\frac{\pi \sigma^{2}}{2}} \operatorname{Li}_{1 / 2}\left(-\frac{1}{\epsilon \sqrt{2 \pi \sigma^{2}}}\right)+\mathcal{O}\left(\frac{\log n}{n^{2}}\right) \quad \text { as } n \rightarrow \infty
$$

where $\mathrm{Li}_{1 / 2}$ denotes the polylogarithm of order $\frac{1}{2}$; see, e.g. [26, §25.12]. Though slower than (4.15), this rate of convergence again is considerably faster than its counterpart (3.9) for best uniform approximations. It should be noted, however, that such a hierarchy of rates, though observed for many familiar distributions, is by no means universal: as mentioned already in Introduction, for the 1-Pareto distribution both sequences $\left(n d_{\epsilon}\left(\mu, \delta_{\bullet}^{u_{n}}\right)\right)$ and $\left(n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)\right)$ converge to their respective limits $\frac{1}{2}$ and $\frac{\pi}{8}$ at the same rate $\left(n^{-2}\right)$, as is evident from (1.7) and (1.8). For the $\frac{1}{2}$-Pareto distribution, i.e. $F_{\mu}(x)=1-x^{-1 / 2}$ for all $x \geq 1$, (3.10) yields

$$
n d_{1}\left(\mu, \delta_{\bullet}^{u_{n}}\right)=0.5000-0.03125 n^{-3}+o\left(n^{-3}\right) \quad \text { as } n \rightarrow \infty,
$$

whereas (4.15), with $c_{1}=\int_{0}^{1}\left(2+t^{3}\right)^{-1} \mathrm{~d} t=0.4508$ and $c_{2}=6 \int_{0}^{1} t\left(2-t^{3}\right)(2+$ $\left.t^{3}\right)^{-2} \mathrm{~d} t=0.9102$, reads

$$
n d_{1}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=0.4508+0.01541 n^{-2}+o\left(n^{-2}\right) \quad \text { as } n \rightarrow \infty
$$

Here, $\left(n d_{1}\left(\mu, \delta_{\bullet}^{u_{n}}\right)\right)$ converges at an even faster rate than $\left(n d_{1}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)\right)$.
Example 6 For the Cantor distribution $\mu$ and its inverse $\mu^{-1}$ in Example 5, Theorem 4.1 yields $\lim _{n \rightarrow \infty} n d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)=\lim _{n \rightarrow \infty} n d_{\epsilon}\left(\mu^{-1}, \delta_{\bullet}^{\bullet, n}\right)=0$. An elementary analysis shows that $\left(n^{1 / c} d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet, n}\right)\right)$ and $\left(n^{1 / c} d_{\epsilon}\left(\mu^{-1}, \delta_{\bullet}^{\bullet}, n\right)\right)$ both are divergent, yet bounded above and below by positive constants, where $c=\log 2 / \log 3<1$ is the Hausdorff dimension of both the set $C=G_{\mu}$ and the measure $\mu$. It seems plausible that Theorem 4.1 can similarly be refined for a wide class of self-similar (singular) distributions, thus complementing known $d_{\mathrm{W}}$-quantization results [18,19,25,29].

To establish one further interesting property of best $d_{\epsilon}$-approximations, recall from Proposition 2.5 that if $d_{\epsilon}\left(\mu, \delta_{x_{n}}^{p_{n}}\right)=\ell_{\bullet}^{\bullet, n}$, then $p_{n}$ can easily be determined from $x_{n}$ (or vice versa). Thus, $x_{n}$ (or $p_{n}$ ) alone already captures $\delta_{x_{n}}^{p_{n}}$ to a large extent, and it is natural to ask, for instance, whether $x_{n, 1}, \ldots, x_{n, n}$, i.e. the locations of best $d_{\epsilon}$-approximations of $\mu \in \mathcal{P}$ conform to an asymptotic point distribution as $n \rightarrow \infty$, referred to as the point density measure of $\mu$ in [19]. In the context of best $d_{\mathrm{W}}$-approximations (or quantizations), and under the appropriate assumptions, this question has a positive answer; see, e.g. the "empirical measure theorem" $[18$, Thm. 7.5] and variants thereof [19]. As is the case with Theorem 4.1 and Proposition 4.3, the result for best $d_{\epsilon^{-}}$ approximations again is simpler than its $d_{\mathrm{W}}$-counterpart in that the asymptotic point distribution is readily identified whenever $\mu \in \mathcal{P}$ is non-singular, and no further assumptions on $\mu$ are needed. In fact, it even is possible to allow for slightly more general $x_{n}$. To concisely state the result, for every $\mu \in \mathcal{P}$ with $\mu_{\mathrm{A}} \neq 0$, define $\mu_{\epsilon}^{*} \in \mathcal{P}$ via

$$
\frac{\mathrm{d} \mu_{\epsilon}^{*}}{\mathrm{~d} \lambda}=\frac{\Omega\left(\frac{1}{\epsilon} \cdot \frac{\mathrm{~d} \mu_{\mathrm{A}}}{\mathrm{~d} \lambda}\right)}{\int_{\mathbb{R}} \Omega\left(\frac{1}{\epsilon} \cdot \frac{\mathrm{~d} \mu_{\mathrm{A}}}{\mathrm{~d} \lambda}\right) \mathrm{d} \lambda} \quad \forall \epsilon>0
$$

Clearly, $\mu_{\epsilon}^{*}$ is absolutely continuous, and $\mu_{\epsilon}^{*}=\mu$ for some (in fact, all) $\epsilon>0$ if and only if $\mu$ is uniform, i.e. $\mu=\lambda(\cdot \cap B) / \lambda(B)$ for some Borel set $B$ with $\lambda(B) \in \mathbb{R}^{+}$. Also, given any $\mu \in \mathcal{P} \backslash \mathcal{P}_{\infty}^{*}$, i.e. \# supp $\mu=\infty$, call a sequence ( $x_{n}$ ), with $x_{n} \in \Xi_{n}$ for every $n \in \mathbb{N}$, asymptotically $d_{\epsilon}$-minimal for $\mu$ if

$$
\lim _{n \rightarrow \infty} \frac{d_{\epsilon}\left(\mu, \delta_{x_{n}}^{\bullet}\right)}{d_{\epsilon}\left(\mu, \delta_{\bullet}^{\bullet}, n\right.}=1
$$

Thus, for instance, $\left(x_{n}\right)$ is asymptotically $d_{\epsilon}$-minimal for $\mu \in \mathcal{P} \backslash \mathcal{P}_{\infty}^{*}$ whenever $\delta_{x_{n}}^{p_{n}}$, with $x_{n} \in \Xi_{n}, p_{n} \in \Pi_{n}$, is a best $d_{\epsilon}$-approximation of $\mu$ for every $n \in \mathbb{N}$.

Theorem 4.5 Let $\mu \in \mathcal{P}$ and $\epsilon>0$. If $\mu_{\mathrm{A}} \neq 0$ and $\left(x_{n}\right)$ is asymptotically $d_{\epsilon}$-minimal for $\mu$, then

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \frac{\#\left\{1 \leq j \leq n: x_{n, j} \in I\right\}}{n}=\mu_{\epsilon}^{*}(I) \quad \forall I \subset \mathbb{R}, I \text { an interval. } \tag{4.16}
\end{equation*}
$$

Proof For convenience, write $f=F_{\mu}=f_{\mathrm{A}}+f_{\mathrm{S}}$ with $f_{\mathrm{A}}, f_{\mathrm{S}} \in \mathcal{F}$ such that $\lambda_{f_{\mathrm{A}}}=\mu_{\mathrm{A}}$ and $\lambda_{f_{\mathrm{S}}}=\mu_{\mathrm{S}}$ (the functions $f_{\mathrm{A}}, f_{\mathrm{S}}$ can be made unique, for instance, by requiring that $\left.f_{\mathrm{A}}(-\infty)=f_{\mathrm{S}}(-\infty)=0\right)$. Also, let $G=G_{\mu}, \ell_{n}=d_{\epsilon}\left(\mu, \delta_{x_{n}}^{\bullet}\right)$ for all $n \in \mathbb{N}$, and define $G^{\dagger}$ as in (4.2). Once again, it suffices to consider the case of $\epsilon=1$. Note that $\mu_{\mathrm{A}} \neq 0$ implies $\ell_{n}>0$ for every $n$, and $\lim _{n \rightarrow \infty} n \ell_{n}=\int_{\mathbb{R}} \Omega\left(\mathrm{d} \mu_{\mathrm{A}} / \mathrm{d} \lambda\right)>0$, by Proposition 4.3 and the assumed asymptotic $d_{\epsilon}$-minimality of $\left(x_{n}\right)$.

Fix a non-empty interval $I=] y, z]$ with $y, z \in \mathbb{R}$. Perturbing $x_{n}$ slightly if necessary, without altering $\#\left\{1 \leq j \leq n: x_{n, j} \in I\right\}$ or increasing $d_{1}\left(\mu, \delta_{x_{n}}^{\bullet}\right)$ for any $n$, it may be assumed that $x_{n, j}<x_{n, j+1}$ for all $n \in \mathbb{N}$ and $j=0, \ldots, n$. Thus, for every $x \in \mathbb{R}$, there exists a unique $j_{n}(x) \in\{0, \ldots, n\}$ with $x_{n, j_{n}(x)} \leq x<x_{n, j_{n}(x)+1}$. By Proposition 2.5,
$f_{\mathrm{A}}\left(x_{n, j+1}-\ell_{n}\right)+f_{\mathrm{S}-}\left(x_{n, j+1}-\ell_{n}\right)-\ell_{n} \leq f_{\mathrm{A}}\left(x_{n, j}+\ell_{n}\right)+f_{\mathrm{S}}\left(x_{n, j}+\ell_{n}\right)+\ell_{n}$,
for every $j=0, \ldots, n$, and consequently also

$$
\begin{equation*}
f_{\mathrm{A}}\left(x_{n, j+1}-\ell_{n}\right)-f_{\mathrm{A}}\left(x_{n, j}+\ell_{n}\right) \leq 2 \ell_{n} \quad \forall j=0, \ldots, n \tag{4.17}
\end{equation*}
$$

Fix any $a>0$, and recalling that $f_{\mathrm{A}}$ is differentiable $\lambda$-almost everywhere with $f_{\mathrm{A}}^{\prime} \geq 0$ and $0<\int_{\mathbb{R}} f_{\mathrm{A}}^{\prime} \mathrm{d} \lambda=\mu_{\mathrm{A}}(\mathbb{R}) \leq 1$, pick a continuous function $g: \mathbb{R} \rightarrow \mathbb{R}^{+}$ with $\int_{\mathbb{R}}\left|f_{\mathrm{A}}^{\prime}-g\right| \mathrm{d} \lambda<a$.

Let $K_{n}=\left\{0 \leq j \leq n: x_{n, j}+\ell_{n}<x_{n, j+1}-\ell_{n}\right\}$ which may not be all of the set $\{0, \ldots, n\}$ but does contain $0, n$ in any case. On the one hand, if $j \in K_{n} \backslash\{0, n\}$ let $J_{n, j}=\left[x_{n, j}+\ell_{n}, x_{n, j+1}-\ell_{n}\right]$ and $\lambda_{n, j}=\lambda\left(J_{n, j}\right)=x_{n, j+1}-x_{n, j}-2 \ell_{n}>0$, and deduce from (4.17) that

$$
\frac{1}{\lambda_{n, j}} \int_{J_{n, j}} g \mathrm{~d} \lambda \leq \frac{2 \ell_{n}-\int_{J_{n, j}}\left(f_{\mathrm{A}}^{\prime}-g\right) \mathrm{d} \lambda}{x_{n, j+1}-x_{n, j}-2 \ell_{n}},
$$

and consequently

$$
\begin{equation*}
\ell_{n} \geq \Omega\left(\frac{1}{\lambda_{n, j}} \int_{J_{n, j}} g \mathrm{~d} \lambda\right)\left(x_{n, j+1}-x_{n, j}\right)-\frac{1}{2} \int_{J_{n, j}}\left|f_{\mathrm{A}}^{\prime}-g\right| \mathrm{d} \lambda ; \tag{4.18}
\end{equation*}
$$

with the usual convention $0 \cdot( \pm \infty)=0$, (4.18) is correct also for $j=0, n$. On the other hand, if $j \notin K_{n}$ then clearly $\ell_{n} \geq \frac{1}{2}\left(x_{n, j+1}-x_{n, j}\right)$. With (4.18) as well as the definitions of $j_{n}(x)$ and $K_{n}$, therefore,

$$
\begin{align*}
(1 & \left.+j_{n}(z)-j_{n}(y)\right) \ell_{n} \\
& \geq \sum_{j \in\left\{j_{n}(y), \ldots, j_{n}(z)\right\} \cap K_{n}} \Omega\left(\frac{1}{\lambda_{n, j}} \int_{J_{n, j}} g \mathrm{~d} \lambda\right)\left(x_{n, j+1}-x_{n, j}\right) \\
& -\frac{1}{2} \sum_{j \in\left\{j_{n}(y), \ldots, j_{n}(z)\right\} \cap K_{n}} \int_{J_{n, j}}\left|f_{\mathrm{A}}^{\prime}-g\right| \mathrm{d} \lambda \\
& +\sum_{j \in\left\{j_{n}(y), \ldots, j_{n}(z)\right\} \backslash K_{n}} \frac{1}{2}\left(x_{n, j+1}-x_{n, j}\right) \\
& \geq \int_{\left[x_{\left.n, j_{n}(y), x_{n, j_{n}(z)+1}\right]} h_{n} \mathrm{~d} \lambda-\frac{1}{2} \int_{\left[x_{n, j_{n}(y)}, x_{n, j_{n}(z)+1}\right]}\left|f_{\mathrm{A}}^{\prime}-g\right| \mathrm{d} \lambda\right.} \quad \geq \int_{I} h_{n} \mathrm{~d} \lambda-\frac{a}{2},
\end{align*}
$$

where the piecewise constant function $h_{n}: \mathbb{R} \rightarrow \mathbb{R}^{+}$is given by

$$
h_{n}(x)= \begin{cases}\Omega\left(\frac{1}{\lambda_{n, j_{n}(x)}} \int_{J_{n, j_{n}(x)}} g \mathrm{~d} \lambda\right) & \text { if } j_{n}(x) \in K_{n} \\ \frac{1}{2} & \text { if } j_{n}(x) \notin K_{n}\end{cases}
$$

If $j_{n}(x) \notin K_{n}$ for all sufficiently large $n$, then $\lim _{n \rightarrow \infty} h_{n}(x)=\frac{1}{2}$, whereas if $x \in$ $G \backslash G^{\dagger}$ and $j_{n}(x) \in K_{n}$ for infinitely many $n$, then $\liminf _{n \rightarrow \infty} h_{n}(x) \geq \Omega(g(x))$ because similarly to (4.3),

$$
\lim _{n \rightarrow \infty}\left[x_{n, j_{n}(x)}, x_{n, j_{n}(x)+1}\right]=\{x\} \quad \forall x \in G \backslash G^{\dagger}
$$

In summary, therefore,

$$
\begin{equation*}
\liminf _{n \rightarrow \infty} h_{n}(x) \geq \Omega\left(g(x) \mathbf{1}_{G}(x)\right) \text { for } \lambda \text {-almost every } x \in \mathbb{R} . \tag{4.20}
\end{equation*}
$$

Note that $j_{n}(z)-j_{n}(y)=\#\left\{1 \leq j \leq n: x_{n, j} \in I\right\}$. Consequently, (4.19), Proposition 4.3 with $\int_{\mathbb{R}} \Omega\left(f_{\mathrm{A}}^{\prime}\right) \mathrm{d} \lambda>0$, and Fatou's lemma applied to (4.20), together yield

$$
\liminf _{n \rightarrow \infty} \frac{\#\left\{1 \leq j \leq n: x_{n, j} \in I\right\}}{n} \geq \frac{\int_{I} \Omega\left(g \mathbf{1}_{G}\right) \mathrm{d} \lambda-\frac{a}{2}}{\int_{\mathbb{R}} \Omega\left(f_{\mathrm{A}}^{\prime}\right) \mathrm{d} \lambda}
$$

Recall that $f_{\mathrm{A}}^{\prime}=0$ on $\mathbb{R} \backslash G$; hence,

$$
\begin{aligned}
\int_{I} \Omega\left(g \mathbf{1}_{G}\right) \mathrm{d} \lambda & =\int_{I} \Omega\left(f_{\mathrm{A}}^{\prime}\right) \mathrm{d} \lambda+\int_{I \cap G}\left(\Omega\left(g \mathbf{1}_{G}\right)-\Omega\left(f_{\mathrm{A}}^{\prime}\right)\right) \mathrm{d} \lambda \\
& \geq \int_{I} \Omega\left(f_{\mathrm{A}}^{\prime}\right) \mathrm{d} \lambda-\frac{1}{2} \int_{I}\left|f_{\mathrm{A}}^{\prime}-g\right| \mathrm{d} \lambda \geq \int_{I} \Omega\left(f_{\mathrm{A}}^{\prime}\right) \mathrm{d} \lambda-\frac{a}{2}
\end{aligned}
$$

and consequently

$$
\liminf _{n \rightarrow \infty} \frac{\#\left\{1 \leq j \leq n: x_{n, j} \in I\right\}}{n} \geq \mu_{1}^{*}(I)-\frac{a}{\int_{\mathbb{R}} \Omega\left(f_{\mathrm{A}}^{\prime}\right) \mathrm{d} \lambda}
$$

Since the number $a>0$ and the interval $I \subset \mathbb{R}$ have been arbitrary, and since $\mu_{1}^{*}(\mathbb{R})=1$,

$$
\lim _{n \rightarrow \infty} \frac{\#\left\{1 \leq j \leq n: x_{n, j} \in I\right\}}{n}=\mu_{1}^{*}(I) \quad \forall I \subset \mathbb{R}, I \text { an interval, }
$$

i.e. (4.16) holds as claimed.

Note that Theorem 4.5 in particular asserts that if $\mu \in \mathcal{P}$ is non-singular and $\left(\delta_{x_{n}}^{p_{n}}\right)$, with $x_{n} \in \Xi_{n}$ and $p_{n} \in \Pi_{n}$ for every $n \in \mathbb{N}$, is any sequence of best $d_{\epsilon}{ }^{-}$ approximations of $\mu$, then the sequence $\left(\delta_{x_{n}}^{u_{n}}\right)$, obtained by "forgetting" the optimal weights and instead assigning equal weight $1 / n$ to each atom, converges weakly to $\mu_{\epsilon}^{*}$. It seems rather remarkable that $\left(\delta_{x_{n}}^{u_{n}}\right)$ always converges, and to a limit that is independent of $\left(x_{n}\right)$. By contrast, simple examples show that $\left(\delta_{x_{n}}^{u_{n}}\right)$ may diverge if $\mu$ is singular; cf. [19].

Example 7 Let $\mu=\exp (a)$ with $a>0$. With $\ell_{\bullet}^{\bullet, n}$ and the (unique) best $d_{\epsilon}$ approximation $\delta_{x_{n}}^{p_{n}}$ of $\mu$ found in Example 1, it is readily confirmed that for any $n \in \mathbb{N}$ and $x \in \mathbb{R}^{+}$the number $\#\left\{1 \leq j \leq n: x_{n, j} \leq x\right\}$ equals the largest integer not larger than

$$
n-\frac{\epsilon}{2 a \ell_{\bullet}^{\bullet, n}} \log \left(1+\left(\frac{e^{-a\left(x+\ell_{\bullet}^{\bullet, n} / \epsilon\right)}}{\ell_{\bullet}^{\bullet, n}}-1\right) \tanh \frac{a \ell_{\bullet}^{\bullet, n}}{\epsilon}\right)
$$

From this, a straightforward calculation utilizing (2.7) yields

$$
\lim _{n \rightarrow \infty} \frac{\#\left\{1 \leq j \leq n: x_{n, j} \leq x\right\}}{n}=1-\frac{\log \left(1+a e^{-a x} / \epsilon\right)}{\log (1+a / \epsilon)} \quad \forall x \in \mathbb{R}^{+}
$$

Thus, the asymptotic point density of $\left(x_{n}\right)$ is

$$
\frac{a^{2}}{\log (1+a / \epsilon)} \cdot \frac{1}{a+\epsilon e^{a x}}=\frac{\Omega\left(a e^{-a x} / \epsilon\right)}{\int_{\mathbb{R}} \Omega\left(a e^{-a y} / \epsilon\right) \mathrm{d} y}=\frac{\mathrm{d} \mu_{\epsilon}^{*}}{\mathrm{~d} \lambda}(x) \quad \forall x \in \mathbb{R}^{+}
$$

in perfect agreement with Theorem 4.5. Note that unlike for best $d_{\mathrm{w}}$-approximations [18, Thm. 7.5], this asymptotic point distribution is not exponential; see also Fig. 3.


Fig. 3 Comparing the standard exponential (left) and standard normal densities (broken curves) to the asymptotic point densities of their respective best $d_{1}$-approximations (solid curves); see Example 7

For another simple example, let $\mu$ be a normal distribution with mean 0 and variance $\sigma^{2}>0$. While no explicit formula is available for the (unique) best $d_{\epsilon}$-approximation $\delta_{x_{n}}^{p_{n}}$ of $\mu$ in general, Theorem 4.5 yields

$$
\frac{\mathrm{d} \mu_{\epsilon}^{*}}{\mathrm{~d} \lambda}(x)=-\frac{1}{\mathrm{Li}_{1 / 2}\left(-\frac{1}{\epsilon \sqrt{2 \pi \sigma^{2}}}\right)} \cdot \frac{1}{\sqrt{2 \pi \sigma^{2}}+2 \pi \epsilon \sigma^{2} e^{x^{2} /\left(2 \sigma^{2}\right)}} \quad \forall x \in \mathbb{R}
$$

as the asymptotic point density of $\left(x_{n}\right)$. Again, this asymptotic point distribution is not normal, unlike its $d_{\mathrm{w}}$-counterpart.

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