# MODIFIED BDF2 SCHEMES FOR SUBDIFFUSION MODELS WITH A SINGULAR SOURCE TERM

MINGHUA CHEN, JIANKANG SHI\*, AND ZHI ZHOU

Abstract. The aim of this paper is to study the time stepping scheme for approximately solving the subdiffusion equation with a weakly singular source term. In this case, many popular time stepping schemes, including the correction of high-order BDF methods, may lose their high-order accuracy. To fill in this gap, in this paper, we develop a novel time stepping scheme, where the source term is regularized by using an m-fold integral-derivative and the equation is discretized by using a modified BDF2 convolution quadrature. We prove that the proposed time stepping scheme is second-order, even if the source term is nonsmooth in time and incompatible with the initial data. Numerical results are presented to support the theoretical results. The proposed approach is applicable for stochastic subdiffusion equation.

Key words. Subdiffusion, modified BDF2 schemes, singular source term, error estimate.

#### 1. Introduction

For anomalous, non-Brownian diffusion, a mean squared displacement often follows the following power-law

$$\langle x^2(t)\rangle \simeq K_{\alpha}t^{\alpha}.$$

Prominent examples for subdiffusion include the classical charge carrier transport in amorphous semiconductors, tracer diffusion in subsurface aquifers, porous systems, dynamics of a bead in a polymeric network, or the motion of passive tracers in living biological cells [22, 23]. Subdiffusion of this type is characterised by a long-tailed waiting time probability density function  $\psi(t) \simeq t^{-1-\alpha}$ , corresponding to the time-fractional diffusion equation with and without an external force field [23, Eq. (88)]

$$\partial_t u(x,t) - \partial_t^{1-\alpha} A u(x,t) = f(x,t), \ 0 < \alpha < 1.$$

Here f is a given source function, and the operator  $A = \Delta$  denotes Laplacian on a polyhedral domain  $\Omega \subset \mathbb{R}^d$  (d = 1, 2, 3) with a homogenous Dirichlet boundary condition. The fractional derivative is taken in the Riemann-Liouville sense, that is,  $\partial_t^{1-\alpha}f = \partial_t J^{\alpha}f$  with the fractional integration operator

$$J^{\alpha}f(t) = \frac{1}{\Gamma(\alpha)} \int_0^t (t-\tau)^{\alpha-1} f(\tau) d\tau = \frac{1}{\Gamma(\alpha)} t^{\alpha-1} * f(t),$$

and \* denotes the Laplace convolution:  $(f * g)(t) = \int_0^t f(t - \tau)g(\tau)d\tau$ .

Since the Riemann-Liouvile fractional derivative and the Caputo fractional derivative can be written in the form [26, p. 76]

$$\partial_t^{\alpha} u(x,t) = {^C\!D}_t^{\alpha} u(x,t) + \frac{1}{\Gamma(1-\alpha)} t^{-\alpha} u(x,0),$$

Received by the editors on November 17, 2024 and, accepted on May 23, 2025. 2000 Mathematics Subject Classification. 26A33, 65M55, 65T50.

<sup>\*</sup>Corresponding author.

which implies that the equivalent form of  $(\spadesuit)$  can be rewritten as

$$(\heartsuit) \qquad \partial_t u(x,t) - {^CD_t}^{1-\alpha} Au(x,t) = f(x,t) + \frac{Au(x,0)}{\Gamma(\alpha)} t^{-(1-\alpha)}, \ 0 < \alpha < 1$$

with the Caputo fractional derivative

$${}^{C}D_{t}^{\alpha}u(t) = \frac{1}{\Gamma(1-\alpha)} \int_{0}^{t} (t-s)^{-\alpha}u'(s)ds, \ 0 < t \le T.$$

Applying the fractional integration operator  $J^{1-\alpha}$  to both sides of  $(\spadesuit)$ , we obtain the equivalent form of  $(\spadesuit)$  as, see [21, Eq. (1.6)] or [31, Eq. (2.3)], namely,

$${}^{C}D_{t}^{\alpha}u(x,t)-Au(x,t)=\frac{1}{\Gamma(1-\alpha)}t^{-\alpha}*f(x,t),\ 0<\alpha<1.$$

As another example, the fractal mobile/immobile models for solute transport associated with power law decay PDF describing random waiting times in the immobile zone, lead to the following models [29, Eq. (15)]

$$(\diamondsuit) \qquad \partial_t u(x,t) + {^CD_t^{\alpha}} u(x,t) - Au(x,t) = -\frac{1}{\Gamma(1-\alpha)} t^{-\alpha} u(x,0), \ 0 < \alpha < 1.$$

Note that the right hand side in the aforementioned PDE models  $(\spadesuit)$ - $(\diamondsuit)$  might be nonsmooth in the time variable. In this paper, we consider the subdiffusion model with weakly singular source term:

(1) 
$${}^{C}D_{t}^{\alpha}u(x,t) - Au(x,t) = g(x,t) := t^{\mu} \circ f(x,t)$$

with the initial condition  $u(x,0) = u_0(x) := v$ , and the homogeneous Dirichlet boundary conditions. The symbol  $\circ$  can be either the convolution \* or the product, and  $\mu$  is a parameter such that

$$\mu > -1$$
 if  $\circ$  denotes convolution, and  $\mu > -\alpha$  if  $\circ$  denotes product.

The well-posedness could be proved using the separation of variables and Mittag–Leffler functions, see e.g. [27, Eq. (2.11)].

Note that many existing time stepping schemes may lose their high-order accuracy when the source term is nonsmooth in the time variable. As an example, it was reported in [11, Section 4.1] that the convolution quadrature generated by k step BDF method (with initial correction) converges with order  $O(\tau^{1+\mu})$ , provided that the source term behaves like  $t^{\mu}$ ,  $\mu > 0$ , see Lemma 3.2 in [35], also see Table 1. The aim of this paper is to fill in this gap.

It is well-known that the smoothness of all the data of (1) (e.g., f=0) does not imply the smoothness of the solution u which has an initial layer at  $t\to 0^+$  (i.e., unbounded near t=0) [26, 27, 33]. There are already two predominant discretization techniques in time direction to restore the desired convergence rate for subdiffusion under appropriate regularity source function. The first type is that the nonuniform time meshes/graded meshes are employed to compensate/capture the singularity of the continuous solution near t=0 under the appropriate regularity source function and initial data, see [3, 15, 17, 20, 25, 24, 33]. See also spectral method with specially designed basis functions [4, 8, 38]. The second type is based on correction of high-order BDFk or  $L_k$  approximation, and then the desired high-order convergence rates can be restored even for nonsmooth initial data [5, 19, 18, 9, 11, 35]. For fractional ODEs, one idea is to use starting quadrature weights to correct the fractional integrals [18] (or fractional substantial calculus [1])

$$J^{\alpha}g(t) = \frac{1}{\Gamma(\alpha)} \int_0^t (t-\tau)^{\alpha-1} g(\tau) d\tau \text{ with } g(t) = t^{\mu} f(t), \ \mu > -1,$$

where the algorithms rely on expanding the solution into power series of t. For fractional PDEs, a common practice is to split the source term into

$$g(t) = g(0) + \sum_{l=1}^{k-1} \frac{t^l}{l!} \partial_t^l g(0) + \frac{t^{k-1}}{(k-1)!} * \partial_t^k g.$$

Then approximating g(0) by  $\partial_{\tau}J^{1}g(0)$  may a modified BDF2 scheme with correction in the first step [5]. The correction of high-order BDFk or  $L_{k}$  convolution quadrature is well developed in [11, 32, 37] when the source term is sufficiently smooth in the time variable. How to deal with a more general source term, which might be nonsmooth in the time variable, is still scarcely discussed in the literature. We recommend interested readers to refer to the concise overview [10] and the monograph [13] for a comprehensive understanding of the topic.

In this paper, we develop a novel second-order time stepping scheme (IDm-BDF2) for solving the subdiffusion (1) with a weakly singular source term, where the low regular source term is regularized by using an m-fold integral-derivative (IDm) and the equation is discretized by using a modified BDF2 convolution quadrature. We prove that the proposed time stepping scheme is second-order, even if the source term is nonsmooth in time and incompatible with the initial data. Numerical results are presented to support the theoretical results. Note that in [39] second-order time-stepping schemes are proposed for nonsmooth source data. In comparison with [39], the current paper offers significant contributions and improvements in the following aspects:

(i) Theoretically, the argument in [39] highly relies on the assumption on the source term g(x,t) that (cf. Assumption 1 in [39])

$$\|\widehat{g}(s)\|_{L^2(\Omega)} \le c|s|^{-\mu-1} \text{ with } -1 < \mu < 0,$$

for all  $s \in \Sigma_{\theta} = \{s \in \mathbb{C} \setminus \{0\} : |\arg s| < \theta\}$ . This assumption is indeed restrictive and only covers the case where the singular source function is given by  $g(x,t) = t^{\mu}f(x)$ , with f being time-independent. To address this limitation, in our work, we propose a refined error analysis that extends to a more general set of problem data, specifically  $t^{\mu} * f(t)$  and  $t^{\mu}f(t)$ . This broader analysis allows for a wider range of source functions and enhances the applicability of our approach.

(ii) The numerical scheme employed in our work differs significantly from the one presented in [39]. In [39], the solution u(t) is approximated using  $\partial_{\tau}J^{1}u(t)$ . To facilitate the approximation, an auxiliary function  $U(t)=J^{1}u(t)$  is introduced, which requires solving a time stepping scheme to address an evolution problem derived by integrating the original subdiffusion equation. Once an approximate solution for U(t) is obtained, a numerical differentiation method is applied to derive the numerical approximation for u(t). In contrast, in our current paper, we adopt a more direct approach. We directly solve the numerical approximation for u(t) using the approximate source  $g(t) \approx \partial_{\tau}J^{1}g$ . This method is more efficient and practical, eliminating the need for an intermediate auxiliary function.

Therefore, the current paper represents a significant improvement over the results presented in [39], both in terms of theoretical advancements and numerical aspects.

The paper is organized as follows. In Section 2, we introduce the development of the IDm-BDF2 scheme for model (1). In Section 3 and 4, based on operational calculus, the detailed convergence analysis of IDm-BDF2 is provided, respectively, for general source function f(x,t) and certain form  $t^{\mu}f(x)$ . Then the desired results

with the low regularity source term  $t^{\mu} \circ f(x,t)$  are obtained in Section 5. To show the effectiveness of the presented schemes, the results of numerical experiments are reported in Section 6. Finally, we conclude the paper with some remarks in the last section.

## 2. IDm-BDF2 Method

In this section, we first provide IDm-BDF2 method for solving subdiffusion (1) if the source term g(x,t) possess the mild regularity. Let V(t) = u(t) - v with V(0) = 0. Then the model (1) can be rewritten as

(2) 
$$\partial_t^{\alpha} V(t) - AV(t) = Av + g(t), \quad 0 < t \le T.$$

From [19] and [34], we know that the operator A satisfies the following resolvent estimate

$$||(z-A)^{-1}|| \le c_{\phi}|z|^{-1} \quad \forall z \in \Sigma_{\phi}$$

for all  $\phi \in (\pi/2, \pi)$ , where  $\Sigma_{\theta} := \{z \in \mathbb{C} \setminus \{0\} : |\arg z| < \theta\}$  is a sector of the complex plane  $\mathbb{C}$ . Hence,  $z^{\alpha} \in \Sigma_{\theta'}$  with  $\theta' = \alpha\theta < \theta < \pi$  for all  $z \in \Sigma_{\theta}$ . Then, there exists a positive constant c such that

(3) 
$$||(z^{\alpha} - A)^{-1}|| \le c|z|^{-\alpha} \quad \forall z \in \Sigma_{\theta}.$$

**2.1. Discretization schemes.** Let  $G(t) = J^1g(t)$  and  $G(t) = J^2g(t)$ . By first fundamental theorem of calculus, we may rewrite (2) as

(4) ID1 Method: 
$$\partial_t^{\alpha} V(t) - AV(t) = \partial_t (tAv + G(t)), \quad 0 < t \le T,$$

(5) ID2 Method: 
$$\partial_t^{\alpha} V(t) - AV(t) = \partial_t^2 \left( \frac{t^2}{2} Av + \mathcal{G}(t) \right), \quad 0 < t \le T.$$

Let  $t_n = n\tau, n = 0, 1, ..., N$ , be a uniform partition of the time interval [0, T] with the step size  $\tau = \frac{T}{N}$ , and let  $u^n$  denote the approximation of u(t) and  $g^n = g(t_n)$ . The convolution quadrature generated by BDF2 approximates the Riemann-Liouville fractional derivative  $\partial_t^{\alpha} \varphi(t_n)$  by

(6) 
$$\partial_{\tau}^{\alpha} \varphi^{n} := \frac{1}{\tau^{\alpha}} \sum_{j=0}^{n} \omega_{j} \varphi^{n-j}$$

with  $\varphi^n = \varphi(t_n)$ . Here the weights  $\omega_i$  are the coefficients in the series expansion

(7) 
$$\delta_{\tau}^{\alpha}(\xi) = \frac{1}{\tau^{\alpha}} \sum_{j=0}^{\infty} \omega_{j} \xi^{j} \quad \text{with} \quad \delta_{\tau}(\xi) := \frac{1}{\tau} \left( \frac{3}{2} - 2\xi + \frac{1}{2} \xi^{2} \right).$$

Then IDm-BDF2 method for (4) and (5) are, respectively, designed by

(8) ID1 – BDF2 Method : 
$$\partial_{\tau}^{\alpha} V^n - AV^n = \partial_{\tau} (t_n Av + G^n)$$
.

(9) 
$$ID2 - BDF2 \text{ Method} : \partial_{\tau}^{\alpha} V^n - AV^n = \partial_{\tau}^2 \left( \frac{t_n^2}{2} Av + \mathcal{G}^n \right).$$

Remark 2.1. In the time semidiscrete approximation (8) and (9), we require  $v \in \mathcal{D}(A)$ , i.e., the initial data v is reasonably smooth. However one can use the schemes (8) and (9) to prove the error estimates with the nonsmooth data  $v \in L^2(\Omega)$ , see Theorems 5.2 and 5.3. Here, we mainly focus on the time semidiscrete approximation (8) and (9), since the spatial discretization is well understood. For example, we choose  $v_h = R_h v$  if  $v \in \mathcal{D}(A)$  and  $v_h = P_h v$  if  $v \in L^2(\Omega)$  following [34, 36].

**2.2. Solution representation for** (4) **and** (5). Taking the Laplace transform in both sides of (4), it leads to

$$\widehat{V}(z) = (z^{\alpha} - A)^{-1} \left( z^{-1} A v + z \widehat{G}(z) \right).$$

By the inverse Laplace transform, there exists [11]

(10) 
$$V(t) = \frac{1}{2\pi i} \int_{\Gamma_{0,r}} e^{zt} (z^{\alpha} - A)^{-1} \left( z^{-1} A v + z \widehat{G}(z) \right) dz$$

with

(11) 
$$\Gamma_{\theta,\kappa} = \{z \in \mathbb{C} : |z| = \kappa, |\arg z| \le \theta\} \cup \{z \in \mathbb{C} : z = re^{\pm i\theta}, r \ge \kappa\}$$
 and  $\theta \in (\pi/2, \pi), \kappa > 0$ .

Similarly, applying the Laplace transform in both sides of (5), it yields

$$\widehat{V}(z) = (z^{\alpha} - A)^{-1} \left( z^{-1} A v + z^2 \widehat{\mathcal{G}}(z) \right).$$

By the inverse Laplace transform, we obtain

(12) 
$$V(t) = \frac{1}{2\pi i} \int_{\Gamma_{\alpha, r}} e^{zt} (z^{\alpha} - A)^{-1} \left( z^{-1} A v + z^2 \widehat{\mathcal{G}}(z) \right) dz.$$

**2.3.** Discrete solution representation for (8) and (9). Given a sequence  $(\kappa_n)_0^{\infty}$  and take  $\widetilde{\kappa}(\zeta) = \sum_{n=0}^{\infty} \kappa_n \zeta^n$  to be its generating power series. Let us first introduce the elementary identities

$$\gamma_l(\xi) = \sum_{n=1}^{\infty} n^l \xi^n = \left(\xi \frac{d}{d\xi}\right)^l \frac{1}{1-\xi}, \ l \ge 1 \text{ and } \gamma_0(\xi) = \sum_{n=1}^{\infty} \xi^n = \frac{\xi}{1-\xi}.$$

**Lemma 2.1.** Let  $\delta_{\tau}(\xi)$  be given in (7) and  $\gamma_1(\xi) = \sum_{n=1}^{\infty} n \xi^n$ ,  $G(t) = J^1 g(t)$ . Then the discrete solution of (8) is represented by

$$V^n = \frac{1}{2\pi i} \int_{\Gamma^\tau_{A_r}} e^{zt_n} (\delta^\alpha_\tau(e^{-z\tau}) - A)^{-1} \delta_\tau(e^{-z\tau}) \tau \left( \gamma_1(e^{-z\tau}) \tau A v + \widetilde{G}(e^{-z\tau}) \right) dz$$

with  $\Gamma_{\theta,\kappa}^{\tau} = \{ z \in \Gamma_{\theta,\kappa} : |\Im z| \le \pi/\tau \}.$ 

*Proof.* Multiplying the (8) by  $\xi^n$  and summing over n with  $V^0 = 0$ , we obtain

$$\sum_{n=1}^{\infty} \partial_{\tau}^{\alpha} V^n \xi^n - \sum_{n=1}^{\infty} A V^n \xi^n = \sum_{n=1}^{\infty} \partial_{\tau} (t_n A v + G^n) \xi^n.$$

From (6) and (7), we have

$$\begin{split} \sum_{n=1}^{\infty} \partial_{\tau}^{\alpha} V^n \xi^n &= \sum_{n=1}^{\infty} \frac{1}{\tau^{\alpha}} \sum_{j=0}^{n} \omega_j V^{n-j} \xi^n = \sum_{n=0}^{\infty} \frac{1}{\tau^{\alpha}} \sum_{j=0}^{n} \omega_j V^{n-j} \xi^n = \sum_{j=0}^{\infty} \frac{1}{\tau^{\alpha}} \sum_{n=j}^{\infty} \omega_j V^{n-j} \xi^n \\ &= \sum_{j=0}^{\infty} \frac{1}{\tau^{\alpha}} \sum_{n=0}^{\infty} \omega_j V^n \xi^{n+j} = \frac{1}{\tau^{\alpha}} \sum_{j=0}^{\infty} \omega_j \xi^j \sum_{n=0}^{\infty} V^n \xi^n = \delta_{\tau}^{\alpha}(\xi) \widetilde{V}(\xi). \end{split}$$

Similarly, one has

$$\sum_{n=1}^{\infty} \partial_{\tau} t_n A v \xi^n = \delta_{\tau}(\xi) \gamma_1(\xi) \tau A v, \quad \sum_{n=1}^{\infty} \partial_{\tau} G^n \xi^n = \delta_{\tau}(\xi) \widetilde{G}(\xi)$$

with  $\gamma_1(\xi) = \frac{\xi}{(1-\xi)^2}$ . It leads to

(13) 
$$\widetilde{V}(\xi) = \left(\delta_{\tau}^{\alpha}(\xi) - A\right)^{-1} \delta_{\tau}(\xi) \left(\gamma_{1}(\xi)\tau A v + \widetilde{G}(\xi)\right).$$

According to Cauchy's integral formula, and the change of variables  $\xi = e^{-z\tau}$ , and Cauchy's theorem, one has [11]

$$V^n = \frac{\tau}{2\pi i} \int_{\Gamma_{a}^{\tau}} e^{zt_n} \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) \left( \gamma_1(e^{-z\tau}) \tau A v + \widetilde{G}(e^{-z\tau}) \right) dz$$

with  $\Gamma_{\theta,\kappa}^{\tau} = \{z \in \Gamma_{\theta,\kappa} : |\Im z| \le \pi/\tau\}$ . The proof is completed.

**Lemma 2.2.** Let  $\delta_{\tau}(\xi)$  be given in (7) and  $\gamma_2(\xi) = \sum_{n=1}^{\infty} n^2 \xi^n$ ,  $\mathcal{G}(t) = J^2 g(t)$ . Then the discrete solution of (9) is represented by

$$V^n = \frac{\tau}{2\pi i} \int_{\Gamma^\tau_{\theta,\kappa}} e^{zt_n} \left( \delta^\alpha_\tau(e^{-z\tau}) - A \right)^{-1} \delta^2_\tau(e^{-z\tau}) \left( \frac{\gamma_2(e^{-z\tau})}{2} \tau^2 A v + \widetilde{\mathcal{G}}(e^{-z\tau}) \right) dz$$

with  $\Gamma_{\theta,\kappa}^{\tau} = \{z \in \Gamma_{\theta,\kappa} : |\Im z| \le \pi/\tau\}.$ 

*Proof.* Multiplying the (9) by  $\xi^n$  and summing over n with  $V^0 = 0$ , we obtain

$$\sum_{n=1}^{\infty} \partial_{\tau}^{\alpha} V^n \xi^n - \sum_{n=1}^{\infty} A V^n \xi^n = \sum_{n=1}^{\infty} \partial_{\tau}^2 \left( \frac{t_n^2}{2} A v + \mathcal{G}^n \right) \xi^n.$$

The similar arguments can be performed as Lemma 2.1, it yields

$$\sum_{n=1}^{\infty} \partial_{\tau}^{\alpha} V^n \xi^n = \delta_{\tau}^{\alpha}(\xi) \widetilde{V}(\xi), \quad \sum_{n=1}^{\infty} \partial_{\tau}^2 t_n^2 A v \xi^n = \delta_{\tau}^2(\xi) \gamma_2(\xi) \tau^2 A v,$$

$$\sum_{n=1}^{\infty} \partial_{\tau}^{2} \mathcal{G}^{n} \xi^{n} = \delta_{\tau}^{2}(\xi) \widetilde{\mathcal{G}}(\xi), \quad \gamma_{2}(\xi) = \frac{\xi + \xi^{2}}{(1 - \xi)^{3}},$$

and

$$(15) \qquad \widetilde{V}(\xi) = \left(\delta_{\tau}^{\alpha}(\xi) - A\right)^{-1} \delta_{\tau}^{2}(\xi) \left(\frac{\gamma_{2}(\xi)}{2} \tau^{2} A v + \widetilde{\mathcal{G}}(\xi)\right).$$

Using Cauchy's integral formula, and the change of variables  $\xi = e^{-z\tau}$ , and Cauchy's theorem, one has

$$(16) \quad V^n = \frac{\tau}{2\pi i} \int_{\Gamma_{\theta}^{\tau}} e^{zt_n} \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}^2(e^{-z\tau}) \left( \frac{\gamma_2(\xi)}{2} \tau^2 A v + \widetilde{\mathcal{G}}(e^{-z\tau}) \right) dz$$

with 
$$\Gamma_{\theta,\kappa}^{\tau} = \{z \in \Gamma_{\theta,\kappa} : |\Im z| \le \pi/\tau\}$$
. The proof is completed.

# 3. Convergence analysis: General source function g(x,t)

In this section, we provide the detailed convergence analysis of ID1-BDF2 in (8) approximation for the subdiffusion (4), and ID2-BDF2 can be similarly augmented.

## **3.1.** A few technical lemmas. First, we give some lemmas that will be used.

**Lemma 3.1.** [11] Let  $\delta_{\tau}(\xi)$  be given in (7). Then there exist the positive constants  $c_1, c_2, c$  and  $\theta \in (\pi/2, \theta_{\varepsilon})$  with  $\theta_{\varepsilon} \in (\pi/2, \pi), \forall \varepsilon > 0$  such that

$$c_1|z| \le |\delta_{\tau}(e^{-z\tau})| \le c_2|z|, \quad |\delta_{\tau}(e^{-z\tau}) - z| \le c\tau^2|z|^3,$$
$$|\delta_{\tau}^{\alpha}(e^{-z\tau}) - z^{\alpha}| \le c\tau^2|z|^{2+\alpha}, \quad \delta_{\tau}(e^{-z\tau}) \in \Sigma_{\pi/2+\varepsilon} \ \forall z \in \Gamma_{\theta,\kappa}^{\tau}.$$

**Lemma 3.2.** Let  $\delta_{\tau}(\xi)$  be given in (7) and  $\gamma_{l}(\xi) = \sum_{n=1}^{\infty} n^{l} \xi^{n}$  with l = 0, 1, 2. Then there exists a positive constants c such that

$$\left| \frac{\gamma_l(e^{-z\tau})}{l!} \tau^{l+1} - z^{-l-1} \right| \le c\tau^{l+1}, \quad \forall z \in \Gamma_{\theta,\kappa}^{\tau},$$

where  $\theta \in (\pi/2, \pi)$  is sufficiently close to  $\pi/2$ .

*Proof.* The arguments can be performed in [32] for l=1,2. For l=0, we have  $\frac{e^{-z\tau}}{1-e^{-z\tau}}\tau - \frac{1}{1-e^{-z\tau}}\tau = -\tau$ . On the other hand, using

$$\frac{1}{1 - e^{-z\tau}}\tau - z^{-1} = \frac{z - (1 - e^{-z\tau})\tau^{-1}}{(1 - e^{-z\tau})\tau^{-1}z},$$

and Lemma 3.1, it yields  $|1 - e^{-z\tau}| \ge c_1 |z|\tau$  and

(17) 
$$\left| (1 - e^{-z\tau})\tau^{-1}z \right| \ge c|z|^2 \ \forall z \in \Gamma^{\tau}_{\theta,\kappa}.$$

Since

$$\begin{aligned} \left| z - \left( 1 - e^{-z\tau} \right) \tau^{-1} \right| &= \left| z - \left( 1 - \sum_{j=0}^{\infty} \frac{(-z\tau)^j}{j!} \right) \tau^{-1} \right| = \left| z - \left( -\sum_{j=1}^{\infty} \frac{(-z\tau)^j}{j!} \right) \tau^{-1} \right| \\ &= \left| z - z \sum_{j=0}^{\infty} \frac{(-z\tau)^j}{(j+1)!} \right| = \left| \tau z^2 \sum_{j=0}^{\infty} \frac{(-z\tau)^j}{(j+2)!} \right| \le c\tau |z|^2 \,. \end{aligned}$$

Thus we have

$$\left| \frac{1}{1 - e^{-z\tau}} \tau - z^{-1} \right| \le c\tau.$$

The proof is completed.

**Lemma 3.3.** Let  $\delta_{\tau}(\xi)$  be given in (7) and  $\gamma_{l}(\xi) = \sum_{n=1}^{\infty} n^{l} \xi^{n}$  with l = 0, 1, 2. Then there exists a positive constants c such that

(18) 
$$\left| \delta_{\tau}(e^{-z\tau}) \frac{\gamma_{l}(e^{-z\tau})}{l!} \tau^{l+1} - z^{-l} \right| \leq c\tau^{l+1} |z| + c\tau^{2} |z|^{2-l}, \quad \forall z \in \Gamma_{\theta,\kappa}^{\tau},$$

where  $\theta \in (\pi/2, \pi)$  is sufficiently close to  $\pi/2$ .

*Proof.* Let

$$\delta_{\tau}(e^{-z\tau})\frac{\gamma_{l}(e^{-z\tau})}{l!}\tau^{l+1} - z^{-l} = J_{1} + J_{2}$$

with

$$J_1 = \delta_{\tau}(e^{-z\tau}) \frac{\gamma_l(e^{-z\tau})}{l!} \tau^{l+1} - \delta_{\tau}(e^{-z\tau}) z^{-l-1} \quad \text{and} \quad J_2 = \delta_{\tau}(e^{-z\tau}) z^{-l-1} - z^{-l}.$$

According to Lemma 3.1 and 3.2, we have

$$|J_1| = \left| \delta_{\tau}(e^{-z\tau}) \left( \frac{\gamma_l(e^{-z\tau})}{l!} \tau^{l+1} - z^{-l-1} \right) \right| \le c_2 |z| c\tau^{l+1} \le c\tau^{l+1} |z|$$

and

$$|J_2| = |(\delta_\tau(e^{-z\tau}) - z) z^{-l-1}| \le c\tau^2 |z|^{2-l}.$$

By the triangle inequality, the desired result is obtained.

**Lemma 3.4.** Let  $\delta_{\tau}^{\alpha}(\xi)$  be given by (7) and  $\gamma_{l}(\xi) = \sum_{n=1}^{\infty} n^{l} \xi^{n}$  with l = 0, 1, 2. Then there exists a positive constants c such that

$$\left\| \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) \frac{\gamma_{l}(e^{-z\tau})}{l!} \tau^{l+1} - (z^{\alpha} - A)^{-1} z^{-l} \right\|$$

$$\leq c\tau^{l+1} |z|^{1-\alpha} + c\tau^{2} |z|^{2-l-\alpha}.$$

Proof. Let

$$\left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} \delta_{\tau}(e^{-z\tau}) \frac{\gamma_{l}(e^{-z\tau})}{l!} \tau^{l+1} - (z^{\alpha} - A)^{-1} z^{-l} = I + II$$

with

$$I = \left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} \left[\delta_{\tau}(e^{-z\tau}) \frac{\gamma_{l}(e^{-z\tau})}{l!} \tau^{l+1} - z^{-l}\right],$$
$$II = \left[\left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} - (z^{\alpha} - A)^{-1}\right] z^{-l}.$$

The resolvent estimate (3) and Lemma 3.1 imply directly

(19) 
$$\| \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \| \le c|z|^{-\alpha}.$$

From (19) and Lemma 3.3, we obtain

$$||I|| \le c\tau^{l+1} |z|^{1-\alpha} + c\tau^2 |z|^{2-l-\alpha}$$

Using Lemma 3.1, (19) and the identity (20)

$$\left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} - (z^{\alpha} - A)^{-1} = \left(z^{\alpha} - \delta_{\tau}^{\alpha}(e^{-z\tau})\right) \left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} (z^{\alpha} - A)^{-1},$$

we estimate II as following

$$||II|| \le c\tau^2 |z|^{2+\alpha} c|z|^{-\alpha} c|z|^{-\alpha} |z|^{-l} \le c\tau^2 |z|^{2-l-\alpha}.$$

By the triangle inequality, the desired result is obtained.

**Lemma 3.5.** Let  $\delta_{\tau}^{\alpha}(\xi)$  be given by (7) and  $\gamma_1(\xi) = \sum_{n=1}^{\infty} n\xi^n = \left(\xi \frac{d}{d\xi}\right) \frac{1}{1-\xi} = \frac{\xi}{(1-\xi)^2}$ . Then there exists a positive constants c such that

$$\left\| \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) \gamma_{1}(e^{-z\tau}) \tau^{2} A - (z^{\alpha} - A)^{-1} z^{-1} A \right\| \leq c\tau^{2} |z|.$$

*Proof.* Using identical  $(z^{\alpha}-A)^{-1}z^{-l}A=-z^{-1}+(z^{\alpha}-A)^{-1}z^{\alpha}z^{-1}$  and

$$\left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} \delta_{\tau}(e^{-z\tau})A = -\delta_{\tau}(e^{-z\tau}) + \left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} \delta_{\tau}^{\alpha}(e^{-z\tau})\delta_{\tau}(e^{-z\tau})A,$$

 $\left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} \delta_{\tau}(e^{-z\tau}) \gamma_{1}(e^{-z\tau}) \tau^{2} A - (z^{\alpha} - A)^{-1} z^{-1} A = J_{1} + J_{2} + J_{3} + J_{4} + J_{5} +$ 

$$\begin{split} J_1 &= \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}^{\alpha}(e^{-z\tau}) \left( \delta_{\tau}(e^{-z\tau}) \gamma_1(e^{-z\tau}) \tau^{l+1} - z^{-1} \right), \\ J_2 &= \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - z^{\alpha} \right) z^{-1}, \\ J_3 &= \left( \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} - (z^{\alpha} - A)^{-1} \right) z^{\alpha - 1}, \quad J_4 = z^{-1} - \delta_{\tau}(e^{-z\tau}) \gamma_1(e^{-z\tau}) \tau^2. \end{split}$$

According to (19) and Lemmas 3.1, 3.3 with l = 1, we estimate  $J_1$ ,  $J_2$  and  $J_4$  as following

$$||J_1|| \le c|z|^{-\alpha}|z|^{\alpha}\tau^2|z| \le c\tau^2|z|,$$
  
$$||J_2|| \le c|z|^{-\alpha}\tau^2|z|^{2+\alpha}|z|^{-1} \le c\tau^2|z|, \quad ||J_4|| \le c\tau^2|z|.$$

From Lemma 3.1, (19) and the identity (20), we estimate  $J_3$  as following

$$||J_3|| \le c\tau^2 |z|^{2+\alpha} |z|^{-\alpha} |z|^{-\alpha} |z|^{-\alpha} |z|^{\alpha-1} \le c\tau^2 |z|.$$

By the triangle inequality, the desired result is obtained.

**3.2.** Error analysis for general source function g(x,t). From  $G(t) = J^1g(t)$ , the Taylor expansion of source function with the remainder term in integral form:

$$1 * g(t) = G(t) = G(0) + tG'(0) + \frac{t^2}{2}G''(0) + \frac{t^2}{2} * G'''(t)$$
$$= J^1 g(0) + tg(0) + \frac{t^2}{2}g'(0) + \frac{t^2}{2} * g''(t).$$

Then we obtain the following results with  $g^{(-1)}(0) = J^1g(0)$ .

**Lemma 3.6.** Let  $V(t_n)$  and  $V^n$  be the solutions of (4) and (8), respectively. Let v=0 and  $G(t):=\frac{t^l}{l!}g^{(l-1)}(0)$  with l=0,1,2. Then

(21) 
$$||V(t_n) - V^n|| \le \left(c\tau^{l+1}t_n^{\alpha-2} + c\tau^2t_n^{\alpha+l-3}\right) ||g^{(l-1)}(0)||.$$

*Proof.* Using (10) and (14), there exist

$$V(t_n) = \frac{1}{2\pi i} \int_{\Gamma_{a,r}} e^{zt_n} (z^{\alpha} - A)^{-1} \frac{1}{z^l} g^{(l-1)}(0) dz,$$

and

$$V^{n} = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa}^{\tau}} e^{zt_{n}} \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) \frac{\gamma_{l}(e^{-z\tau})}{l!} \tau^{l+1} g^{(l-1)}(0) dz,$$

where  $\theta \in (\pi/2, \pi)$  is sufficiently close to  $\pi/2$ , and  $\gamma_l(\xi) = \sum_{n=1}^{\infty} n^l \xi^n$ . Let

$$V(t_n) - V^n = J_1 + J_2$$

with

$$J_{1} = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa}^{\tau}} e^{zt_{n}} \left[ \frac{(z^{\alpha} - A)^{-1}}{z^{l}} - \left( \delta_{\tau}^{\alpha} (e^{-z\tau}) - A \right)^{-1} \delta_{\tau} (e^{-z\tau}) \frac{\gamma_{l}(e^{-z\tau})}{l!} \tau^{l+1} \right] g^{(l-1)}(0) dz,$$

and

$$J_2 = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa} \setminus \Gamma_{\theta,\kappa}^{\tau}} e^{zt_n} \left( z^{\alpha} - A \right)^{-1} \frac{1}{z^l} g^{(l-1)}(0) dz.$$

According to the triangle inequality, (3) and Lemma 3.4, one has

$$||J_1|| \le c \int_{\kappa}^{\frac{\pi}{\tau \sin \theta}} e^{rt_n \cos \theta} \left( \tau^{l+1} r^{1-\alpha} + \tau^2 r^{2-l-\alpha} \right) dr \, ||g^{(l-1)}(0)||$$

$$+ c \int_{-\theta}^{\theta} e^{\kappa t_n \cos \psi} \left( \tau^{l+1} \kappa^{2-\alpha} + \tau^2 \kappa^{3-l-\alpha} \right) d\psi \, ||g^{(l-1)}(0)||$$

$$\le \left( c \tau^{l+1} t_n^{\alpha-2} + c \tau^2 t_n^{\alpha+l-3} \right) \, ||g^{(l-1)}(0)|| \, ,$$

for the last inequality, we use

(22) 
$$\int_{\kappa}^{\frac{\pi}{\tau \sin \theta}} e^{rt_n \cos \theta} r^{2-l-\alpha} dr = t_n^{\alpha+l-3} \int_{t_n \kappa}^{\frac{t_n \pi}{\tau \sin \theta}} e^{s \cos \theta} s^{2-l-\alpha} ds \le c t_n^{\alpha+l-3},$$

$$\int_{-\theta}^{\theta} e^{\kappa t_n \cos \psi} \kappa^{3-l-\alpha} d\psi = t_n^{\alpha+l-3} \int_{-\theta}^{\theta} e^{\kappa t_n \cos \psi} (\kappa t_n)^{3-l-\alpha} d\psi \le c t_n^{\alpha+l-3}.$$

From (3), it yields

$$||J_{2}|| \leq c ||g^{(l-1)}(0)|| \int_{\frac{\pi}{\tau \sin \theta}}^{\infty} e^{rt_{n} \cos \theta} r^{-l-\alpha} dr$$

$$\leq c\tau^{2} ||g^{(l-1)}(0)|| \int_{\frac{\pi}{\tau \sin \theta}}^{\infty} e^{rt_{n} \cos \theta} r^{2-l-\alpha} dr$$

$$\leq c\tau^{2} t_{n}^{\alpha+l-3} ||g^{(l-1)}(0)||.$$

Here we using  $1 \leq (\frac{\sin \theta}{\pi})^2 \tau^2 r^2$  with  $r \geq \frac{\pi}{\tau \sin \theta}$ . The proof is completed.

**Lemma 3.7.** Let  $V(t_n)$  and  $V^n$  be the solutions of (4) and (8), respectively. Let  $v=0, G(t):=\frac{t^2}{2}*g''(t)$  and  $\int_0^t (t-s)^{\alpha-1}\|g''(s)\|ds<\infty$ . Then

$$||V(t_n) - V^n|| \le c\tau^2 \int_0^{t_n} (t_n - s)^{\alpha - 1} ||g''(s)|| ds.$$

*Proof.* By (10), we obtain

(23) 
$$V(t_n) = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa}} e^{zt_n} (z^{\alpha} - A)^{-1} z \widehat{G}(z) dz = (\mathscr{E}(t) * G(t))(t_n)$$
$$= \left(\mathscr{E}(t) * \left(\frac{t^2}{2} * g''(t)\right)\right) (t_n) = \left(\left(\mathscr{E}(t) * \frac{t^2}{2}\right) * g''(t)\right) (t_n)$$

with

(24) 
$$\mathscr{E}(t) = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa}} e^{zt} (z^{\alpha} - A)^{-1} z dz.$$

From (13), it yields

$$\begin{split} \widetilde{V}(\xi) &= \left(\delta_{\tau}^{\alpha}(\xi) - A\right)^{-1} \delta_{\tau}(\xi) \widetilde{G}(\xi) = \widetilde{\mathscr{E}_{\tau}}(\xi) \widetilde{G}(\xi) = \sum_{n=0}^{\infty} \mathscr{E}_{\tau}^{n} \xi^{n} \sum_{j=0}^{\infty} G^{j} \xi^{j} \\ &= \sum_{n=0}^{\infty} \sum_{j=0}^{\infty} \mathscr{E}_{\tau}^{n} G^{j} \xi^{n+j} = \sum_{j=0}^{\infty} \sum_{n=j}^{\infty} \mathscr{E}_{\tau}^{n-j} G^{j} \xi^{n} = \sum_{n=0}^{\infty} \sum_{j=0}^{n} \mathscr{E}_{\tau}^{n-j} G^{j} \xi^{n} = \sum_{n=0}^{\infty} V^{n} \xi^{n} \end{split}$$

with

$$V^n = \sum_{j=0}^n \mathscr{E}_\tau^{n-j} G^j := \sum_{j=0}^n \mathscr{E}_\tau^{n-j} G(t_j).$$

Here  $\sum_{n=0}^{\infty} \mathscr{E}_{\tau}^{n} \xi^{n} = \widetilde{\mathscr{E}_{\tau}}(\xi) = (\delta_{\tau}^{\alpha}(\xi) - A)^{-1} \delta_{\tau}(\xi)$ . From the Cauchy's integral formula and the change of variables  $\xi = e^{-z\tau}$ , we obtain the representation of the  $\mathscr{E}_{\tau}^{n}$  as following

$$\mathscr{E}_{\tau}^{n} = \frac{1}{2\pi i} \int_{|\xi| = \rho} \xi^{-n-1} \widetilde{\mathscr{E}_{\tau}}(\xi) d\xi = \frac{\tau}{2\pi i} \int_{\Gamma_{\theta,\kappa}^{\tau}} e^{zt_{n}} \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) dz,$$

where  $\theta \in (\pi/2, \pi)$  is sufficiently close to  $\pi/2$  and  $\kappa = t_n^{-1}$  in (11). According to (19), Lemma 3.1 and  $\tau t_n^{-1} = \frac{1}{n} \leq 1$ , there exists (25)

$$\|\mathscr{E}_{\tau}^{n}\| \le c\tau \left( \int_{\kappa}^{\frac{\pi}{\tau \sin \theta}} e^{rt_{n} \cos \theta} r^{1-\alpha} dr + \int_{-\theta}^{\theta} e^{\kappa t_{n} \cos \psi} \kappa^{2-\alpha} d\psi \right) \le c\tau t_{n}^{\alpha-2} \le ct_{n}^{\alpha-1}.$$

Let  $\mathscr{E}_{\tau}(t) = \sum_{n=0}^{\infty} \mathscr{E}_{\tau}^{n} \delta_{t_{n}}(t)$ , with  $\delta_{t_{n}}$  being the Dirac delta function at  $t_{n}$ . Then

(26) 
$$(\mathcal{E}_{\tau}(t) * G(t))(t_n) = \left(\sum_{j=0}^{\infty} \mathcal{E}_{\tau}^j \delta_{t_j}(t) * G(t)\right)(t_n)$$

$$= \sum_{j=0}^n \mathcal{E}_{\tau}^j G(t_n - t_j) = \sum_{j=0}^n \mathcal{E}_{\tau}^{n-j} G(t_j) = V^n.$$

Moreover, using the above equation, there exist

$$(\widetilde{\mathcal{E}_{\tau} * t^{l}})(\xi) = \sum_{n=0}^{\infty} \sum_{j=0}^{n} \mathcal{E}_{\tau}^{n-j} t_{j}^{l} \xi^{n} = \sum_{j=0}^{\infty} \sum_{n=j}^{\infty} \mathcal{E}_{\tau}^{n-j} t_{j}^{l} \xi^{n} = \sum_{j=0}^{\infty} \sum_{n=0}^{\infty} \mathcal{E}_{\tau}^{n} t_{j}^{l} \xi^{n+j}$$

$$= \sum_{n=0}^{\infty} \mathcal{E}_{\tau}^{n} \xi^{n} \sum_{j=0}^{\infty} t_{j}^{l} \xi^{j} = \widetilde{\mathcal{E}_{\tau}}(\xi) \tau^{l} \sum_{j=0}^{\infty} j^{l} \xi^{j} = \widetilde{\mathcal{E}_{\tau}}(\xi) \tau^{l} \gamma_{l}(\xi), \quad l = 1, 2.$$

In particular, we have

$$(\widetilde{\mathscr{E}_{\tau} * 1})(\xi) = \widetilde{\mathscr{E}_{\tau}}(\xi) \sum_{j=0}^{\infty} \xi^{j} = \widetilde{\mathscr{E}_{\tau}}(\xi) \frac{1}{1-\xi}, \ l = 0.$$

From (23), (26) and (21), we have the following estimate

(27) 
$$\left\| \left( (\mathscr{E}_{\tau} - \mathscr{E}) * \frac{t^{l}}{l!} \right) (t_{n}) \right\| \leq c\tau^{l+1} t_{n}^{\alpha-2} + c\tau^{2} t_{n}^{\alpha+l-3} \leq c\tau^{l} t_{n}^{\alpha-1} \quad l = 0, 1, 2.$$

Next, we prove the following inequality (28) for t > 0

(28) 
$$\left\| \left( (\mathscr{E}_{\tau} - \mathscr{E}) * \frac{t^2}{2} \right) (t) \right\| \le c\tau^2 t^{\alpha - 1}, \quad \forall t \in (t_{n-1}, t_n).$$

By Taylor series expansion of  $\mathscr{E}(t)$  at  $t=t_n$ , we get

$$\left(\mathscr{E} * \frac{t^2}{2}\right)(t) = \left(\mathscr{E} * \frac{t^2}{2}\right)(t_n) + (t - t_n)\left(\mathscr{E} * t\right)(t_n) + \frac{(t - t_n)^2}{2}\left(\mathscr{E} * 1\right)(t_n) + \frac{1}{2}\int_t^t (t - s)^2 \mathscr{E}(s)ds,$$

which also holds for  $(\mathscr{E}_{\tau} * t^2)(t)$ . Therefore, using (27), it yields

$$\left\| \left( (\mathscr{E}_{\tau} - \mathscr{E}) * \frac{t^{l}}{l!} \right) (t_{n}) \right\| \leq c\tau^{l+1} t_{n}^{\alpha - 2} + c\tau^{2} t_{n}^{\alpha + l - 3} \leq c\tau^{l} t_{n}^{\alpha - 1} \leq c\tau^{l} t^{\alpha - 1}.$$

According to (24), (3) and (22), one has

$$\|\mathscr{E}(t)\| \leq c \left( \int_{\kappa}^{\infty} e^{rt\cos\theta} r^{1-\alpha} dr + \int_{-\theta}^{\theta} e^{\kappa t\cos\psi} \kappa^{2-\alpha} d\psi \right) \leq ct^{\alpha-2}.$$

Moreover, we get

$$\left\|\int_{t_n}^t (t-s)^2 \mathscr{E}(s) ds\right\| \leq c \int_t^{t_n} (s-t)^2 s^{\alpha-2} ds \leq c \int_t^{t_n} (s-t) s^{\alpha-1} ds \leq c \tau^2 t^{\alpha-1}.$$

Using the definition of  $\mathscr{E}_{\tau}(t) = \sum_{n=0}^{\infty} \mathscr{E}_{\tau}^{n} \delta_{t_{n}}(t)$  in (26) and (25), we deduce

$$\left\| \int_{t_n}^t (t-s)^2 \mathscr{E}_{\tau}(s) ds \right\| \le (t_n - t)^2 \| \mathscr{E}_{\tau}^n \| \le c\tau^3 t_n^{\alpha - 2}$$
$$\le c\tau^2 t_n^{\alpha - 1} \le c\tau^2 t^{\alpha - 1}, \ \forall \ t \in (t_{n-1}, t_n).$$

By (27) and the above inequalities, it yields the inequality (28). The proof is completed.  $\hfill\Box$ 

**Theorem 3.1** (ID1-BDF2). Let  $V(t_n)$  and  $V^n$  be the solutions of (4) and (8), respectively. Let  $v \in L^2(\Omega)$ ,  $g \in C^1([0,T];L^2(\Omega))$  and  $\int_0^t (t-s)^{\alpha-1} \|g''(s)\| ds < \infty$ . Then the following error estimate holds for any  $t_n > 0$ :

$$||V^{n} - V(t_{n})|| \le c\tau^{2} \left( t_{n}^{-2} ||v|| + t_{n}^{\alpha-2} ||g(0)|| + t_{n}^{\alpha-1} ||g'(0)|| + \int_{0}^{t_{n}} (t_{n} - s)^{\alpha-1} ||g''(s)|| ds \right).$$

*Proof.* Subtracting (10) from (14), we obtain

$$V^n - V(t_n) = I_1 - I_2 + I_3$$

with

$$\begin{split} I_1 = & \frac{1}{2\pi i} \int_{\Gamma^{\tau}_{\theta,\kappa}} e^{zt_n} \left[ \left( \delta^{\alpha}_{\tau}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) \gamma_1(e^{-z\tau}) \tau^2 - (z^{\alpha} - A)^{-1} z^{-1} \right] Av dz, \\ I_2 = & \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa} \setminus \Gamma^{\tau}_{\theta,\kappa}} e^{zt_n} (z^{\alpha} - A)^{-1} z^{-1} Av dz, \\ I_3 = & \frac{\tau}{2\pi i} \int_{\Gamma^{\tau}_{\theta,\kappa}} e^{zt_n} \left( \delta^{\alpha}_{\tau}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) \widetilde{G}(e^{-z\tau}) dz \\ & - \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa}} e^{zt_n} (z^{\alpha} - A)^{-1} z \widehat{G}(z) dz. \end{split}$$

According to the Lemma 3.5, we estimate the first term  $I_1$  as following

(29) 
$$||I_1|| \leq c\tau^2 ||v|| \int_{\Gamma_{\theta,\kappa}^{\tau}} |e^{zt_n}| |z||dz|$$

$$\leq c\tau^2 ||v|| \left( \int_{\kappa}^{\frac{\pi}{\tau \sin \theta}} e^{rt_n \cos \theta} r dr + \int_{-\theta}^{\theta} e^{\kappa t_n \cos \psi} \kappa^2 d\psi \right) \leq c\tau^2 t_n^{-2} ||v|| .$$

Using the resolvent estimate (3), we estimate the second term  $I_2$  as following

(30) 
$$||I_2|| \le c \int_{\Gamma_{\theta, r} \setminus \Gamma_{\theta, r}^{\tau}} \left| e^{zt_n} \right| |z|^{-1} ||v||_{L^2(\Omega)} |dz| \le c\tau^2 t_n^{-2} ||v||_{L^2(\Omega)},$$

since

$$\int_{\Gamma_{\theta,\kappa}\backslash\Gamma_{\theta,\kappa}^{\tau}} \left| e^{zt_n} \right| |z|^{-1} |dz| = \int_{\frac{\pi}{\tau\sin\theta}}^{\infty} e^{rt_n\cos\theta} r^{-1} dr \le c\tau^2 \int_{\frac{\pi}{\tau\sin\theta}}^{\infty} e^{rt_n\cos\theta} r dr \le c\tau^2 t_n^{-2}$$

with  $1 \le \left(\frac{\sin \theta}{\pi}\right)^2 \tau^2 r^2$ ,  $r\tau \ge \frac{\pi}{\sin \theta}$ .

From Lemmas 3.6 and 3.7 with  $G(t)=tg(0)+\frac{t^2}{2}g'(0)+\frac{t^2}{2}*g''(t)$ , there exist

$$||I_3|| \le c\tau^2 t_n^{\alpha - 2} ||g(0)|| + c\tau^2 t_n^{\alpha - 1} ||g'(0)|| + c\tau^2 \int_0^{t_n} (t_n - s)^{\alpha - 1} ||g''(s)|| ds.$$

The proof is completed.

**Theorem 3.2** (ID2-BDF2). Let  $V(t_n)$  and  $V^n$  be the solutions of (5) and (9), respectively. Let  $v \in L^2(\Omega)$ ,  $g \in C^1([0,T]; L^2(\Omega))$  and  $\int_0^t (t-s)^{\alpha-1} \|g''(s)\| ds < \infty$ .

Then the following error estimate holds for any  $t_n > 0$ :

$$||V^{n} - V(t_{n})|| \le c\tau^{2} \left( t_{n}^{-2} ||v|| + t_{n}^{\alpha-2} ||g(0)|| + t_{n}^{\alpha-1} ||g'(0)|| \int_{0}^{t_{n}} (t_{n} - s)^{\alpha-1} ||g''(s)|| ds \right).$$

*Proof.* Similar arguments can be performed as Theorem 3.1, we omit it here.  $\Box$ 

# 4. Convergence analysis: Singular source function $t^{\mu}q(x)$ , $\mu > -\alpha$

From Theorem 3.1 and Theorem 3.2, it seems that there are no difference between ID1-BDF2 and ID2-BDF2 for general source function. However, both of them are very different for the singular source function with the form  $t^{\mu}q(x)$ .

**4.1. Low regularity source term.** In the section, we first consider low regularity source term  $g(x,t)=t^{\mu}q(x)$  with  $\mu>0$  for subdiffusion (4). We introduce the polylogarithm function or Bose-Einstein integral

(32) 
$$Li_p(\xi) = \sum_{j=1}^{\infty} \frac{\xi^j}{j^p}, \ p \notin \mathbb{N}.$$

**Lemma 4.1.** [9, 37] Let  $|z\tau| \leq \frac{\pi}{\sin \theta}$  and  $\theta > \pi/2$  be close to  $\pi/2$ , and  $p \neq 1, 2, \ldots$  The series

(33) 
$$Li_p(e^{-z\tau}) = \Gamma(1-p)(z\tau)^{p-1} + \sum_{j=0}^{\infty} (-1)^j \zeta(p-j) \frac{(z\tau)^j}{j!}$$

converges absolutely. Here  $\zeta$  denotes the Riemann zeta function, namely,  $\zeta(p) = Li_p(1)$ .

Let 
$$G(t)=J^1g(t)=rac{t^{\mu+1}}{\mu+1}q$$
. Using  $\widehat{G}(z)=rac{\Gamma(\mu+1)}{z^{\mu+2}}q$  and (10), we have

(34) 
$$V(t) = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa}} e^{zt} (z^{\alpha} - A)^{-1} \left( z^{-1} A v + \frac{\Gamma(\mu + 1)}{z^{\mu + 1}} q \right) dz.$$

From (14), the discrete solution for the subdiffusion (8) is

$$(35) \ V^n = \frac{1}{2\pi i} \int_{\Gamma^\tau_{\theta,\kappa}} e^{zt_n} (\delta^\alpha_\tau(e^{-z\tau}) - A)^{-1} \delta_\tau(e^{-z\tau}) \tau \left( \gamma_1(e^{-z\tau}) \tau A v + \widetilde{G}(e^{-z\tau}) \right) dz$$

with 
$$\gamma_1(e^{-z\tau}) = \frac{e^{-z\tau}}{(1-e^{-z\tau})^2}$$
 and  $\Gamma^{\tau}_{\theta,\kappa} = \{z \in \Gamma_{\theta,\kappa} : |\Im z| \le \pi/\tau\}$ . Here

$$\widetilde{G}(\xi) = \sum_{n=1}^{\infty} G^n \xi^n = q \frac{\tau^{\mu+1}}{\mu+1} \sum_{n=1}^{\infty} \frac{\xi^n}{n^{-\mu-1}} = q \frac{\tau^{\mu+1}}{\mu+1} Li_{-\mu-1}(\xi) \quad \text{with} \quad 0 < \mu < 1.$$

**Lemma 4.2.** Let  $\delta_{\tau}^{\alpha}(\xi)$  is given by (7) and  $\gamma_{l}(\xi) = \sum_{n=1}^{\infty} n^{l} \xi^{n} = \left(\xi \frac{d}{d\xi}\right)^{l} \frac{1}{1-\xi}$  with l=1,2 are given by Lemma 3.3. Then there exist a positive constants c such that

$$\begin{split} & \left\| \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}^{l}(e^{-z\tau}) - (z^{\alpha} - A)^{-1}z^{l} \right\| \leq c\tau^{2}|z|^{l+2-\alpha}, \\ & \left\| \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}^{l}(e^{-z\tau}) \frac{\gamma_{l}(e^{-z\tau})}{l!} \tau^{l+1} - (z^{\alpha} - A)^{-1}z^{-1} \right\| \\ & \leq c\tau^{2}|z|^{1-\alpha} \ \forall z \in \Gamma_{\theta,\kappa}^{\tau}, \end{split}$$

where  $\theta \in (\pi/2, \pi)$  is sufficiently close to  $\pi/2$ .

*Proof.* First we consider

$$(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A)^{-1} \delta_{\tau}^{l}(e^{-z\tau}) - (z^{\alpha} - A)^{-1} z^{l} = I + II$$

with

$$\begin{split} I &= \left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} \left(\delta_{\tau}^{l}(e^{-z\tau}) - z^{l}\right), \\ II &= \left(\left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} - (z^{\alpha} - A)^{-1}\right) z^{l}. \end{split}$$

According to (19) and Lemma 3.1, we obtain

$$||I|| \le c\tau^2 |z|^{l+2-\alpha}.$$

Using the Lemma 3.1, (19), (3) and the identity

$$\left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} - (z^{\alpha} - A)^{-1} = \left(z^{\alpha} - \delta_{\tau}^{\alpha}(e^{-z\tau})\right) \left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} (z^{\alpha} - A)^{-1},$$

we estimate II as following

$$\|II\| \le c\tau^2 |z|^{2+\alpha} c|z|^{-\alpha} c|z|^{-\alpha} |z|^l \le c\tau^2 |z|^{l+2-\alpha}.$$

According to the triangle inequality, the desired result is obtained.

Next we consider

$$\left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} \delta_{\tau}^{l}(e^{-z\tau}) \frac{\gamma_{l}(e^{-z\tau})}{l!} \tau^{l+1} - (z^{\alpha} - A)^{-1} z^{-1} = J_{1} + J_{2}$$

with

$$J_{1} = \left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} \delta_{\tau}^{l}(e^{-z\tau}) \left[\frac{\gamma_{l}(e^{-z\tau})}{l!} \tau^{l+1} - z^{-l-1}\right],$$

$$J_{2} = \left[\left(\delta_{\tau}^{\alpha}(e^{-z\tau}) - A\right)^{-1} \delta_{\tau}^{l}(e^{-z\tau}) - (z^{\alpha} - A)^{-1} z^{l}\right] z^{-l-1}.$$

According to (19) and Lemmas 3.1, 3.2 with l = 1, 2, we obtain

$$||J_1|| \le c\tau^{l+1}|z|^{l-\alpha} \le c\tau^2|z|^{1-\alpha}.$$

From I and II, we have

$$||J_2|| \le c\tau^2 |z|^{l+2-\alpha} |z|^{-l-1} = c\tau^2 |z|^{1-\alpha}.$$

According to the triangle inequality, the desired result is obtained.

**Lemma 4.3.** Let 
$$\widehat{G}(z) = \frac{1}{\mu+1} \frac{\Gamma(\mu+2)}{z^{\mu+2}} q$$
 and  $\widetilde{G}(e^{-z\tau}) = q \frac{\tau^{\mu+1}}{\mu+1} Li_{-\mu-1}(e^{-z\tau})$ . Then  $\left\| \tau \widetilde{G}(e^{-z\tau}) - \widehat{G}(z) \right\| \leq c\tau^{\mu+2} \|q\|, \ \mu \notin \mathbb{N}.$ 

*Proof.* Using the definitions of  $\widehat{G}(z)$  and  $\widetilde{G}(e^{-z\tau})$  and Lemma 4.1 with  $p=-\mu-1$ , we have

$$\begin{split} \left\| \tau \widetilde{G}(e^{-z\tau}) - \widehat{G}(z) \right\| &= \left\| \frac{\tau^{\mu+2}}{(\mu+1)} \left( Li_{-\mu-1}(e^{-z\tau}) - \frac{\Gamma(\mu+2)}{(z\tau)^{\mu+2}} \right) q \right\| \\ &\leq \frac{\tau^{\mu+2}}{(\mu+1)} \left| \sum_{j=0}^{\infty} (-1)^j \zeta(-\mu-1-j) \frac{(z\tau)^j}{j!} \right| \|q\| \leq c\tau^{\mu+2} \|q\| \, . \end{split}$$

The proof is completed.

**Theorem 4.1** (ID1-BDF2). Let  $V(t_n)$  and  $V^n$  be the solutions of (4) and (8), respectively. Let  $v \in L^2(\Omega)$  and  $g(x,t) = t^{\mu}q(x)$ ,  $\mu > 0$ ,  $q(x) \in L^2(\Omega)$ . Then

$$\|V^n - V(t_n)\| \leq c\tau^2 t_n^{-2} \|v\| + c\tau^{\mu+2} t_n^{\alpha-2} \|q\| + c\tau^2 t_n^{\alpha+\mu-2} \|q\| \, .$$

*Proof.* From Theorem 3.1, the desired results is obtained with  $\mu \in \mathbb{N}$ . We next prove the case  $\mu \notin \mathbb{N}$ . Subtracting (34) from (35), we obtain

$$V^n - V(t_n) = I_1 - I_2 + I_3 - I_4$$

$$\begin{split} I_1 = & \frac{1}{2\pi i} \int_{\Gamma^{\tau}_{\theta,\kappa}} e^{zt_n} \left[ \left( \delta^{\alpha}_{\tau}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) \gamma_1(e^{-z\tau}) \tau^2 - (z^{\alpha} - A)^{-1} z^{-1} \right] Avdz, \\ I_2 = & \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa} \backslash \Gamma^{\tau}_{\theta,\kappa}} e^{zt_n} \left( z^{\alpha} - A \right)^{-1} z^{-1} Avdz, \\ I_3 = & \frac{1}{2\pi i} \int_{\Gamma^{\tau}_{\theta,\kappa}} e^{zt_n} \left[ \left( \delta^{\alpha}_{\tau}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) \tau \widetilde{G}(e^{-z\tau}) - (z^{\alpha} - A)^{-1} z \widehat{G}(z) \right] dz, \\ I_4 = & \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa} \backslash \Gamma^{\tau}_{\theta}} e^{zt_n} \left( z^{\alpha} - A \right)^{-1} z \widehat{G}(z) dz. \end{split}$$

According to (29) and (30), we estimate  $I_1$  and  $I_2$  as following

$$||I_1|| \le c\tau^2 t_n^{-2} ||v||$$
 and  $||I_2|| \le c\tau^2 t_n^{-2} ||v||$ .

From (31), we estimate that  $I_4$  is similar to  $I_2$  as following

$$||I_4|| \le c \int_{\Gamma_{\theta,\kappa} \setminus \Gamma_{\theta,\kappa}^{\tau}} \left| e^{zt_n} \right| |z|^{-\alpha} \left\| z \widehat{G}(z) \right\| |dz|$$

$$\le c \int_{\Gamma_{\theta,\kappa} \setminus \Gamma_{\theta,\kappa}^{\tau}} \left| e^{zt_n} \right| |z|^{-\alpha} |z|^{-\mu-1} \|q\| |dz| \le c\tau^2 t_n^{\alpha+\mu-2} \|q\|.$$

Finally we consider  $I_3 = I_{31} + I_{32}$  with

$$\begin{split} I_{31} = & \frac{1}{2\pi i} \int_{\Gamma^{\tau}_{\theta,\kappa}} e^{zt_n} \left( \delta^{\alpha}_{\tau}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) \left( \tau \widetilde{G}(e^{-z\tau}) - \widehat{G}(z) \right) dz, \\ I_{32} = & \frac{1}{2\pi i} \int_{\Gamma^{\tau}_{\theta,\kappa}} e^{zt_n} \left( \left( \delta^{\alpha}_{\tau}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}(e^{-z\tau}) - (z^{\alpha} - A)^{-1} z \right) \widehat{G}(z) dz. \end{split}$$

According to (19) and Lemmas 3.1 and 4.3, there exists

$$||I_{31}|| \le c\tau^{\mu+2} ||q|| \int_{\Gamma_{\theta,\kappa}^{\tau}} |e^{zt_n}| |z|^{1-\alpha} |dz| \le c\tau^{\mu+2} t_n^{\alpha-2} ||q||.$$

From Lemma 4.2 and  $\hat{G}(z) = \frac{1}{\mu+1} \frac{\Gamma(\mu+2)}{z^{\mu+2}} q$ , we estimate  $I_{32}$  as following

$$||I_{32}|| \le c\tau^2 ||q|| \int_{\Gamma_{n}^{\tau}} |e^{zt_n}| |z|^{3-\alpha} |z|^{-\mu-2} |dz| \le c\tau^2 t_n^{\alpha+\mu-2} ||q||.$$

By the triangle inequality, the desired result is obtained.

**4.2.** Singular source term. In this subsection, we consider the singular source

term 
$$g(x,t) = t^{\mu}q(x)$$
 with  $\mu > -\alpha$  for subdiffusion (5).  
Let  $\mathcal{G}(t) = J^2g(t) = \frac{t^{\mu+2}}{(\mu+1)(\mu+2)}q$ . Using  $\widehat{\mathcal{G}}(z) = \frac{\Gamma(\mu+1)}{z^{\mu+3}}q$  and (12), we have

(36) 
$$V(t) = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa}} e^{zt} (z^{\alpha} - A)^{-1} \left( z^{-1} A v + \frac{\Gamma(\mu + 1)}{z^{\mu + 1}} q \right) dz.$$

From (16), it yields

$$V^n = \frac{\tau}{2\pi i} \int_{\Gamma^\tau_{a,r}} e^{zt_n} \left( \delta^\alpha_\tau(e^{-z\tau}) - A \right)^{-1} \delta^2_\tau(e^{-z\tau}) \left( \frac{\gamma_2(e^{-z\tau})}{2} \tau^2 A v + \widetilde{\mathcal{G}}(e^{-z\tau}) \right) dz$$

with 
$$\frac{\gamma_2(e^{-z\tau})}{2} = \frac{e^{-z\tau} + e^{-2z\tau}}{2(1 - e^{-z\tau})^3}$$
 and  $\Gamma_{\theta,\kappa}^{\tau} = \{z \in \Gamma_{\theta,\kappa} : |\Im z| \le \pi/\tau\}$ . Here

$$\widetilde{\mathcal{G}}(\xi) = \sum_{n=1}^{\infty} \mathcal{G}^n \xi^n = q \frac{\tau^{\mu+2}}{(\mu+2)(\mu+1)} \sum_{n=1}^{\infty} \frac{\xi^n}{n^{-\mu-2}} = q \frac{\tau^{\mu+2}}{(\mu+2)(\mu+1)} Li_{-\mu-2}(\xi).$$

**Lemma 4.4.** Let 
$$\widehat{\mathcal{G}}(z) = q \frac{\Gamma(\mu+1)}{z^{\mu+3}}$$
 and  $\widetilde{\mathcal{G}}(e^{-z\tau}) = q \frac{\tau^{\mu+2}}{(\mu+2)(\mu+1)} Li_{-\mu-2}(e^{-z\tau})$ . Then  $\left\|\tau \widetilde{\mathcal{G}}(e^{-z\tau}) - \widehat{\mathcal{G}}(z)\right\| \leq c\tau^{\mu+3} \|q\|, \ \mu \notin \mathbb{N}.$ 

*Proof.* From Lemma 4.1, we have

$$\begin{split} \left\| \tau \widetilde{\mathcal{G}}(e^{-z\tau}) - \widehat{\mathcal{G}}(z) \right\| &= \left\| \frac{\tau^{\mu+3}}{(\mu+2)(\mu+1)} \left( Li_{-\mu-2}(e^{-z\tau}) - \frac{\Gamma(\mu+3)}{(z\tau)^{\mu+3}} \right) q \right\| \\ &\leq \frac{\tau^{\mu+3}}{(\mu+2)(\mu+1)} \left| \sum_{j=0}^{\infty} (-1)^j \zeta(-\mu-2-j) \frac{(z\tau)^j}{j!} \right| \|q\| \\ &\leq c\tau^{\mu+3} \|q\| \, . \end{split}$$

The proof is completed.

**Theorem 4.2** (ID2-BDF2). Let  $V(t_n)$  and  $V^n$  be the solutions of (5) and (9), respectively. Let  $v \in L^2(\Omega)$  and  $g(x,t) = t^{\mu}q(x)$ ,  $\mu > -\alpha$ ,  $q(x) \in L^2(\Omega)$ . Then

$$\|V^n - V(t_n)\| \le c\tau^2 t_n^{-2} \|v\| + c\tau^{\mu+3} t_n^{\alpha-3} \|q\| + c\tau^2 t_n^{\alpha+\mu-2} \|q\|.$$

*Proof.* From Theorem 3.1, the desired results is obtained with  $\mu \in \mathbb{N}$ . We next prove the case  $\mu \notin \mathbb{N}$ . Subtracting (12) from (16), we obtain

$$V^n - V(t_n) = I_1 - I_2 + I_3 - I_4$$

with

$$\begin{split} I_1 = & \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa}^{\tau}} e^{zt_n} \left[ \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}^2(e^{-z\tau}) \frac{e^{-z\tau} + e^{-2z\tau}}{2(1 - e^{-z\tau})^3} \tau^3 \right. \\ & - \left( z^{\alpha} - A \right)^{-1} z^{-1} \right] A v dz, \\ I_2 = & \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa} \setminus \Gamma_{\theta,\kappa}^{\tau}} e^{zt_n} \left( z^{\alpha} - A \right)^{-1} z^{-1} A v dz, \\ I_3 = & \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa}^{\tau}} e^{zt_n} \left[ \left( \delta_{\tau}^{\alpha}(e^{-z\tau}) - A \right)^{-1} \delta_{\tau}^2(e^{-z\tau}) \tau \widetilde{\mathcal{G}}(e^{-z\tau}) - (z^{\alpha} - A)^{-1} z^2 \widehat{\mathcal{G}}(z) \right] dz, \\ I_4 = & \frac{1}{2\pi i} \int_{\Gamma_{\theta,\kappa} \setminus \Gamma_{\theta,\kappa}^{\tau}} e^{zt_n} \left( z^{\alpha} - A \right)^{-1} z^2 \widehat{\mathcal{G}}(z) dz. \end{split}$$

Using (29), (30) and Lemma 4.2, we estimate  $I_1$  and  $I_2$  as following

$$||I_1|| \le c\tau^2 t_n^{-2} ||v||$$
 and  $||I_2|| \le c\tau^2 t_n^{-2} ||v||$ .

By (31), we estimate that  $I_4$  is similar to  $I_2$  as following

$$||I_4|| \le c \int_{\Gamma_{\theta,\kappa} \setminus \Gamma_{\theta,\kappa}^{\tau}} \left| e^{zt_n} \right| |z|^{-\alpha} ||z|^{2} \widehat{\mathcal{G}}(z) ||dz|$$

$$\le c ||q|| \int_{\Gamma_{\theta,\kappa} \setminus \Gamma_{\theta,\kappa}^{\tau}} \left| e^{zt_n} \right| |z|^{-\alpha} |z|^{-\mu-1} |dz| \le c\tau^2 t_n^{\alpha+\mu-2} ||q||.$$

Finally we consider  $I_3 = I_{31} + I_{32}$  with

$$\begin{split} I_{31} = & \frac{1}{2\pi i} \int_{\Gamma^{\tau}_{\theta,\kappa}} e^{zt_n} \left( \delta^{\alpha}_{\tau}(e^{-z\tau}) - A \right)^{-1} \delta^{2}_{\tau}(e^{-z\tau}) \left( \tau \widetilde{\mathcal{G}}(e^{-z\tau}) - \widehat{\mathcal{G}}(z) \right) dz, \\ I_{32} = & \frac{1}{2\pi i} \int_{\Gamma^{\tau}_{\theta,\kappa}} e^{zt_n} \left( \left( \delta^{\alpha}_{\tau}(e^{-z\tau}) - A \right)^{-1} \delta^{2}_{\tau}(e^{-z\tau}) - (z^{\alpha} - A)^{-1} z^2 \right) \widehat{\mathcal{G}}(z) dz. \end{split}$$

According to (19) and Lemmas 3.1 and 4.4, there exists

$$||I_{31}|| \le c\tau^{\mu+3} ||q|| \int_{\Gamma_{a,\mu}^{\tau}} |e^{zt_n}| |z|^{2-\alpha} |dz| \le c\tau^{\mu+3} t_n^{\alpha-3} ||q||.$$

From Lemma 4.2, we estimate  $I_{32}$  as following

$$||I_{32}|| \le c\tau^2 ||q|| \int_{\Gamma_{\theta,\kappa}^{\tau}} |e^{zt_n}| |z|^{4-\alpha} |z|^{-\mu-3} |dz|$$

$$\le c\tau^2 ||q|| \int_{\Gamma_{\theta,\kappa}^{\tau}} |e^{zt_n}| |z|^{1-\alpha-\mu} |dz| \le c\tau^2 t_n^{\alpha+\mu-2} ||q||.$$

By the triangle inequality, the desired result is obtained.

# 5. Convergence analysis: Source function $t^{\mu} \circ f(x,t)$ with $\mu > -1$

Based on the discussion of Section 3 and 4, we now analyse the error estimates for subdiffusion (1) with the singular source term  $t^{\mu} \circ f(x,t)$ .

**5.1.** Convergence analysis: Convolution source function  $t^{\mu} * f(t)$ ,  $\mu > -1$ . Let f(t) = f(0) + tf'(0) + t \* f''(t). Then we obtain

$$g(t) = t^{\mu} * f(t) = \frac{t^{\mu+1}f(0)}{\mu+1} + \frac{t^{\mu+2}f'(0)}{(\mu+1)(\mu+2)} + t^{\mu} * t * f''(t).$$

Let  $G(t) = J^1 g(t) = \frac{1}{\mu+1} t^{\mu+1} * f(t)$  with G(0) = 0. It yields

$$G(t) = \frac{t^{\mu+2}f(0)}{(\mu+1)(\mu+2)} + \frac{t^{\mu+3}f'(0)}{(\mu+1)(\mu+2)(\mu+3)} + \frac{1}{\mu+1}t^{\mu+1} * t * f''(t)$$

$$= \frac{t^{\mu+2}f(0)}{(\mu+1)(\mu+2)} + \frac{t^{\mu+3}f'(0)}{(\mu+1)(\mu+2)(\mu+3)} + \frac{t^2}{2} * (t^{\mu} * f''(t)),$$

where we use

$$t^{\mu+1}*t = \int_0^t (t-s)^{\mu+1} s ds = \frac{\mu+1}{2} \int_0^t (t-s)^{\mu} s^2 ds = \frac{\mu+1}{2} t^2 * t^{\mu}.$$

**Lemma 5.1.** Let  $V(t_n)$  and  $V^n$  be the solutions of (4) and (8), respectively. Let v = 0,  $G(t) := \frac{t^2}{2} * (t^{\mu} * f''(t))$  with  $\mu > -1$  and  $\int_0^t (t-s)^{\alpha-1} s^{\mu} * ||f''(s)|| ds < \infty$ . Then

$$||V(t_n) - V^n|| \le c\tau^2 \int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\mu} * ||f''(s)|| \, ds \le c\tau^2 \int_0^{t_n} (t_n - s)^{\alpha + \mu} ||f''(s)|| \, ds.$$

*Proof.* By Lemma 3.7 with  $g''(t) = t^{\mu} * f''(t)$ , we obtain

$$||V(t_n) - V^n|| \le c\tau^2 \int_0^{t_n} (t_n - s)^{\alpha - 1} ||s^{\mu} * f''(s)|| ds$$

$$\le c\tau^2 \int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\mu} * ||f''(s)|| ds$$

$$= c\tau^2 \left( t^{\alpha - 1} * t^{\mu} \right) * ||f''(t)||_{t = t_n} \le c\tau^2 \int_0^{t_n} (t_n - s)^{\alpha + \mu} ||f''(s)|| ds.$$

The proof is completed.

**Theorem 5.1** (ID1-BDF2). Let  $V(t_n)$  and  $V^n$  be the solutions of (4) and (8), respectively. Let  $v \in L^2(\Omega)$ ,  $g(t) = t^{\mu} * f(t)$  with  $\mu > -1$  and  $f \in C^1([0,T]; L^2(\Omega))$ ,  $\int_0^t (t-s)^{\alpha-1} s^{\mu} * \|f''(s)\| ds < \infty$ . Then

$$||V^n - V(t_n)||$$

$$\leq c\tau^{2} \left( t_{n}^{-2} \|v\| + t_{n}^{\alpha+\mu-1} \|f(0)\| + t_{n}^{\alpha+\mu} \|f'(0)\| + \int_{0}^{t_{n}} (t_{n} - s)^{\alpha-1} s^{\mu} * \|f''(s)\| \, ds \right)$$

$$\leq c\tau^{2} \left( t_{n}^{-2} \|v\| + t_{n}^{\alpha+\mu-1} \|f(0)\| + t_{n}^{\alpha+\mu} \|f'(0)\| + \int_{0}^{t_{n}} (t_{n} - s)^{\alpha+\mu} \|f''(s)\| \, ds \right).$$

*Proof.* According to Theorem 4.1, Lemma 5.1, and similar treatment of the initial data v in Theorem 3.1, the desired result is obtained.

**5.2. Convergence analysis: product source function**  $t^{\mu}f(t)$ ,  $\mu > 0$ . Let  $G(t) = J^1g(t)$  and f(t) = f(0) + tf'(0) + t \* f''(t). Then we have

$$G(t) = 1 * (t^{\mu} f(t)) = \frac{t^{\mu+1} f(0)}{\mu+1} + \frac{t^{\mu+2} f'(0)}{\mu+2} + 1 * [t^{\mu} (t * f''(t))].$$

Let  $h(t) = t^{\mu} (t * f''(t))$  with h(0) = 0. It leads to

$$h'(t) = \mu t^{\mu - 1} \left( t * f''(t) \right) + t^{\mu} \left( 1 * f''(t) \right)$$

with h'(0) = 0, since

$$|h'(t)| \le \left| \mu t^{\mu - 1} \int_0^t (t - s) f''(s) ds \right| + \left| t^{\mu} \int_0^t f''(s) ds \right|$$
  
$$\le (\mu + 1) t^{\mu} \int_0^t |f''(s)| ds, \ \mu > 0.$$

Moreover, there exists

(37) 
$$h''(t) = \mu (\mu - 1) t^{\mu - 2} (t * f''(t)) + 2\mu t^{\mu - 1} (1 * f''(t)) + t^{\mu} f''(t).$$

Thus one has

(38) 
$$1 * h(t) = th(0) + \frac{t^2}{2}h'(0) + \frac{t^2}{2} * h''(t) = \frac{t^2}{2} * h''(t).$$

**Lemma 5.2.** Let  $V(t_n)$  and  $V^n$  be the solutions of (4) and (8), respectively. Let v = 0,  $G(t) = 1 * [t^{\mu}(t * f''(t))]$  with  $\mu > 0$  and  $f \in C^1([0,T];L^2(\Omega))$ ,  $\int_0^t ||f''(s)|| ds < \infty$ ,  $\int_0^t (t-s)^{\alpha-1} s^{\mu} ||f''(s)|| ds < \infty$ . Then

$$||V(t_n) - V^n|| \le c\tau^2 \left( t_n^{\alpha + \mu - 1} \int_0^{t_n} ||f''(s)|| \, ds + \int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\mu} \, ||f''(s)|| \, ds \right).$$

*Proof.* Let  $h(t) = t^{\mu} (t * f''(t))$ . From (38), we have  $G(t) = 1 * h(t) = \frac{t^2}{2} * h''(t)$ . According to Lemma 3.7 and (37), it yields

$$||V(t_n) - V^n|| \le c\tau^2 \int_0^{t_n} (t_n - s)^{\alpha - 1} ||h''(s)|| ds \le c\tau^2 (I_1 + I_2 + I_3)$$

with

$$I_{1} = \int_{0}^{t_{n}} (t_{n} - s)^{\alpha - 1} \| s^{\mu - 2} (s * f''(s)) \| ds,$$

$$I_{2} = \int_{0}^{t_{n}} (t_{n} - s)^{\alpha - 1} \| s^{\mu - 1} (1 * f''(s)) \| ds \text{ and } I_{3} = \int_{0}^{t_{n}} (t_{n} - s)^{\alpha - 1} \| s^{\mu} f''(s) \| ds.$$

We estimate  $I_1$  as following

$$I_{1} = \int_{0}^{t_{n}} (t_{n} - s)^{\alpha - 1} s^{\mu - 1} \left\| \int_{0}^{s} \frac{s - w}{s} f''(w) dw \right\| ds$$

$$\leq \int_{0}^{t_{n}} (t_{n} - s)^{\alpha - 1} s^{\mu - 1} \int_{0}^{t_{n}} \|f''(w)\| dw ds = B(\alpha, \mu) t_{n}^{\alpha + \mu - 1} \int_{0}^{t_{n}} \|f''(w)\| dw,$$

since

$$\int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\mu - 1} ds = t_n^{\alpha + \mu - 1} \int_0^1 (1 - s)^{\alpha - 1} s^{\mu - 1} ds = B(\alpha, \mu) t_n^{\alpha + \mu - 1}.$$

Similarly, we estimate  $I_2$  as following

$$I_2 \le \int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\mu - 1} \int_0^{t_n} \|f''(w)\| \, dw ds = B(\alpha, \mu) t_n^{\alpha + \mu - 1} \int_0^{t_n} \|f''(w)\| \, dw.$$

By the triangle inequality, we obtain

$$||V(t_n) - V^n|| \le c\tau^2 \left( t_n^{\alpha + \mu - 1} \int_0^{t_n} ||f''(s)|| \, ds + \int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\mu} \, ||f''(s)|| \, ds \right).$$

The proof is completed.

**Theorem 5.2** (ID1-BDF2). Let  $V(t_n)$  and  $V^n$  be the solutions of (4) and (8), respectively. Let  $v \in L^2(\Omega)$ ,  $g(t) = t^{\mu} f(t)$  with  $\mu > 0$  and  $f \in C^1([0,T];L^2(\Omega))$ ,  $\int_0^t \|f''(s)\| \, ds < \infty$ ,  $\int_0^t (t-s)^{\alpha-1} s^{\mu} \|f''(s)\| \, ds < \infty$ . Then

$$||V^{n} - V(t_{n})|| \leq c\tau^{2} \left(t_{n}^{-2}||v|| + t_{n}^{\alpha+\mu-2}||f(0)|| + t_{n}^{\alpha+\mu-1}||f'(0)||\right) + c\tau^{2} \left(t_{n}^{\alpha+\mu-1} \int_{0}^{t_{n}} ||f''(s)|| ds + \int_{0}^{t_{n}} (t_{n} - s)^{\alpha-1} s^{\mu} ||f''(s)|| ds\right).$$

*Proof.* According to Theorem 4.1, Lemma 5.2, and similar treatment of the initial data v in Theorem 3.1, the desired result is obtained.

**5.3. Convergence analysis: product source function**  $t^{\mu}f(t)$ ,  $-\alpha < \mu < 0$ . Let  $\mathcal{G}(t) = J^2g(t)$  and f(t) = f(0) + tf'(0) + t \* f''(t). Then we have

$$\mathcal{G}(t) = t * (t^{\mu} f(t)) = \frac{t^{\mu+2} f(0)}{(\mu+1)(\mu+2)} + \frac{t^{\mu+3} f'(0)}{(\mu+2)(\mu+3)} + t * [t^{\mu} (t * f''(t))].$$

Let  $h(t) = t^{\mu} (t * f''(t))$  with h(0) = 0. It leads to

$$h'(t) = \mu t^{\mu - 1} (t * f''(t)) + t^{\mu} (1 * f''(t)),$$

which implies

$$|h'(0)| \le (\mu + 1) \int_0^t s^{\mu} |f''(s)| ds,$$

since

$$|h'(t)| \le (\mu+1)t^{\mu} \int_0^t |f''(s)| \, ds \le (\mu+1) \int_0^t s^{\mu} |f''(s)| \, ds \text{ with } -1 < \mu < 0.$$

Thus we get

(39) 
$$t * h(t) = \frac{t^2}{2}h(0) + \frac{t^3}{6}h'(0) + \frac{t^3}{6} * h''(t) = \frac{t^3}{6}h'(0) + \frac{t^3}{6} * h''(t).$$

**Lemma 5.3.** Let  $V(t_n)$  and  $V^n$  be the solutions of (5) and (9), respectively. Let v=0,  $\mathcal{G}(t)=t*[t^{\mu}(t*f''(t))]$  with  $-\alpha<\mu<0$  and  $f\in C^1([0,T];L^2(\Omega))$ ,  $\int_0^t s^{\frac{\mu-1}{2}}\|f''(s)\|\,ds<\infty$ ,  $\int_0^t (t-s)^{\alpha-1}s^{\mu}\|f''(s)\|\,ds<\infty$ . Then

$$||V(t_n) - V^n|| \le c\tau^2 \left( t_n^{\alpha + \frac{\mu - 1}{2}} \int_0^{t_n} s^{\frac{\mu - 1}{2}} ||f''(s)|| \, ds + \int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\mu} ||f''(s)|| \, ds \right).$$

*Proof.* Let  $h(t) = t^{\mu} (t * f''(t))$ . From (39), we have

$$G(t) = t * h(t) = \frac{t^3}{6}h'(0) + \frac{t^3}{6} * h''(t).$$

According to Theorems 4.2, 3.2 and (37), it yields

$$||V(t_n) - V^n|| \le c\tau^2 \left( t_n^{\alpha - 1} ||h'(0)|| + \int_0^{t_n} (t_n - s)^{\alpha - 1} ||h''(s)|| ds \right)$$
  
 
$$\le c\tau^2 \left( I_1 + I_2 + I_3 + I_4 \right)$$

with

$$I_{1} = t_{n}^{\alpha-1} \|h'(0)\|, \quad I_{2} = \int_{0}^{t_{n}} (t_{n} - s)^{\alpha-1} \|s^{\mu-2} (s * f''(s))\| ds,$$

$$I_{3} = \int_{0}^{t_{n}} (t_{n} - s)^{\alpha-1} \|s^{\mu-1} (1 * f''(s))\| ds \text{ and } I_{4} = \int_{0}^{t_{n}} (t_{n} - s)^{\alpha-1} \|s^{\mu} f''(s)\| ds.$$

Since

$$I_1 = t_n^{\alpha - 1} \|h'(0)\| \le c t_n^{\alpha - 1} \int_0^{t_n} s^{\mu} \|f''(s)\| \, ds \le c \int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\mu} \|f''(s)\| \, ds,$$

and

$$I_{2} = \int_{0}^{t_{n}} (t_{n} - s)^{\alpha - 1} s^{\frac{\mu - 1}{2}} \left\| \int_{0}^{s} \frac{s - w}{s} s^{\frac{\mu - 1}{2}} f''(w) dw \right\| ds$$

$$\leq \int_{0}^{t_{n}} (t_{n} - s)^{\alpha - 1} s^{\frac{\mu - 1}{2}} \int_{0}^{t_{n}} w^{\frac{\mu - 1}{2}} \|f''(w)\| dw ds$$

$$\leq c t_{n}^{\alpha + \frac{\mu - 1}{2}} \int_{0}^{t_{n}} w^{\frac{\mu - 1}{2}} \|f''(w)\| dw,$$

where we use

$$\int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\frac{\mu - 1}{2}} ds = t_n^{\alpha + \frac{\mu - 1}{2}} \int_0^1 (1 - s)^{\alpha - 1} s^{\frac{\mu - 1}{2}} ds = B\left(\alpha, \frac{\mu + 1}{2}\right) t_n^{\alpha + \frac{\mu - 1}{2}}.$$

Similarly, we estimate  $I_3$  as following

$$I_3 \le \int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\mu - 1} \int_0^{t_n} \|f''(w)\| \, dw ds \le c t_n^{\alpha + \frac{\mu - 1}{2}} \int_0^{t_n} w^{\frac{\mu - 1}{2}} \|f''(w)\| \, dw.$$

By the triangle inequality, we obtain

$$||V(t_n) - V^n|| \le c\tau^2 \left( t_n^{\alpha + \frac{\mu - 1}{2}} \int_0^{t_n} s^{\frac{\mu - 1}{2}} ||f''(s)|| \, ds + \int_0^{t_n} (t_n - s)^{\alpha - 1} s^{\mu} ||f''(s)|| \, ds \right).$$

The proof is completed.

**Theorem 5.3** (ID2-BDF2). Let  $V(t_n)$  and  $V^n$  be the solutions of (5) and (9), respectively. Let  $v \in L^2(\Omega)$ ,  $g(t) = t^{\mu} f(t)$  with  $-\alpha < \mu < 0$  and  $f \in C^1([0,T];L^2(\Omega))$ ,  $\int_0^t s^{\frac{\mu-1}{2}} \|f''(s)\| ds < \infty$ ,  $\int_0^t (t-s)^{\alpha-1} s^{\mu} \|f''(s)\| ds < \infty$ . Then

$$||V^{n} - V(t_{n})|| \leq c\tau^{2} \left(t_{n}^{-2}||v|| + t_{n}^{\alpha+\mu-2}||f(0)|| + t_{n}^{\alpha+\mu-1}||f'(0)||\right) + c\tau^{2} \left(t_{n}^{\alpha+\frac{\mu-1}{2}} \int_{0}^{t_{n}} s^{\frac{\mu-1}{2}} ||f''(s)|| ds + \int_{0}^{t_{n}} (t_{n} - s)^{\alpha-1} s^{\mu} ||f''(s)|| ds\right).$$

*Proof.* According to Theorem 4.2, Lemma 5.3, and similar treatment of the initial data v in Theorem 3.1, the desired result is obtained.

**Remark 5.1.** Theorems 4.2 and 5.3 are naturally extended to  $\mu > -1$ . However, it is crucial to emphasize that when  $\mu \in (-1, -\alpha)$ ,

$$||u(t)||_{L^2(\Omega)} \sim ct^{\mu+\alpha} \to \infty \quad as \quad t \to 0^+,$$

and consequently, the initial condition cannot be satisfied in the conventional sense. As a result, the weak solution to the initial value problem (1) needs to be redefined in a weaker sense to fulfill the initial condition. In order to avoid redundant discussions of the model, we have decided to present our argument for the source term  $t^{\mu} \circ f(x,t)$ , where  $\mu > -1$  if  $\circ$  denotes convolution, while  $\mu > -\alpha$  if  $\circ$  denotes product.

# 6. Numerical results

We numerically verify the above theoretical results and the discrete  $L^2$ -norm is used to measure the numerical errors. In the space direction, it is discretized with the spectral collocation method with the Chebyshev-Gauss-Lobatto points [30]. Here we mainly focus on the time direction convergence order, since the convergence rate of the spatial discretization is well understood. Since the analytic solutions are unknown, the order of the convergence of the numerical results is computed by the following formula

Convergence Rate = 
$$\frac{\ln \left( \|u^{N/2} - u^N\| / \|u^N - u^{2N}\| \right)}{\ln 2}$$

with  $u^N = V^N + v$  in (8).

**6.1.** Subdiffusion model with Dirichlet and Neumann boundary conditions. In the experiment, several algorithms including the correction BDF2 methods [11], FBDF22 Method [39] are carried out and compared with IDm-method:

(40) BDF2 Method: 
$$\partial_{\tau}^{\alpha} V^n - AV^n = Av + g^n$$
.

(41) 
$$\operatorname{Corr} - \operatorname{BDF2} \operatorname{Method} : \partial_{\tau}^{\alpha} V^{n} - AV^{n} = \frac{3}{2} Av + \frac{1}{2} g^{0} + g^{n}.$$

(42) FBDF22 Method : 
$$\begin{cases} \partial_{\tau}^{\alpha} U^{n} - AU^{n} = \partial_{\tau} \left[ J^{2} \left( g(t) + \frac{t^{-\alpha} v}{\Gamma(1-\alpha)} \right) \Big|_{t=t_{n}} \right], \\ u^{n} = \partial_{\tau} U^{n}, \quad u^{n} = V^{n} + v. \end{cases}$$

**Example 6.1** (Dirichlet boundary conditions). Let T = 1 and  $\Omega = (-1, 1)$ . Consider subdiffusion (1) with

$$v(x) = \sin(x)\sqrt{1-x^2}$$
 and  $g(x,t) = (1+t^{\mu}+t^{\alpha\mu})\circ(1-t)^{\beta}e^x(1+\chi_{(0,1)}(x)).$ 

Here  $J^mg(x,t)=t^{m-1}*g(x,t),\ m=1,2$  is calculated by JacobiGL Algorithm [2, 7], which generates the nodes and weights of Gauss-Labatto integral with the weighting function of the form  $(1-t)^\mu$  or  $(1+t)^\mu$ .

Table 1. The discrete  $L^2$ -norm  $\|u^N-u^{2N}\|$  and convergent order of schemes (40), (41) and (8), (9) with  $\beta=0,\ \alpha=0.7$ . Here  $\circ$  denotes the dot product.

Scheme	$\mu$	N = 50	N = 100	N = 200	N = 400	N = 800
	0.8	2.4743e-03	1.1981e-03	5.8732e-04	2.9005e-04	1.4390e-04
BDF2			1.0462	1.0286	1.0178	1.0113
DDF 2	-0.8	1.5948e-01	1.3256 e-01	1.1109e-01	9.3707e-02	7.9450 e-02
			0.26679	0.25489	0.24549	0.23811
	0.8	9.4381e-05	3.6107e-05	1.3189e-05	4.6888e-06	1.6386e-06
Corr-BDF2			1.3862	1.4529	1.4921	1.5168
	-0.8	NaN	NaN	NaN	NaN	NaN
	0.8	8.4501e-05	2.0876e-05	5.1883e-06	1.2933e-06	3.2283e-07
FBDF22			2.0171	2.0085	2.0042	2.0022
r DDF 22	-0.8	5.8209 e-04	1.4211e-04	3.5057e-05	8.6949 e - 06	2.1627e-06
			2.0343	2.0192	2.0114	2.0073
	0.8	1.6660e-04	4.1216e-05	1.0249e-05	2.5553e-06	6.3792e-07
ID1-BDF2			2.0151	2.0077	2.0040	2.0021
1D1-DD1 2	-0.8	6.7744e-03	3.0380e-03	1.3367e-03	5.8281e-04	2.5299e-04
			1.1570	1.1844	1.1976	1.2039
ID2-BDF2	0.8	3.2389e-04	7.9995e-05	1.9879e-05	4.9539e-06	1.2374e-06
			2.0175	2.0087	2.0046	2.0013
1D2-DDF 2	-0.8	2.1611e-03	5.2769 e-04	1.3018e-04	3.2292 e- 05	8.0280 e-06
			2.0340	2.0192	2.0112	2.0081

Table 2. The discrete  $L^2$ -norm  $\|u^N-u^{2N}\|$  and convergent order of schemes (8) and (9) with  $\beta=1.9$ , respectively. Here  $\circ$  denotes the dot product.

Scheme	$\alpha$	$\mu$	N = 50	N = 100	N = 200	N = 400	N = 800
		0.5	1.5025e-03	3.9778e-04	1.0433e-04	2.7198e-05	7.0660e-06
	0.3			1.9174	1.9307	1.9396	1.9445
	0.5	-0.9	4.9903e-03	2.7664e-03	1.4020 e - 03	6.8259 e-04	3.2574 e-04
ID1-BDF2				0.85109	0.98050	1.0384	1.0673
IDI-DDI 2		0.5	6.8462 e-04	1.8033e-04	4.6484e-05	1.1840e-05	2.9948e-06
	0.7			1.9247	1.9558	1.9731	1.9831
	0.7	-0.9	2.0722e-02	1.0219e-02	4.8849e-03	2.3017e-03	1.0770e-03
				1.0199	1.0648	1.0856	1.0956
		0.5	3.1810e-03	8.4340e-04	2.2164e-04	5.7938e-05	1.5180e-05
	0.3			1.9152	1.9280	1.9356	1.9323
	0.5	-0.9	4.6179e-03	1.1806e-03	3.0298e-04	7.7857e-05	2.0182e-05
ID2-BDF2				1.9677	1.9622	1.9603	1.9478
1D2-DD1 2		0.5	1.9266e-03	5.0536e-04	1.3015e-04	3.3167e-05	8.4027e-06
	0.7			1.9307	1.9571	1.9724	1.9808
	0.7	-0.9	7.2846e-03	1.8010e-03	4.4808e-04	1.1179e-04	2.7922e-05
				2.0161	2.0070	2.0030	2.0013

Table 3. The discrete  $L^2$ -norm  $||u^N - u^{2N}||$  and convergent order of schemes (40) and (8) with  $\beta = 1.9$ , respectively. Here  $\circ$  denotes the Laplace convolution.

Scheme	α	$\mu$	N = 50	N = 100	N = 200	N = 400	N = 800
		-0.2	6.4420e-05	1.2431e-05	2.6710e-06	6.1586e-07	1.4766e-07
	0.3			2.3735	2.2185	2.1167	2.0603
	0.5	-0.8	1.6132 e-03	4.2435 e-04	1.0992e-04	2.8213 e-05	7.2033e-06
ID1-BDF2				1.9266	1.9487	1.9621	1.9696
1D1-DDF 2		-0.2	2.8145e-04	6.7873e-05	1.6649 e - 05	4.1218e-06	1.0253e-06
	0.7			2.0520	2.0274	2.0141	2.0072
	0.7	-0.8	6.3566e-04	1.7068e-04	4.4407e-05	1.1358e-05	2.8782e-06
				1.8969	1.9425	1.9671	1.9805

**Example 6.2** (Neumann boundary conditions). Let T=1 and  $\Omega=(-1,1)$ . Consider subdiffusion with the Neumann boundary conditions

$$\begin{cases} {}^{C}D_{t}^{\alpha}u(x,t) - Au(x,t) = g(x,t), & x \in \Omega, \ t \in (0,T], \\ u(x,0) = \sqrt{1 - x^{2}/2}, & x \in \Omega, \\ u_{x}(0,t) = e^{-t}, \ u(1,t) = 1, & x \in \partial\Omega, \ t \in [0,T], \end{cases}$$

where  $g(x,t) = (1 + t^{\mu} + t^{\alpha\mu})e^x(1 + \chi_{(0,1)}(x)).$ 

Table 4. The discrete  $L^2$ -norm  $||u^N - u^{2N}||$  and convergent order for Example 6.2 with  $\alpha = 0.7$ .

Scheme	$\mu$	N = 50	N = 100	N = 200	N = 400	N = 800
	0.8	1.5408e-02	3.9107e-03	9.8528e-04	2.4730e-04	6.1950 e-05
ID1-BDF2			1.9782	1.9888	1.9943	1.9971
1D1-DD1 2	-0.8	2.7374e-01	1.2905 e-01	5.8398e-02	2.5866e-02	1.1329e-02
			1.0848	1.1440	1.1749	1.1911
	0.8	6.9166e-03	1.7704e-03	4.4795e-04	1.1266e-04	2.8253 e-05
ID2-BDF2			1.9660	1.9827	1.9913	1.9955
	-0.8	4.8497e-02	1.2342e-02	3.1180e-03	7.8456e-04	1.9699e-04
			1.9743	1.9849	1.9907	1.9938

**Example 6.3** (Two-dimensional subdiffusion model). Let T=1 and  $\Omega=(-1,1)\times (-1,1)$ . Consider the two-dimensional subdiffusion model with the inhomogeneous Dirichlet boundary conditions

(43) 
$$\begin{cases} {}^{C}D_{t}^{\alpha}u(x,y,t) - Au(x,y,t) = g(x,y,t), & (x,y) \in \Omega, \ t \in (0,T], \\ u(x,y,0) = \sqrt{(1-x^{2}/2)(1-y^{2}/2)}, & (x,y) \in \Omega, \\ u(x,y,t) = \sqrt{(1-x^{2}/2)(1-y^{2}/2)}e^{-t}, & (x,y) \in \partial\Omega, \ t \in [0,T], \end{cases}$$

where  $g(x, y, t) = (1 + t^{\mu} + t^{\alpha \mu})e^{(x+y)}(1 + \chi_{(0,1)}(x)).$ 

For subdiffusion PDEs model (1), it is natural to appear the low regularity or singular term such as

$$t^{\mu}f(x,t)$$
 or  $t^{\mu}*f(x,t), \ \mu > -1$ .

In this case, many popular time stepping schemes, including the correction of high-order BDF methods may lose their high-order accuracy, see [11, Section 4.1]

Table 5. The discrete  $L^2$ -norm  $||u^N - u^{2N}||$  and convergent order for Example 6.3 with  $\alpha = 0.7$ .

Scheme	$\mu$	N = 50	N = 100	N = 200	N = 400	N = 800
	0.8	8.4127e-02	2.1353e-02	5.3801e-03	1.3504 e-03	3.3829e-04
ID1-BDF2			1.9781	1.9887	1.9942	1.9971
ID1-DDF 2	-0.8	1.6331e+00	7.6864e-01	3.4752 e-01	1.5385e-01	6.7365 e-02
			1.0872	1.1452	1.1755	1.1915
	0.8	3.0394e-02	7.7886e-03	1.9718e-03	4.9605e-04	1.2442e-04
ID2-BDF2			1.9643	1.9819	1.9909	1.9953
	-0.8	2.4180e-01	6.1541e-02	1.5548e-02	3.9124e-03	9.8238e-04
			1.9742	1.9848	1.9906	1.9937

and Lemma 3.2 in [35], also see Table 1. The correction BDF2 methods recover superlinear convergence order  $\mathcal{O}(\tau^{1+\alpha\mu})$ , provided that the source term behaves like  $t^{\alpha\mu}$ , which is invalid for  $\mu < 0$ , since it is required the source function  $g \in C([0,T];L^2(\Omega))$ .

To fill in this gap, the desired second-order convergence rate can be achieved by ID1-BDF2 with  $\mu>0$  but it is still likely to exhibit an order reduction with  $\mu<0$ . Furthermore, ID2-BDF2 method has filled a gap with  $-1<\mu<0$ , see Tables 1 and 2. Table 3 shows that ID1-BDF2 recovers second order convergence and this is in agreement with the order of convergence for  $t^{\mu}*f(x,t), \ \mu>-1$ .

The proposed methods can be extended to inhomogeneous Dirichlet and Neumann boundary conditions including two-dimensional case, see Examples 6.2 and 6.3.

Remark 6.1. For Hadamard's finite-part integral [6, p. 233]

$$\int_0^t s^{\mu} ds = \frac{1}{1+\mu} t^{1+\mu}, \ \mu < -1,$$

of course the limit does not exist, and so Hadamard suggested simply to ignore the unbounded contribution. In this case, we can similar provide

ID3 – BDF2 Method : 
$$\partial_{\tau}^{\alpha}V^{n} - AV^{n} = \partial_{\tau}^{3} \left(\frac{t_{n}^{3}}{6}Av + \mathbb{G}^{n}\right), \ \mathbb{G} = J^{3}g(x,t),$$

which also recovers the high-order accuracy even for the hypersingul source term, see Table 6.

TABLE 6. The discrete  $L^2$ -norm  $||u^N - u^{2N}||$  and convergent order with  $\beta = 0$ ,  $\alpha = 0.7$ . Here  $\circ$  denotes the dot product.

Scheme	$\mu$	N = 50	N = 100	N = 200	N = 400	N = 800
ID2-BDF2	-1.8	1.7275e-02	8.1527e-03	3.6909e-03	1.6393e-03	7.2110e-04
			1.0834	1.1433	1.1709	1.1848
ID3-BDF2	-1.8	7.7995e-03	1.8929e-03	4.6855e-04	9.5882e-05	2.2325e-05
			2.0428	2.0143	2.2889	2.1026

**Remark 6.2.** It is easy to extend the higher order schemes, e.g., ID2-BDF3, ID3-BDF4, see Table 7.

TABLE 7. The discrete maximum-norm  $||u^N - u^{2N}||$  and convergent order of ID2-BDF3 and ID3-BDF4 scheme for Example 6.1 with g(x,t) = 0.

Scheme	$\alpha$	N = 40	N = 80	N = 160	N = 320	N = 640
	0.3	2.2976e-07	2.7210e-08	3.3127e-09	4.0871e-10	5.0758e-11
ID2-BDF3			3.0779	3.0380	3.0188	3.0093
1D2-DDF3	0.7	7.0505e- $07$	8.2623e-08	$1.0008e ext{-}08$	1.2317e-09	1.5278e-10
			3.0930	3.0453	3.0224	3.0111
	0.3	2.5000e-08	1.4195e-09	8.4663e-11	5.1885e-12	3.9823e-13
ID3-BDF4			4.1384	4.0675	4.0283	3.7036
	0.7	8.5711e-08	4.8005e-09	2.8439e-10	1.7327e-11	1.14674e-12
			4.1582	4.0772	4.0367	3.9174

**6.2. Stochastic subdiffusion model.** In this subsection, we are interested in the ID2-BDF2 method for solving stochastic fractional subdiffusion model

(44) 
$$\begin{cases} {}^{C}D_{t}^{\alpha}u(t) - \Delta u = \partial_{t}^{-\gamma}\frac{dW(t)}{dt} & \text{in } \Omega \times (0,T), \\ u = 0, & \text{on } \partial\Omega \times (0,T), \end{cases}$$

with the initial condition  $u(0) = u_0$ 

**Example 6.4.** Let T=1 and  $\Omega=(0,1)$ . Consider stochastic fractional subdiffusion equation (44) with  $u_0(x)=\sin(x)\sqrt{x(1-x)}$ .

First, an approximation of trace class noise  $\frac{dW(t)}{dt}$  in (44) is defined by

$$\dot{W}_{\ell}(t) = \frac{dW_{\ell}(t)}{dt} = \sum_{j=1}^{\ell} \gamma_j^{1/2} \dot{\beta}_j(t) \varphi_j \text{ with } \gamma_j^{1/2} = j^{-2}$$

where  $\beta_j(t)$  and  $\varphi_j$  (j=1,2,...) are the independently identically distributed Brownian motions and the eigenfunctions of the operator A, respectively. In particular, in the one-dimensional case, we have  $\varphi_j(x) = \sqrt{2}\sin(j\pi x)$ , j=1,2,...

Let  $V(t) = u(t) - u_0$  with V(0) = 0. Substituting  $\frac{dW_{\ell}(t)}{dt}$  for  $\frac{dW(t)}{dt}$  in (44), we obtain, with  $V_{\ell}(0) = 0$ ,

(45) 
$$\partial_t^{\alpha} V_{\ell}(t) - A V_{\ell}(t) = \partial_t \left( t A u_0 \right) + \partial_t^{2-\gamma} \left( t * \frac{dW_{\ell}(t)}{dt} \right).$$

Next, we design the ID2-BDF2 time discretization scheme of (45) as following

$$(46) \qquad \partial_{\tau}^{\alpha} V_{\ell}^{n} - A V_{\ell}^{n} = \partial_{\tau} \left( t_{n} A u_{0} \right) + \partial_{\tau}^{2-\gamma} \left( t_{n} * \frac{dW_{\ell}(t_{n})}{dt} \right), \quad n = 1, 2, \dots, N.$$

Using BDF2 integrals convolution quadrature formula [1, 18], it yields

$$t_n * \frac{dW_{\ell}(t_n)}{dt} = \partial_t^{-2} \frac{dW_{\ell}(t_n)}{dt} = \partial_t^{-1} W_{\ell}(t_n) \approx \overline{\tau} \sum_{k=1}^{\overline{n}} w_{\overline{n}-k}^{(-1)} \sum_{j=1}^{\ell} j^{-2} \beta_j(\overline{t}_k) \varphi_j$$

with  $\ell=100$ . Here Brownian motions  $\{\beta_j\}_{j=1}^\ell$  can be generated by MATLAB code, see [28, p. 395]. The time step size  $\overline{\tau}=2^{-20}$  and  $t_n=n\tau=n\frac{1}{N}=n\frac{\overline{N}}{N}\overline{\tau}=\overline{n}\,\overline{\tau}=\overline{t}_{\overline{n}}$  with  $\overline{n}=n\overline{N}/N$  and  $\overline{N}=2^{20}$ . All the expected values are computed with 1000 trajectories.

Table 8 shows that the ID2-BDF2 method is able to achieve superlinear convergence rate  $O(\tau^{\alpha+\gamma-1/2})$  for  $1/2 < \alpha + \gamma < 2$  in solving the stochastic fractional

TABLE 8. The discrete  $L^2$ -norm  $||u^N - u^{2N}||$  and convergent order of ID2-BDF2 schemes (46) with  $\gamma = 0.9$ .

$\alpha$	N = 64	N = 128	N = 256	N = 512
0.1	9.7379e-03	7.4798e-03	5.6008e-03	4.1725e-03
0.1		0.3806	0.4173	0.4247
0.5	3.5940 e-03	2.1200e-03	1.2159e-03	6.9409e-04
0.5		0.7615	0.8020	0.8087
0.9	3.1885e-04	1.1245e-04	4.1985e-05	1.6693e-05
0.9		1.5035	1.4213	1.3306

subdiffusion models, which breaks the first-order barrier in [14]. The convergence theory of stochastic differential equations is being research.

## 7. Conclusions

Fractional PDEs model naturally imply a less smooth or low regularity source function  $t^{\mu} \circ f(x,t)$  in the right-hand side, which is likely to result in a severe order reduction in most existing time-stepping schemes. To fill in this gap, we provide a new idea to obtain the second-order time-stepping schemes for subdiffusion, where the source term is regularized by using an m-fold integral-derivative and the equation is discretized by using a modified BDF2 convolution quadrature. The detailed theoretical analysis and numerical verifications are presented. In the future studies, we will try to adapt the idea to higher order schemes [11], the nonlinear fractional models [16], and the stochastic fractional evolution model [12].

# Acknowledgments

This work was supported by the National Natural Science Foundation of China under Grant No.12471381 and Science Fund for Distinguished Young Scholars of Gansu Province under Grant No. 23JRRA1020. The work of Z. Zhou is partly supported by Hong Kong Research Grants Council (15303122) and an internal grant of Hong Kong Polytechnic University (Project ID: P0031041, Work Programme: ZZKS).

## References

- M.H. Chen and W.H. Deng, Discretized fractional substantial calculus, ESAIM: Math. Mod. Numer. Anal., 49 (2015), pp. 373–394.
- [2] M.H. Chen and W.H. Deng, High order algorithms for the fractional substantial diffusion equation with truncated Lévy flights, SIAM J. Sci. Comput., 37 (2015), pp. A890–A917.
- [3] M.H. Chen, S.Z. Jiang and W.P. Bu, Two L1 schemes on graded meshes for fractional Feynman-Kac equation, J. Sci. Comput., 88 (2021), No. 58.
- [4] S. Chen, J. Shen, Z. Zhang and Z. Zhou, A spectrally accurate approximation to subdiffusion equations using the log orthogonal functions, SIAM J. Sci. Comput., 42 (2020), pp. A849– A877.
- [5] E. Cuesta, Ch. Lubich and C. Palencia, Convolution quadrature time discretization of fractional diffusion-wave equations, Math. Comp., 75 (2006), pp. 673–696.
- [6] K. Diethelm, The Analysis of Fractional Differential Equations, Springer, 2010.
- [7] J.S. Hesthaven and T. Warburton, Nodal Discontinuous Galerkin Methods: Algorithms, Analysis, and Applications, Springer, 2007.
- [8] D. Hou and C. Xu, A fractional spectral method with applications to some singular problems, Adv. Comput. Math., 43 (2017), pp. 911–944.
- [9] B. Jin, R. Lazarov and Z. Zhou, An analysis of the L1 scheme for the subdiffusion equation with nonsmooth data, IMA J. Numer. Anal., 36 (2016), pp. 197–221.

- [10] B. Jin, R. Lazarov and Z. Zhou, Numerical methods for time-fractional evolution equations with nonsmooth data: a concise overview, Comput. Methods Appl. Mech. Engrg., 346 (2019), pp. 332–358.
- [11] B. Jin, B.Y. Li and Z. Zhou, Correction of high-order BDF convolution quadrature for fractional evolution equations, SIAM J. Sci. Comput., 39 (2017), pp. A3129–A3152.
- [12] B. Jin, Y. Yan and Z. Zhou, Numerical approximation of stochastic time-fractional diffusion, ESAIM Math. Model. Numer. Anal., 53 (2019), pp. 1245–1268.
- [13] B. Jin and Z. Zhou, Numerical treatment and analysis of time-fractional evolution equations, Springer, Cham, 2023.
- [14] W. Kang, B. A. Egwu, Y.B. Yan and A. K. Pani, Galerkin finite element approximation of a stochastic semilinear fractional subdiffusion with fractionally integrated additive noise, IMA J. Numer. Anal., 42 (2022), pp. 2301–2335.
- [15] N. Kopteva, Error analysis of an L2-type method on graded meshes for a fractional-order parabolic problem, Math. Comp., 90 (2021), pp. 19–40.
- [16] W. Li and A. Salgad, Time fractional gradient flows: Theory and numerics, Math. Models Methods Appl. Sci., 33 (2023), pp. 377–453.
- [17] H.-L. Liao, D. Li and J. Zhang, Sharp error estimate of the nonuniform L1 formula for linear reaction-subdiffusion equations, SIAM J. Numer. Anal., 56 (2018), pp. 1112–1133.
- [18] Ch. Lubich, Discretized fractional calculus, SIAM J. Math. Anal., 17 (1986), pp. 704–719.
- [19] Ch. Lubich, I.H. Sloan and V. Thomée, Nonsmooth data error estimates for approximations of an evolution equation with a positive-type memory term, Math. Comp., 65 (1996), pp. 1–17.
- [20] W. McLean and K. Mustapha, Time-stepping error bounds for fractional diffusion problems with non-smooth initial data, J. Comput. Phys., 293 (2015), pp. 201–217.
- [21] W. McLean, K. Mustapha, R. Ali and O. Knio, Well-posedness of time-fractional advectiondiffusion-reaction equations, Fract. Calc. Appl. Anal., 22 (2019), pp. 918–944.
- [22] R. Metzler, Brownian motion and beyond: first-passage, power spectrum, non-Gaussianity, and anomalous diffusion, J. Stat. Mech., (2019), 114003.
- [23] R. Metzler and J. Klafter, The random walk's guide to anomalous diffusion: a fractional dynamics approach, Phys. Rep., 339 (2000), pp. 1–77.
- [24] K. Mustapha, An L1 approximation for a fractional reaction-diffusion equation, a second-order error analysis over time-graded meshes, SIAM J. Numer. Anal., 58 (2020), pp. 1319–1338
- [25] K. Mustapha, B. Abdallah and K.M. Furati, A discontinuous Petrov-Galerkin method for time-fractional diffusion equations, SIAM J. Numer. Anal., 52 (2014), pp. 2512–2529.
- [26] I. Podlubny, Fractional Differential Equations, Academic Press, 1999.
- [27] K. Sakamoto and M. Yamamoto, Initial value/boundary value problems for fractional diffusion-wave equations and applications to some inverse problems, J. Math. Anal. Appl., 382 (2011), pp. 426–447.
- [28] V. Schmidt, Stochastic Geometry, Spatial Statistics and Random Fields, Springer, Cham, 2015.
- [29] R. Schumer, D.A. Benson, M.M. Meerschaert and B. Baeumer, Fractal mobile/immobile solute transport, Water Resour. Res., 39 (2003), pp. 1–12.
- [30] J. Shen, T. Tang and L. Wang, Spectral Methods: Algorithms, Analysis and Applications, Springer, 2011.
- [31] J. Shen, F. Zeng and M. Stynes, Second-order error analysis of the averaged L1 scheme  $\overline{L1}$  for time-fractional initial-value and subdiffusion problems, Sci. China Math., 67 (2024), pp. 1641–1664.
- [32] J.K. Shi and M.H. Chen, Correction of high-order BDF convolution quadrature for fractional Feynman-Kac equation with Lévy flight, J. Sci. Comput., 85 (2020), No. 28.
- [33] M. Stynes, E. O'riordan and J.L. Gracia, Error analysis of a finite difference method on graded meshes for a time-fractional diffusion equation, SIAM J. Numer. Anal., 55 (2017), pp. 1057–1079.
- [34] V. Thomée, Galerkin Finite Element Methods for Parabolic Problems, Springer, 2006.
- [35] K. Wang and Z. Zhou, High-order time stepping schemes for semilinear subdiffusion equations, SIAM J. Numer. Anal., 58 (2020), pp. 3226–3250.
- [36] Y.Y. Wang, Y.B. Yan and Y. Yang, Two high-order time discretization schemes for subdiffusion problems with nonsmooth data, Fract. Calc. Appl. Anal., 23 (2020), pp. 1349–1380.

- [37] Y.B. Yan, M. Khan and N.J. Ford, An analysis of the modified L1 scheme for time-fractional partial differential equations with nonsmooth data, SIAM J. Numer. Anal., 56 (2018), pp. 210-227.
- [38] M. Zayernouri and G.E. Karniadakis, Fractional Sturm-Liouville eigen-problems: theory and numerical approximation, Comput. Phys., 252 (2013), pp. 495–517.
- [39] H. Zhou and W.Y. Tian, Two time-stepping schemes for sub-diffusion equations with singular source terms, J. Sci. Comput., 92 (2022), No. 70.

School of Mathematics and Statistics, Gansu Key Laboratory of Applied Mathematics and Complex Systems, Lanzhou University, Lanzhou 730000, People's Republic of China.

 $E ext{-}mail: \texttt{chenmh@lzu.edu.cn}$ 

School of Mathematics and Computational Science, Xiangtan University, Xiangtan 411105, Hunan, People's Republic of China.

 $E ext{-}mail: shijk@xtu.edu.cn}$ 

Department of Applied Mathematics, The Hong Kong Polytechnic University, Kowloon, Hong Kong, People's Republic of China.

 $E ext{-}mail: {\tt zhizhou@polyu.edu.hk}$