

Solutions to assignment #5

8.2.1 e) equilibrium satisfies $\frac{d}{dt}v(x)=0$, hence

$$\left. \begin{array}{l} kv_{xx} = -k \\ v_x(0) = 0 \\ v_x(L) = 0 \end{array} \right\} \Rightarrow \begin{array}{l} v(x) = -\frac{x^2}{2} + ax + b \\ v_x(x) = -x + a \end{array}$$

$$v_x(0) = 0 = a \Rightarrow a = 0$$

$$v_x(L) = -L \neq 0 \Rightarrow \text{there exists no equilibrium.}$$

To find a problem with homogeneous boundary values, I don't need to do anything, since it has homogeneous b.c. already. But to make the problem homogeneous I can write

$$u(x,t) = w(x,t) + kt, \text{ which gives}$$

$$\left. \begin{array}{l} w_t = kw_{xx} \\ w_x(0) = 0 = w_x(L) \end{array} \right\}.$$

$$\left. \begin{array}{l} \underline{8.2.2.b)} \quad u_t = ku_{xx} + Q(x,t) \\ u(x,0) = f(x) \\ u(0,t) = A(t) \\ u_x(L,t) = B(t) \end{array} \right\}$$

$$\text{try: } u(x,t) = v(x,t) + w(x,t).$$

$$\left. \begin{array}{l} v_t + w_t = kv_{xx} + kw_{xx} + Q \\ v(x,0) + w(x,0) = f(x) \end{array} \right\} \begin{array}{l} v(0,t) + w(0,t) = A(t) \\ v_x(L,t) + w_x(L,t) = B(t) \end{array}$$

v-problem

$$v(0, t) = A(t)$$

$$v_x(L, t) = B(t)$$

w-problem

$$w_t = k w_{xx} + k v_{xx} - v_t + Q$$

$$w(x, 0) = f(x) - v(x, 0)$$

$$w(0, t) = 0$$

$$w_x(L, t) = 0$$

Solve the v-problem: $v(x, t) = A(t) + B(t)x$

Then the remaining problem with homogeneous boundary conditions is

$$w_t = k w_{xx} - \dot{A}(t) - \dot{B}(t)x + Q(x, t)$$

$$w(x, 0) = f(x) - A(0) - B(0)x$$

$$w(0, t) = 0$$

$$w_x(L, t) = 0$$

8.2.3

Step 1: Make homogeneous boundary conditions:

$$u(r, t) = v(r) + w(r, t)$$

$$w_t = \frac{k}{r} \left(\frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) \right) + Q(r)$$

$$v(r) + w(r, 0) = f(r)$$

$$v(a) + w(a, t) = T$$

v-problem

$$\left. \begin{aligned} \frac{k}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v}{\partial r} \right) &= 0 \\ v(a) &= T \end{aligned} \right\}$$

w-problem

$$\left. \begin{aligned} w_t &= \frac{k}{r} \frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) + Q \\ w(r, 0) &= f(r) - v(r) \\ w(a, t) &= 0 \end{aligned} \right\}$$

Solve v-problem: $v(r) = T$ (constant).

Then $u(r, t) = w(r, t) + T$ and

$$\left. \begin{aligned} w_t &= \frac{k}{r} \left(\frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) \right) + Q(r) \\ w(r, 0) &= f(r) - T \\ w(a, t) &= 0 \end{aligned} \right\} (*)$$

Step 2

look at homogeneous problem:

$$\left. \begin{aligned} w_t &= \frac{k}{r} \left(\frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) \right) \\ w(a, t) &= 0, \\ |w(0, t)| &< \infty \end{aligned} \right\}$$

Separation: $w(r, t) = R(r) H(t)$

$$R H' = \frac{k}{r} \left((r R')' \right) H$$

$$\frac{H'}{k H} = \frac{1}{r R} \left((r R')' \right) = -\lambda$$

$$H' = -\lambda k H,$$

$$r(r R')' + \lambda r^2 R = 0, \quad R(a) = 0, \quad |R(0)| < \infty$$

The R-problem is a Bessel equation of order $m=0$. The solution is $J_0(\sqrt{\lambda_{0n}} r)$, where λ_{0n} are given as $\sqrt{\lambda_{0n}} a = z_{0n}$, which are the zeroes of $J_0(z)$.

Step 3 To solve (*) we assume

$$w(r, t) = \sum_{n=1}^{\infty} a_n(t) J_0(\sqrt{\lambda_{0n}} r).$$

We need the Bessel series of $Q(r)$ and of $f(r) - T$:

$$Q(r) = \sum_{n=1}^{\infty} q_n J_0(\sqrt{\lambda_{0n}} r) \quad q_n = \frac{\int_0^a J_0(\sqrt{\lambda_{0n}} r) Q(r) r dr}{\int_0^a J_0^2(\sqrt{\lambda_{0n}} r) r dr}$$

$$f(r) - T = \sum_{n=1}^{\infty} f_n J_0(\sqrt{\lambda_{0n}} r) \quad f_n = \frac{\int_0^a J_0(\sqrt{\lambda_{0n}} r) (f(r) - T) r dr}{\int_0^a J_0^2(\sqrt{\lambda_{0n}} r) r dr}$$

Substitute into (*):

$$\begin{aligned} \sum_{n=1}^{\infty} \dot{a}_n(t) J_0(\sqrt{\lambda_{0n}} r) &= \sum_{n=1}^{\infty} a_n(t) (-\lambda_{0n}) J_0(\sqrt{\lambda_{0n}} r) \\ &\quad + \sum_{n=1}^{\infty} q_n J_0(\sqrt{\lambda_{0n}} r) \end{aligned}$$

Compare coefficients

$$\dot{a}_n(t) = -\lambda_{0n} a_n(t) + q_n$$

$$\begin{aligned}
 a_n(t) &= a_n(0)e^{-\lambda_n t} + \int_0^t e^{-\lambda_n(t-s)} q_n ds \\
 &= a_n(0)e^{-\lambda_n t} + \frac{q_n}{\lambda_n} (1 - e^{-\lambda_n t})
 \end{aligned}$$

$$\begin{aligned}
 \text{At } t=0: \quad w(r, 0) &= \sum_{n=1}^{\infty} a_n(0) J_0(\sqrt{\lambda_n} r) \\
 &= f(r) - T \\
 &= \sum_{n=1}^{\infty} f_n J_0(\sqrt{\lambda_n} r)
 \end{aligned}$$

$$\Rightarrow a_n(0) = f_n.$$

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Step 4 Put everything together:

$$\begin{aligned}
 u(r, t) &= T + \sum_{n=1}^{\infty} \left[f_n e^{-\lambda_n t} - \frac{q_n}{\lambda_n} e^{-\lambda_n t} + \frac{q_n}{\lambda_n} \right] J_0(\sqrt{\lambda_n} r) \\
 &= T + \sum_{n=1}^{\infty} \left(\frac{q_n}{\lambda_n} + \left(f_n - \frac{q_n}{\lambda_n} \right) e^{-\lambda_n t} \right) J_0(\sqrt{\lambda_n} r)
 \end{aligned}$$

8.3.2 Homogeneous problem

$$\left. \begin{aligned} u_t &= k u_{xx} \\ u(0, t) &= 0 = u(L, t) \end{aligned} \right\} \text{ has eigenvalues } \lambda_n = \left(\frac{n\pi}{L}\right)^2$$

and eigenfunctions $\phi_n = \sin\left(\frac{n\pi x}{L}\right)$

Then we try: $u(x, t) = \sum_{n=1}^{\infty} a_n(t) \sin\left(\frac{n\pi x}{L}\right)$.

We will need the same series expansions for Q and f :

$$Q(x) = \sum_{n=1}^{\infty} q_n \sin\left(\frac{n\pi x}{L}\right)$$

$$q_n = \frac{2}{\pi} \int_0^L Q(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

$$f(x) = \sum_{n=1}^{\infty} f_n \sin\left(\frac{n\pi x}{L}\right)$$

$$f_n = \frac{2}{\pi} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

From the non-homogeneous problem we obtain

$$\begin{aligned} \sum_{n=1}^{\infty} a_n'(t) \sin\left(\frac{n\pi x}{L}\right) &= \sum_{n=1}^{\infty} a_n(t) k \left(-\frac{n\pi}{L}\right)^2 \sin\left(\frac{n\pi x}{L}\right) \\ &\quad + \sum_{n=1}^{\infty} q_n \sin\left(\frac{n\pi x}{L}\right) \end{aligned}$$

Compare coefficients

$$a_n'(t) = -k \left(\frac{n\pi}{L}\right)^2 a_n + q_n$$

with initial condition $a_n(0) = f_n$.

The solution is

$$a_n(t) = f_n e^{-k \left(\frac{n\pi}{L}\right)^2 t} + q_n \int_0^t e^{-k \left(\frac{n\pi}{L}\right)^2 (t-s)} ds$$

$$a_n(t) = f_n e^{-k(\frac{n\pi}{L})^2 t} + \frac{q_n}{k(\frac{n\pi}{L})^2} (1 - e^{-k(\frac{n\pi}{L})^2 t})$$

$$= \frac{q_n}{k(\frac{n\pi}{L})^2} + e^{-k(\frac{n\pi}{L})^2 t} \left(f_n - \frac{q_n}{k(\frac{n\pi}{L})^2} \right)$$

$$\xrightarrow{t \rightarrow \infty} \frac{q_n}{k(\frac{n\pi}{L})^2}$$

Solution

$$u(x, t) = \sum_{n=1}^{\infty} \left[\frac{q_n}{k(\frac{n\pi}{L})^2} + e^{-k(\frac{n\pi}{L})^2 t} \left(f_n - \frac{q_n}{k(\frac{n\pi}{L})^2} \right) \right] \sin\left(\frac{n\pi x}{L}\right)$$

$$\xrightarrow{t \rightarrow \infty} \sum_{n=1}^{\infty} \frac{q_n}{k(\frac{n\pi}{L})^2} \sin\left(\frac{n\pi x}{L}\right)$$

$=: r(x)$ which is the steady state.

$$\begin{aligned} \underline{10.3.11} \quad a) \mathcal{F}(c_1 f + c_2 g) &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} (c_1 f + c_2 g) e^{i\omega x} dx \\ &= \frac{c_1}{2\pi} \int_{-\infty}^{+\infty} f e^{i\omega x} dx + \frac{c_2}{2\pi} \int_{-\infty}^{+\infty} g e^{i\omega x} dx \\ &= c_1 \mathcal{F}(f) + c_2 \mathcal{F}(g) \end{aligned}$$

$$b) \mathcal{F}^{-1}(\mathcal{F}(\omega) G(\omega)) = \frac{1}{2\pi} f * g \neq f \cdot g.$$

10.3.5] $f(x), \hat{f}(\omega) = \mathcal{F}(f(x))$

$$\mathcal{F}^{-1}(e^{i\omega\beta} \hat{f}(\omega)) = \int_{-\infty}^{+\infty} \hat{f}(\omega) e^{i\omega\beta} e^{-i\omega x} d\omega$$

$$= \int_{-\infty}^{+\infty} \hat{f}(\omega) e^{-i\omega(x-\beta)} d\omega$$

$$= f(x-\beta)$$

10.4.4]

$$\left. \begin{aligned} u_t &= k u_{xx} - \gamma u & -\infty < x < \infty \\ u(x, 0) &= f(x) \end{aligned} \right\}$$

call $\hat{u}(\omega, t) = \mathcal{F}(u(x, t))$

$$\hat{u}_t = -k\omega^2 \hat{u} - \gamma \hat{u}$$

$$\hat{u}_t = (-k\omega^2 - \gamma) \hat{u} \Rightarrow \hat{u}(\omega, t) = \hat{f}(\omega) e^{-(k\omega^2 + \gamma)t}$$

$$\hat{u}(\omega, 0) = \hat{f}(\omega)$$

$$u(x, t) = \mathcal{F}^{-1}(\hat{u}(\omega, t)) = \mathcal{F}^{-1}(\hat{f}(\omega) e^{-(k\omega^2 + \gamma)t})$$

$$= \mathcal{F}^{-1}(\hat{f}(\omega) e^{-k\omega^2 t} e^{-\gamma t}) \quad \mathcal{F}^{-1} \text{ is linear:}$$

$$= e^{-\gamma t} \mathcal{F}^{-1}(\hat{f}(\omega) e^{-k\omega^2 t})$$

$$= e^{-\gamma t} \frac{1}{2\pi} f * g$$

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$$\Rightarrow u(x,t) = e^{-\gamma t} \int_{-\infty}^{+\infty} f(y) \frac{e^{-\frac{(x-y)^2}{4kt}}}{\sqrt{4\pi kt}} dy$$

We find the transformation if we multiply by $e^{\gamma t}$.

$e^{\gamma t} u(x,t) = f * G$ looks like a solution of a homogeneous heat equation. Indeed, if we define $w(x,t) := e^{\gamma t} u(x,t)$, we

$$\begin{aligned} \text{get } w_t &= \gamma e^{\gamma t} u + e^{\gamma t} u_t \\ &= \gamma w + e^{\gamma t} (k u_{xx} - \gamma u) \\ &= \gamma w + k w_{xx} - \gamma w \end{aligned}$$

$$\text{hence } \left. \begin{aligned} w_t &= k w_{xx} \\ w(x,0) &= f(x) \end{aligned} \right\} \quad -\infty < x < \infty$$

10.3.5 $\hat{f}(\omega) = \mathcal{F}(f(x))$

$$\begin{aligned}\mathcal{F}^{-1}(e^{i\omega\beta} \hat{f}(\omega)) &= \int_{-\infty}^{+\infty} \hat{f}(\omega) e^{i\omega\beta} e^{-i\omega x} d\omega \\ &= \int_{-\infty}^{+\infty} \hat{f}(\omega) e^{-i\omega(x-\beta)} d\omega \\ &= f(x-\beta)\end{aligned}$$

10.4.4.1 $u_t = ku_{xx} - \gamma u \quad -\infty < x < \infty$
 $u(x, 0) = f(x)$

a) Use $\mathcal{F}$ -transform: $\hat{u}(\omega, t) = \mathcal{F}(u(x, t))$

$$\mathcal{F}(u_t) = \hat{u}_t \quad \mathcal{F}(u_{xx}) = -\omega^2 \hat{u}$$

$$\begin{aligned}\hat{u}_t &= -k\omega^2 \hat{u} - \gamma \hat{u} \\ &= -(k\omega^2 + \gamma) \hat{u}\end{aligned}$$

which is solved by

$$\hat{u}(\omega, t) = \hat{u}(\omega, 0) e^{-(k\omega^2 + \gamma)t}$$

$$\hat{u}(\omega, 0) = \hat{f}(\omega)$$

hence $\hat{u}(\omega, t) = \hat{f}(\omega) e^{-(k\omega^2 + \gamma)t}$

$$u(x, t) = \mathcal{F}^{-1} \left(\hat{f}(\omega) e^{-(k\omega^2 + \gamma)t} \right)$$

$$= \mathcal{F}^{-1} \left(\hat{f}(\omega) e^{-k\omega^2 t} e^{-\gamma t} \right)$$

$$= e^{-\gamma t} \mathcal{F}^{-1} \left(\hat{f}(\omega) e^{-k\omega^2 t} \right)$$

solution of heat equation
without γ .

$$= e^{-\gamma t} \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{+\infty} f(y) e^{-\frac{(x-y)^2}{4kt}} dy //$$

b) We multiply by $e^{\gamma t}$:

$$e^{\gamma t} u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{+\infty} f(y) e^{-\frac{(x-y)^2}{4kt}} dy ,$$

which looks like a solution of the heat eq.

Indeed if we transform $w(x, t) = e^{\gamma t} u(x, t)$
then we get

$$\frac{\partial}{\partial t} w(x, t) = \gamma e^{\gamma t} u + e^{\gamma t} u_t$$

$$= \gamma e^{\gamma t} u + e^{\gamma t} (k u_{xx} - \gamma u)$$

$$= k e^{\gamma t} u_{xx}$$

$$= k w_{xx} //$$

$$\underline{10.5.12} \quad \text{Solve } \left. \begin{aligned} u_t &= k u_{xx} & x > 0 \\ u_x(0, t) &= 0 \\ u(x, 0) &= f(x) \end{aligned} \right\}$$

Use cosine-transform: $\tilde{u}(\omega, t) = C(u(x, t))$

$$C(u_t) = \tilde{u}_t$$

$$C(u_{xx}) = -\frac{2}{\pi} u_x(0, t) - \omega^2 \tilde{u}$$

||
0 from boundary condition.

$$\Rightarrow \tilde{u}_t = -k\omega^2 \tilde{u}$$

$$\tilde{u}(\omega, t) = \tilde{u}(\omega, 0) e^{-k\omega^2 t}$$

$$= \tilde{f}(\omega) e^{-k\omega^2 t} \quad \tilde{f}(\omega) = C(f(x)).$$

$$\Rightarrow u(x, t) = \int_0^{\infty} \tilde{f}(\omega) e^{-k\omega^2 t} \cos \omega x \, d\omega$$

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even

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$$= \frac{1}{2} \int_{-\infty}^{+\infty} \tilde{f}(\omega) e^{-k\omega^2 t} (\cos \omega x - i \sin \omega x) \, d\omega$$

$$= \int_{-\infty}^{+\infty} \frac{\tilde{f}(\omega)}{2} e^{-k\omega^2 t} e^{-i\omega x} \, d\omega$$

$$= \mathcal{F}^{-1} \left(\frac{\tilde{f}(\omega)}{2} e^{-k\omega^2 t} \right) \quad (*)$$

$$\frac{\gamma(\omega)}{2} = \frac{1}{2\pi} \int_0^{\infty} f(x) \cos(\omega x) dx \quad \text{use even extension}$$

$$= \frac{1}{2\pi} \int_{-\infty}^{+\infty} f_{\text{even}}(x) (\cos(\omega x) + i \sin \omega x) dx$$

$$= \frac{1}{2\pi} \int_{-\infty}^{+\infty} f_{\text{even}}(x) e^{i\omega x} dx$$

$$= \mathcal{F}(f_{\text{even}}(x)) \quad (**)$$

From (*) and (**) it follows that $u(x, t)$ is a solution of
$$\left. \begin{aligned} u_t &= k u_{xx} & -\infty < x < \infty \\ u(x, 0) &= f_{\text{even}}(x) \end{aligned} \right\}$$

Hence $u(x, t) = f_{\text{even}} * G$

$$= \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{+\infty} f_{\text{even}}(y) e^{-\frac{(x-y)^2}{4kt}} dy$$

$$= \frac{1}{\sqrt{4\pi kt}} \left(\int_0^{\infty} f(y) \left(e^{-\frac{(x+y)^2}{4kt}} + e^{-\frac{(x-y)^2}{4kt}} \right) dy \right)$$